

The characteristics of the 2022 Tonga volcanic tsunami in the Pacific Ocean

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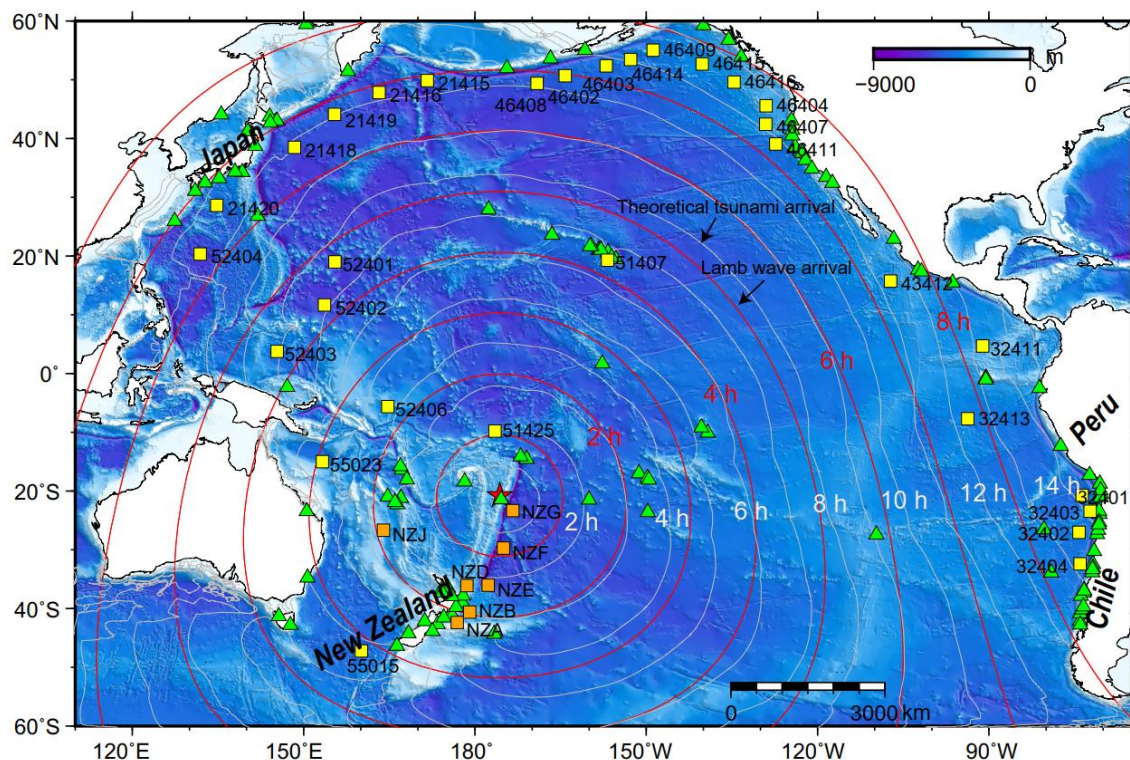
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Abstract. On 15th January 2022, an exceptional eruption of Hunga Tonga–Hunga Ha’apai volcano generated atmospheric and tsunami waves that were widely observed at oceans globally, gaining a remarkable attention to scientists in related fields. The tsunamigenic mechanism of this rare event remains an enigmatic due to its complexity and lacking of direct underwater observations. Here, to explore the tsunamigenic mechanisms of this volcanic tsunami event and its hydrodynamic processes in the Pacific Ocean, we conduct tsunami waveform and spectral analyses of the waveform recordings at 116 coastal gauges and 38 deep-ocean buoys across the Pacific Ocean. Combined with the constraints of some representative barometers, we obtain the plausible tsunamigenic origins during the volcano activity. We identify four distinct tsunami wave components generated by air-sea coupling and seafloor crustal deformation. Those tsunami components are differentiated by their different propagating speeds or period bands. The first-arriving tsunami component with ~80–100 min period was from shock waves spreading at a velocity of ~1000 m/s in vicinity of the eruption. The second component with extraordinary tsunami amplitude in deep sea was from Lamb waves. The Lamb wave with ~30–40 min period radically propagated outward from the eruption site with spatially decreasing propagation velocities from ~340 m/s to ~315m/s. The third component with ~10–30 min period was probably from some atmospheric gravity wave modes propagating faster than 200 m/s but slower than Lamb waves. The last component with ~3–5 min period originated from partial caldera collapse with dimension of ~0.8–1.8 km. Surprisingly, the 2022 Tonga volcanic tsunami produced long oscillation in the Pacific Ocean which is comparable with those of the 2011 Tohoku tsunami. We point out that the long oscillation is not only associated with the resonance effect with the atmospheric acoustic-gravity waves, but more importantly the interactions with local bathymetry. This rare event also calls for more attention to the tsunami hazards produced by atypical tsunamigenic source, e.g., volcanic eruption.

30 **1. Introduction**

31 On 15 January 2022 at 04:14:45 (UTC), a submarine volcano erupted violently at the uninhabited Hunga
32 Tonga-Hunga Ha'apai (HTHH) island at 20.546°S 175.390°W (USGS, 2022). The volcano is located ~67
33 km north of Nuku'alofa, the capital of Tonga (NASA, 2022) (Figure 1). The blasts launched plumes of
34 ash, steam, and gas ~58 km high into stratosphere (Yuen et al., 2022) which not only blanketed nearby
35 islands in ash (Duncombe, 2022; NASA, 2022), but caused various atmospheric acoustic-gravity wave
36 modes (AGWs) of various scales, e.g., Lamb waves from atmospheric surface pressure disturbance
37 associated with the eruption (Liu and Higuera, 2022; Adam, 2022; Kubota et al., 2022; Matoza et al.,
38 2022). Tsunami with conspicuous sea level changes were detected by coastal tide gauges and Deep-ocean
39 Assessment and Reporting of Tsunamis (DART) buoy stations in the Pacific (Figure 1), the Atlantic, and
40 Indian Oceans as well as the Caribbean and Mediterranean seas (Carvajal et al., 2022; Kubota et al., 2022;
41 Ramírez-Herrera et al., 2022), while the large waves were mainly concentrated in the Pacific Ocean, like
42 coastlines of New Zealand, Japan, California, and Chile (Carvajal et al., 2022). The event caused at least
43 3 fatalities in Tonga. Two people drowned in northern Peru when ~2 m destructive tsunami waves
44 inundated an island in the Lambayeque region, Chile (Edmonds, 2022).

45 Satellite images revealed that the elevation of HTHH island has gone through dramatic change before
46 and after the mid-January 2022 eruption. Previously, after the 2015 eruption, the two existing Hunga
47 Tonga and Hunga Ha'apai Islands were linked together. The volcanic island rose 1.8 km from the seafloor
48 where it stretched ~20 km across and topped a underwater caldera ~5 km in diameter (Garvin et al., 2018;
49 NASA, 2022). After the violent explosion on 15 January 2022, the newly formed island during 2015 was
50 completely gone, with only small tips left in far southwestern and northeastern HTHH island (NASA,
51 2022). HTHH volcano lies along the northern part of Tonga–Kermadec arc, where the Pacific Plate
52 subducts under the Indo-Australian Plate (Billen et al., 2003). The convergence rate (15~24 cm/year)
53 between the Tonga-Kermadec subduction system and the Pacific plate is among the fastest recorded plate
54 velocity on Earth, forming the second deepest trench around the globe (Satake, 2010; Bevis et al., 1995).
55 The fast convergence rate contributes to the frequent earthquakes, tsunamis and volcanic eruptions in
56 this region historically (Bevis et al., 1995). The 2022 HTHH volcano is part of a submarine-volcano
57 chain that extends all the way from New Zealand to Fiji (Plank et al., 2020). HTHH volcano had many
58 notable eruptions before 2022 since its first historically recorded eruption in 1912, i.e., in 1937, 1988,



60
 61 **Figure 1. The spatial distribution of the eruption site (red star), DART stations (squares), tide**
 62 **gauges (triangles) and the calculated tsunami arrival times. White contours indicate the modelled**
 63 **arrival times of conventional tsunami. Red contours indicate the estimated arrival times of Lamb**
 64 **waves (see how we derive these contours in section 3.1).**

65 The 2022 HTHH eruption is the first volcanic event which generates worldwide tsunami signatures since
 66 the 1883 Krakatau event (Matoza et al., 2022; Self and Rampino, 1981; Nomanbhoy and Satake, 1995).

67 The tsunamigenic mechanism of this rare volcanic eruption-induced tsunami is still poorly understood
 68 due to its complex nature and the deficiencies of near-field seafloor surveys. Various tsunami generation
 69 mechanisms have been proposed so far based on the observations of ground-based and spaceborne
 70 geophysical instrumentations (Kubota et al., 2022; Matoza et al., 2022; Carvajal et al., 2022). The
 71 mechanisms are closely associated with the air-sea coupling with atmospheric waves. Atmospheric
 72 waves propagating in the atmospheric fluid are generated by different physical mechanisms (Gossard
 73 and Hooke, 1975a). Lamb wave is a horizontally propagating acoustic waves in Lamb mode
 74 which is trapped at the earth's surface with group velocities close to the mean sound velocity of
 75 the lower atmosphere (e.g. Lamb, 1932). Atmospheric gravity wave is triggered when air
 76 molecules in the atmosphere are disturbed vertically other than horizontally (e.g. Le Pichon et
 77 al., 2010). Nonlinear propagation of atmospheric wave may cause period lengthening and the

78 formation of shock-wave (Matoza et al., 2022). The most-mentioned mechanism of the tsunami is
79 the fast-traveling atmospheric Lamb wave generated by the atmospheric pressure rise of ~ 2 hPa during
80 the eruption. The Lamb wave circled the Earth for several times with travelling speed close to that of the
81 sound wave in the lower atmosphere, leading to globally observed sea level fluctuations (Adam, 2022;
82 Duncombe, 2022; Kubota et al., 2022; Matoza et al., 2022) (Figure 1). The second mechanism is
83 suggested to be a variety of other acoustic-gravity wave modes (Adam, 2022; Matoza et al., 2022;
84 Themens et al., 2022; Zhang et al., 2022). The third mechanism may be related to the seafloor crustal
85 deformation induced by one or more volcanic activities in the vicinity of the eruption site (e.g.,
86 pyroclastic flows, partial collapse of the caldera) (Carvajal et al., 2022), which are more responsible for
87 the near-field tsunamis with theoretical tsunami speeds.

88 To investigate the possible tsunamigenic mechanisms and detailed hydrodynamic behaviors of this rare
89 volcanic tsunami event, in this study, we collect, process and analyze the sea level measurements from
90 116 tide gauge and 38 DART buoys in the Pacific Ocean (shown in Figures 1 and 2). We first do statistical
91 analysis of the tsunami waveforms to estimate the propagating speed of the Lamb wave and to understand
92 the tsunami wave characteristics in the Pacific Ocean through demonstrating the tsunami wave properties,
93 i.e., arrival times, wave heights and durations. We then conduct wavelet analysis for representative DART
94 buoys and tide gauges respectively to explore tsunamigenic mechanisms of the event and to better
95 understand its hydrodynamic processes in the Pacific Ocean. Aided by wavelet analysis of corresponding
96 barometers near the selected DART buoys and comparison with tsunami records of the 2011 Tohoku
97 tsunami, we are able to piece together all the analysis and demonstrate that the 2022 HTHH tsunami was
98 generated by air-sea coupling with a wide range of atmospheric waves with different propagating
99 velocities and period bands, and seafloor crustal deformation associated with the volcanic eruption. We
100 demonstrate as well that the tsunami was amplified at the far-field Pacific coastlines where the local
101 bathymetric effects play a dominant role in tsunami scale.

102 **2. Data and Methods**

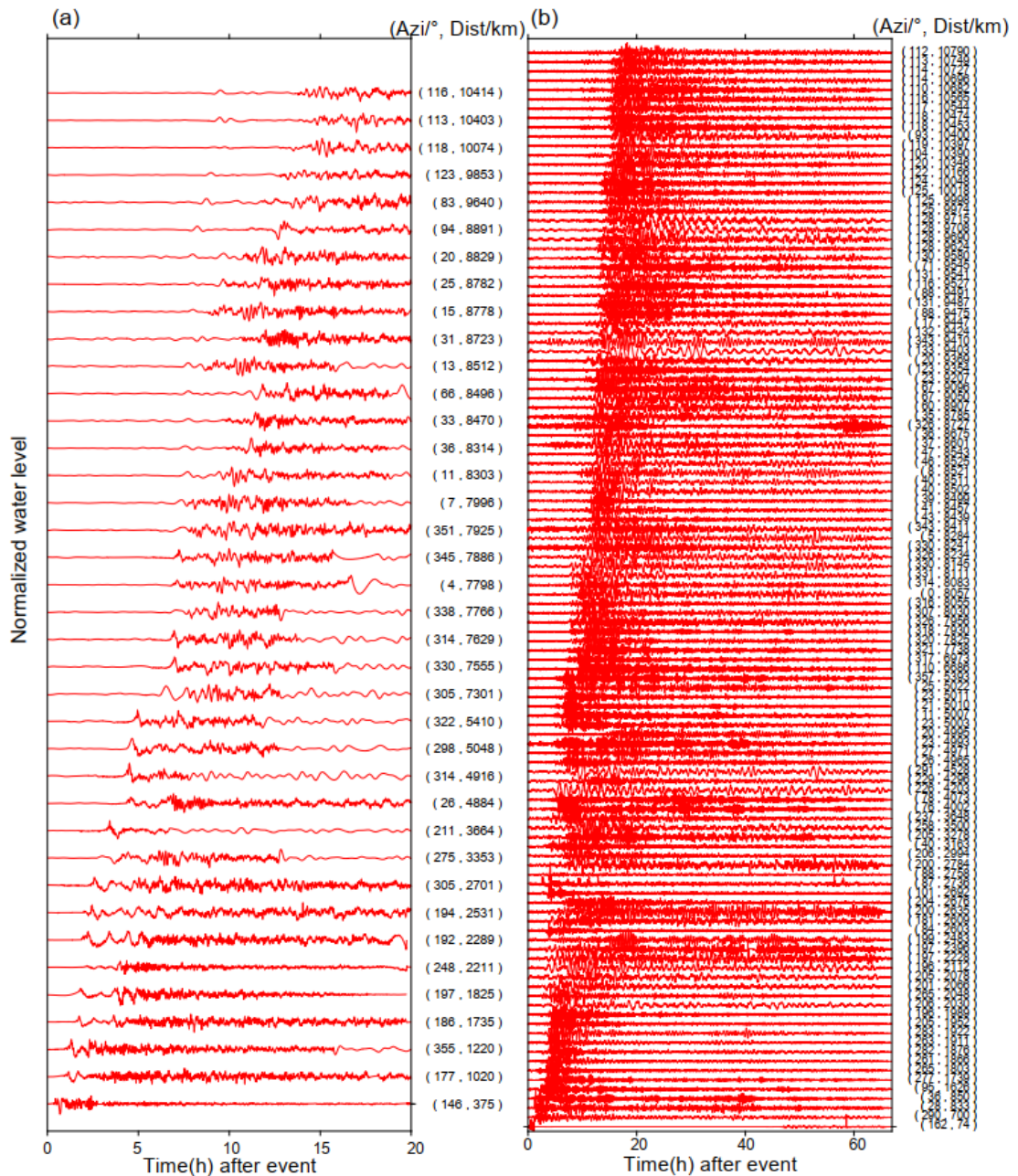
103 **2.1 Data**

104 We collected high-quality sea level records across the Pacific Ocean at 38 DART buoys (in which 31
105 stations from <https://nctr.pmel.noaa.gov/Dart/>, 7 stations from <https://tilde.geonet.org.nz/dashboard/>) and

106 116 tide gages from IOC (The Intergovernmental Oceanographic Commission, [http://www.ioc-](http://www.ioc-sealevelmonitoring.org)
107 [sealevelmonitoring.org](http://www.ioc-sealevelmonitoring.org)) (Figure 1). The epicentral distances of tide gauges and DART buoys range
108 between 74–10790 km and 375–10414 km, respectively. The sampling rates of DART buoys are
109 changing over time. Passing of tsunami event generally can trigger the DART system to enter its high
110 frequency sampling mode (15 seconds or 1 min) from normal frequency mode (15 min)
111 (www.ndbc.noaa.gov/dart). In contrast, sampling rates of normal tide gauges at coasts are uniform with
112 sampling interval of 1 min. The sampling interval of both DART and tide gauges is preprocessed to 15
113 seconds. Firstly, we eliminate abnormal spikes and fill gaps by linear interpolation. Secondly, we applied
114 a fourth-order Butterworth-Highpass filter with a cut-off frequency of 3.5×10^{-5} Hz (~ 8 hours) to remove
115 the tidal components (Figure 2) (Heidarzadeh and Satake, 2013). After the two steps, quality control step
116 is conducted to select high-quality data, in which we delete waveforms with spoiled data or massive data
117 loss due to equipment failure, or with the maximum tsunami heights of tide gauges less than 0.2 m, then
118 the selected data will be ready for further statistics and spectral analysis. We also collect and analyze the
119 atmospheric pressure disturbance data recorded by some representative barometers. The sampling rates
120 of the barometers is generally uniform with a sampling rate of 1 min except for some stations in New
121 Zealand with interval of 10 min. Considering the sample rate, we employ a fourth-order Butterworth-
122 Bandpass filter with period ranging between 2–150 min for wavelet analysis of the barometers with 1
123 min sample rate, while we apply the fourth-order Butterworth-Bandpass filter with range of 30–150 min
124 to long-period waveform display based on two reasons. (1) The barometer data we use for the analysis
125 include some in New Zealand with 10 min sample rate; (2) Filtering out the short-period waves helps
126 highlight long-period tsunami wave components.

127 The tsunami waveforms recorded by DART buoys which are installed offshore in the deep water are
128 expected to contain certain characteristics of the tsunami source (Wang et al., 2020, 2021). The
129 waveforms recorded by tide gauge distributed along coastlines are significantly influenced by local
130 bathymetry/topography which are used for investigating bathymetric effect on tsunami behaviors
131 (Rabinovich et al., 2017, 2006; Rabinovich, 2009). Therefore, we use the DART data for source-related
132 analysis and choose some tide gauge data to investigate the tsunami behaviors at the Pacific coastlines.

133



134

135 **Figure 2. Detided tsunami waveforms at (a) DART buoys and (b) tide gauges. Waveforms in both**
 136 **subplots are shown in ascending distance. Azi stands for azimuth. The data are normalized with**
 137 **respect to the largest amplitude of each tide gauge.**

138 **2.2 Tsunami Modelling**

139 We use a numerical tsunami modelling package JAGURS (Baba et al. 2015) to simulate the tsunami
 140 propagation of the 2022 HTHH event and obtain the theoretical tsunami arrival time based on the shallow
 141 water wave speed (white contours in Figure 1). The code solves linear Boussinesq-type equations in a
 142 spherical coordinate system using a finite difference approximation with the leapfrog method. We specify
 143 a unit Gaussian-shaped vertical sea surface displacement at the volcanic base as the source of

144 conventional tsunami. For a unite source i with center at longitude φ_i and latitude θ_i , the
145 displacement distribution $Zi(\varphi, \theta)$ can be expressed as:

$$146 \quad Zi(\varphi, \theta) = \exp\left[-\frac{(\varphi-\varphi_i)^2+(\theta-\theta_i)^2}{2\sigma}\right] \quad (1)$$

147 Where we set characteristic length σ as 5 km (NASA, 2022). The bathymetric data is resampled from the
148 GEBCO 2019 with 15 arc-sec resolution (The General Bathymetric Chart of the Oceans, downloaded
149 from <https://www.gebco.net>).

150 **2.3 Spectral Analysis of Tsunami Waves**

151 To investigate the temporal changes of the dominant wave periods, we conduct continuous wavelet
152 transformation (frequency-time) analyses for some representative DART buoys, tide gauges and
153 barometers, in which wavelet Morlet mother function is implemented (Kristeková et al., 2006). The first
154 32-hour time series of DART buoys and barometers after the eruption (at 04:14:45 on 15 January 2022)
155 are used for source-related wavelet analysis. The first 48-hour time series of tide gauges after the eruption
156 are employed for hydrodynamics-related wavelet analysis at coastlines. We adopt the Averaged-Root-
157 Mean-Square (ARMS) method as a measure of absolute average tsunami amplitude with a moving time
158 window of 20 min to calculate the tsunami duration (Heidarzadeh and Satake, 2014). We define the time
159 durations as the time period where ARMS levels of tsunami waves are above those prior to the tsunami
160 arrivals.

161 **3. Results**

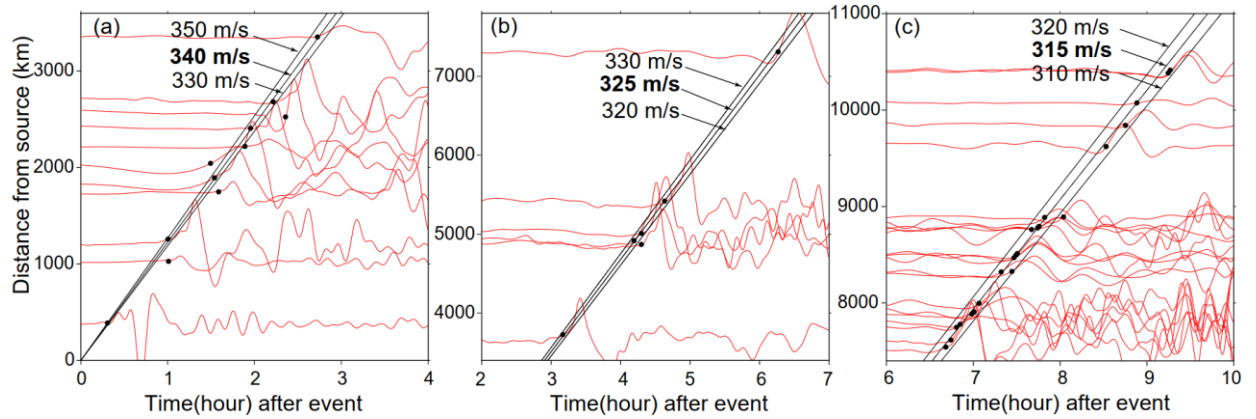
162 **3.1 The decreasing propagation velocities of the Lamb Wave**

163 Although many types of atmospheric waves were generated by the 2022 HTHH eruption, the most
164 prominent signature was the Lamb waves which were globally observed by ground-based and spaceborne
165 geophysical instrumentations (Kulichkov et al., 2022; Liu et al., 2022; Lin et al., 2022; Matoza et al.,
166 2022; Themens et al., 2022; Adam, 2022; Kubota et al., 2022). Interestingly, we notice that a wide range
167 of the velocities from 280 m/s to 340 m/s were proposed through observations and Lamb wave modelling
168 (e.g., Kubota et al., 2022; Lin et al., 2022; Matoza et al., 2022; Themens et al., 2022). The travelling
169 velocity of Lamb waves in real atmosphere is affected by temperature distributions, winds and dissipation
170 (Otsuka, 2022). To investigate whether the propagation speeds of the lamb wave change in space and

171 time, we analyze the waveforms recorded by the DART buoys in the Pacific Ocean. DART buoy with
172 pressure sensor deployed at the ocean's bottom records the sea level change that is transferred from
173 pressure records in Pascals, instead of direct water height. For the 2022 HTHH tsunami event, the
174 pressure fluctuation at DART buoy is a superposition of the pressure changes caused by tsunami and the
175 Lamb wave (Kubota et al., 2022). The Pacific DART buoys recorded the most discernible air-sea
176 coupling pulse in deep ocean with Lamb waves that arrived earlier than the theoretical tsunamis (Figure
177 1). The tsunami waveforms recorded by tide gauges did not clearly detect the tsunami signals associated
178 with the Lamb waves, therefore are not sufficient for further analysis (Figure 2). Thus, we estimate the
179 speed of Lamb waves using the waveforms recorded by the Pacific DART buoys. The Lamb wave
180 arrivals are limited within arrival time range from possible velocities of 280–340 m/s. The time points at
181 which the tsunami amplitudes first exceed 1 e-4 m above sea level are defined as Lamb wave arrivals.
182 Using different velocity values as constraints, we illustrate the velocities of Lamb wave were generally
183 uniform, but slightly decrease with the increase of propagation distance (Figure 3). The Lamb waves
184 initially propagated radially at speed of ~340 m/s before slowing to ~325 m/s after reaching ~3400 km,
185 and further decreasing to ~315 m/s at 7400 km. In an isothermal troposphere assumption, the phase
186 velocity of the Lamb wave (C_L) can be estimated with the following equation (Gossard and Hooke,
187 1975b):

$$188 \quad C_L = \sqrt{\frac{\gamma \cdot R \cdot T}{M}} \quad (2)$$

189 Where $\gamma = 1.4$ (air specific heat ratio corresponding to atmospheric temperature), $R = 8314.36 \text{ J kmol}^{-1}$
190 K^{-1} (the universal gas constant), $M = 28.966 \text{ kg kmol}^{-1}$ (molecular mass for dry air) are constant for the
191 air, T is the absolute temperature in kelvin. Thus, Lamb wave velocity is mainly affected by the air
192 temperature, meaning the travelling velocity of lamb waves might decrease when propagating from
193 regions with high temperature towards those with low temperatures, e.g., the north pole. By assuming a
194 set of possible temperatures in January (Table 1), we calculated the velocities C_L could range between
195 312–343 m/s when temperatures vary between -30–20 °C. Therefore, the decreased velocity of the Lamb
196 waves could be a consequence of cooling of the air temperature.



197
 198 **Figure 3. Identifying the Lamb wave-induced tsunami velocities using different constant velocities**
 199 **as constraints. Black dots mark the arrival times of the Lamb waves. Black lines represent**
 200 **velocities.**

201 **Table 1. Estimated Lamb wave velocities in an isothermal troposphere assumption**

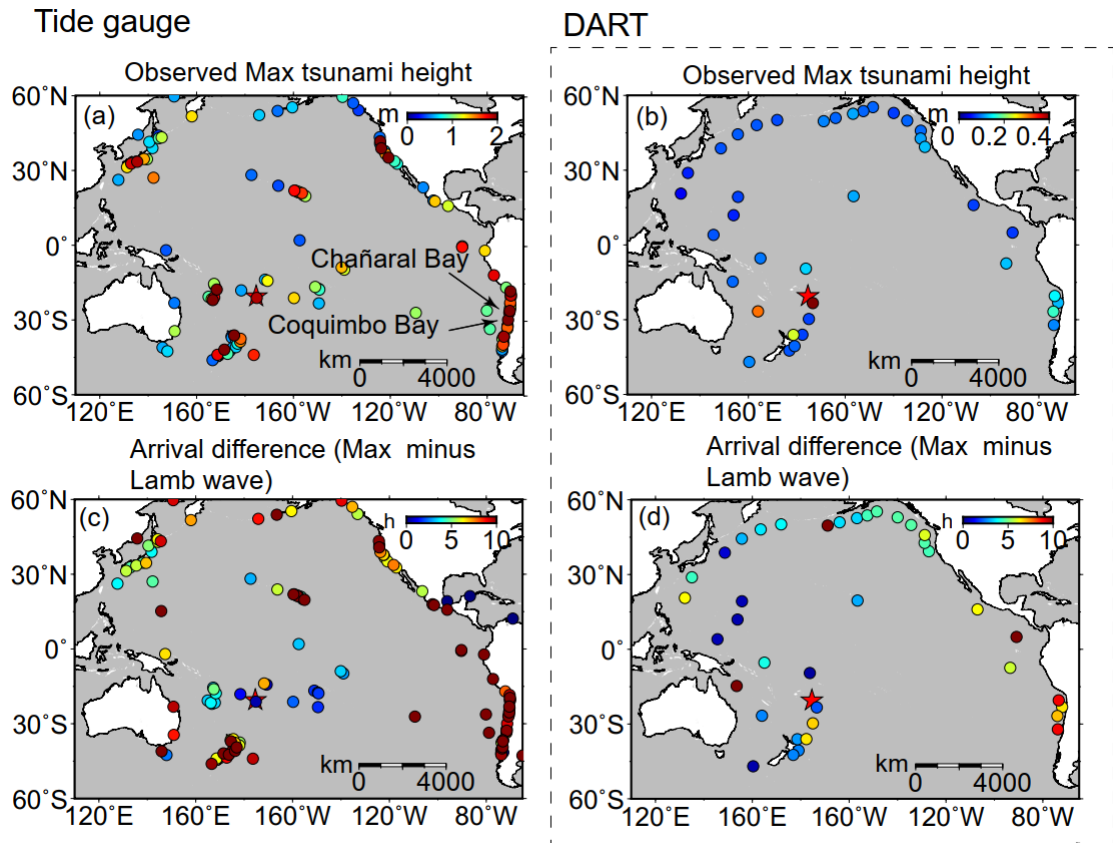
| Celsius temperature (°C) | thermodynamic temperature (K) | C_L (m/s) |
|--------------------------|-------------------------------|-------------|
| 20 | 293.15 | 343.14 |
| 10 | 283.15 | 337.23 |
| 0 | 273.15 | 331.21 |
| -10 | 263.15 | 325.19 |
| -20 | 253.15 | 318.86 |
| -30 | 243.15 | 312.49 |

202 **3.2 Tsunami features observed by DART buoys and Tide gauges**

203 The statistics of tsunami heights and arrival times recorded at 38 DART buoys and 116 tide gauges across
 204 the Pacific Ocean are used to interpret the tsunami characteristics. The comparison of the statistical
 205 characters between DART and tide gauge observations yields some useful information of the
 206 hydrodynamic process of tsunami propagation and help identify tsunami wave components with different
 207 traveling velocities.

208 The average value of the maximum tsunami wave height (trough-to-crest) for the 116 tide gauge stations
 209 is ~ 1.2 m. Figure 4a shows tide gauges with large tsunami heights exceeding 2 m are mainly distributed
 210 in coastlines with complex geometries (Figure S1a), such as gauges at New Zealand, Japan, and north
 211 and south America. For example, the largest tsunami height among tide gauges is 3.6 m at a bay-shaped
 212 coastal area Chañaral in Chile (Figure S1b). In sharp contrast to tide gauges, the maximum tsunami
 213 heights of most Pacific DART buoys are less than 0.2 m. The largest tsunami height in the DART buoys
 214 is only ~ 0.4 m recorded at the nearest one, 375 km from the volcano (Figure 4b). The comparison between
 215 DART buoys and tide gauges indicate that the direct contribution of air-sea coupling to the tsunami

216 heights is probably in the level of tens of centimeters (Kubota et al., 2022). The meter-scale tsunami
217 heights at the coastlines suggest the bathymetric effect could play a major role during tsunami
218 propagation. In respect to the arrival of maximum tsunami waves, the time lags between Lamb waves
219 and the maximum heights of tide gauges mainly range between ~0–10 h (Figure 4c). The delayed times
220 of ~10 h are observed in New Zealand, Hawaii, and west coast of America (Figure 4c), suggesting the
221 interaction between tsunami waves and local topography/bathymetry delays the arrival of the maximum
222 waves (e.g., Hu et al., 2022). For example, the delayed maximum tsunami height can be attributed to the
223 edge waves (Satake et al., 2020) and resonance effect (Wang et al., 2021) from tsunami interplays with
224 bays/harbors, islands, and continental shelves of various sizes. The significant regional dependence of
225 the coastal tsunami heights and the time lags of the maximum tsunami waves can be attributed to the
226 complexity of local bathymetry, such as continental shelves with different slopes, and harbor/bay with
227 different shapes and sizes (Satake et al., 2020). On the other hand, for tsunami events with earthquake
228 origins (e.g. Heidarzadeh and Satake, 2013), the first waves recorded by DART buoys are normally
229 observed as the largest wave since DART buoys are located in the deep sea and less influenced by
230 bathymetric variation. In the case of Tonga tsunami event, we observe the inconsistency between the
231 arrivals of the Lamb wave-induced tsunami waves and the maximum tsunami heights (Figure 4d). The
232 time lags of the maximum waves of DART buoys present a coarsely increasing tendency with the
233 increasing distance from the volcano, which indicates the contribution of other tsunami generation
234 mechanism propagating with a uniform but lower speed than Lamb wave.



235
 236 **Figure 4. The spatiotemporal signatures of the 2022 HTHH tsunami across the Pacific Ocean. (a)**
 237 **Observed the maximum tsunami height (trough-to-crest height) of tide gauges. (c) Arrival**
 238 **differences between the maximum tsunami height of tide gauges and Lamb waves. (b) and (d) are**
 239 **the same as (a) and (c) but for DART buoys.**

240 3.3 Tsunami components identified from wavelet analysis

241 The statistical analysis of tsunami waveforms at tide gauges and DART buoys suggest the tsunami waves
 242 likely contain several components with different source origins. To further identify these tsunami
 243 components, we conduct wavelet analysis for tsunami waveforms recorded by representative DART
 244 buoys and air pressure waveforms recorded by selected barometers. We demonstrate the analysis result
 245 through the frequency-time (f-t) plot of wavelet which shows how energy and period vary at frequency
 246 and time bands (Figure 5 and Figure 6). Tsunami components have clear signatures in all f-t plots as the
 247 energy levels are quite large when they arrive. Figure 5 shows the wavelet analysis of six DART buoys
 248 located in the vicinity of the eruption site (<3664 km). Figure 6 show the wavelet analysis of ten DART
 249 buoys located in the Pacific rim which are far away from the source location. We observe three interesting
 250 phenomena: 1) most of the tsunami wave energy is concentrated in four major period bands, i.e., 3–5
 251 min, ~10–30 min, ~30–40 min, and ~80–100 min; 2) The significant tsunami component with period

252 band of 3-5 mins are recorded by stations between the eruption site and the north tip of the New Zealand;
 253 3) There exists one exceptional tsunami component with longer wave period of ~80–100 min mainly
 254 recorded in the Tonga, the New Zealand and Hawaii, which travels even faster than the Lamb waves.
 255 To further explore the source mechanism of these tsunami components, we take advantage of the
 256 published information related to different propagating velocities of atmospheric gravity waves (Kubota
 257 et al., 2022) and add four kinds of propagating velocities as criteria to differentiate the tsunami arrivals
 258 from different sources (Figure 5 and Figure 6). The first reference speed is 1000 m/s related to the
 259 radically propagating atmospheric shock waves near the source region (Matoza et al., 2022; Themens et
 260 al., 2022). The second one is the velocities of Lamb wave ranging between 315–340 m/s derived from
 261 the aforementioned analysis in section 3.1 (Figure 3). The third one is 200 m/s corresponding to the lower
 262 limit of atmospheric gravity wave modes other than Lamb waves which were also excited by the volcanic
 263 eruptions (Kubota et al., 2022). The last is the arrival time of conventional tsunami given by tsunami
 264 modelling (Figure 1). The theoretical velocity of conventional tsunami is significantly nonuniform
 265 spatially as compared with those of the atmospheric waves. The conventional tsunami propagation speed
 266 is determined by the water depth along the propagation route. The velocity of non-dispersion shallow-
 267 water waves (C_H) in the ocean is given by:

$$268 \quad C_H = \sqrt{g \cdot H} \quad (3)$$

269 Where g is gravity acceleration (9.81m/s^2), H is the water depth. The propagation velocities of tsunami
 270 are ~296–328 m/s in the deepest trenches on earth (i.e., ~11 km in Mariana Trench and ~9 km in Tonga
 271 Trench). The velocities decrease quickly to only ~44 m/s at ~200 m depth along the edge of continental
 272 shelf. With the average depth of ~4–5 km, the average velocities in the Pacific Ocean range between
 273 ~200–224 m/s. Thus, theoretical tsunami velocities present significant slowness and variability. We
 274 delineate the arrival times of the four reference speeds in Figures 5 and 6.

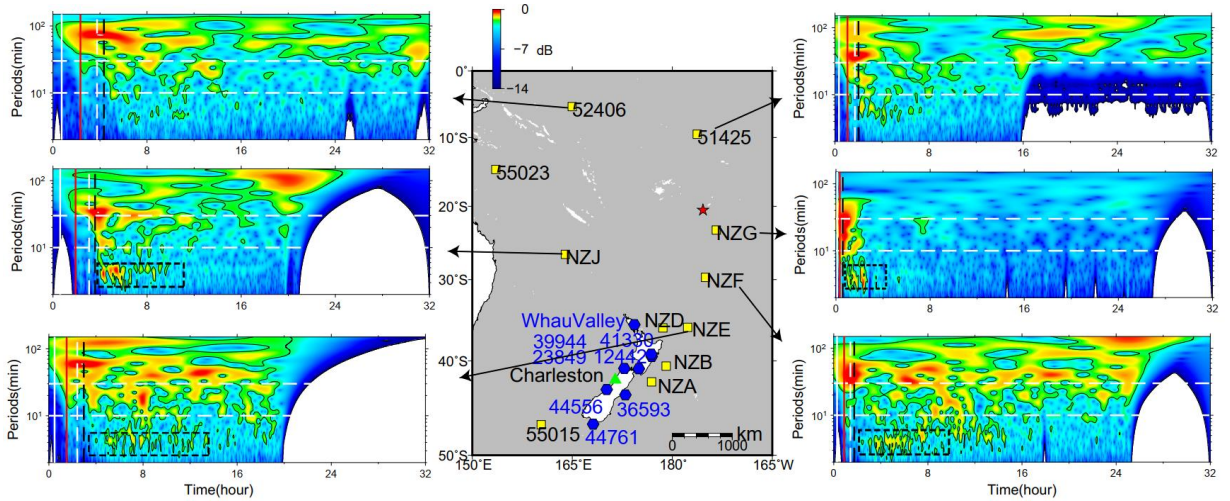
275 One particularly remarkable phenomenon is that the wave component with period of ~80–100 min
 276 propagated at a very fast speed of ~1000 m/s in the vicinity of the HTHH site, i.e., New Zealand and
 277 Hawaii (e.g., stations 52406, NZJ, NZE, 51425 in Figure 5, and 51407 in Fig. 6). We infer that the
 278 tsunami component within ~80–100 min period band was likely produced by the atmospheric shock
 279 waves during the initial stage of the volcanic eruption and spatially only cover the near-source region.
 280 To verify this observation, we select 16 representative barometers located in the near-source region and

281 far-field area for wavelet analysis (see the locations in Figure 5 and Figure 6). Figure 7 shows the
282 waveforms of atmospheric pressure at selected locations and Figure 8 provides the frequency-time (f-t)
283 plot of wavelet analysis of some representative barometers. Interestingly, we are able to discern the air
284 pressure pulses prior to Lamb waves at barometers in New Zealand (the two columns on the left in Figure
285 7), although such signals are not detectable in waveforms recorded by barometers far from the source
286 (the two columns on the right in Figure 7). The spatial distribution of such unusual pressure changes
287 suggest that the fast travelling shock waves were only limited in the near-source region, as reflected in
288 the travelling ionospheric disturbances (Matoza et al., 2022; Themens et al., 2022). Additionally, we also
289 see that the long period signals of ~80–100 min appear in DART buoys far away from the eruption site.
290 Such signals may be related with the long-period gravity waves (Matoza et al., 2022).

291 The tsunami components at period band of ~30–40 min can be readily associated with Lamb waves
292 because the arrival times of the tsunami waves and Lamb waves have excellent match, as shown in the
293 tsunami data recorded by DART buoys (e.g., NZJ and 51425 in Figure 5; 51407, 32401 and 32413 in
294 Figure 6) and pressure data by barometers (Figure 8).

295 For the tsunami components with the period band of ~10–30 min, although the arrivals of ~10–30 min
296 tsunami components cover some theoretical tsunami arrival times, they do not consistently match. The
297 tsunami components occurring within the time period between Lamb waves and the lower gravity waves’
298 velocities has a good agreement with the velocity range of several atmospheric gravity wave modes
299 (Matoza et al., 2022; Themens et al., 2022; Kubota et al., 2022). Similarly, the air pressure data also show
300 energy peaks at ~10–30 min period band, which is consistent with the tsunami data (Figure 8). Such
301 consistency further verifies the contribution of atmospheric gravity waves to the volcanic tsunami.

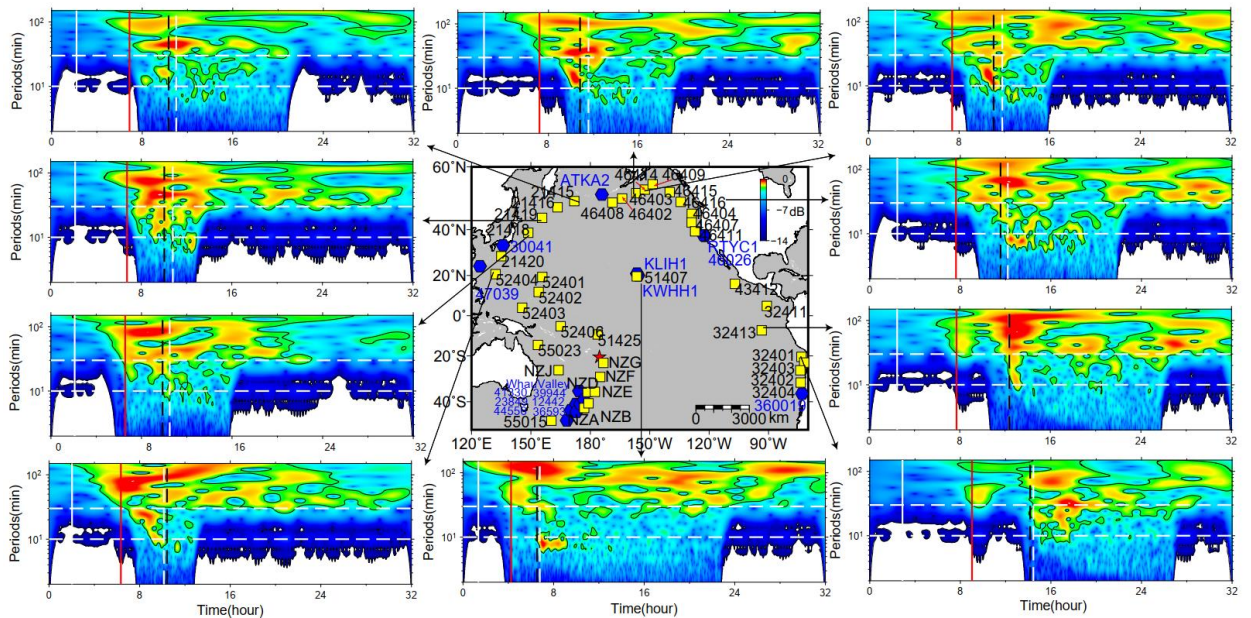
302 The tsunami components with the shortest period of ~3–5 min (stations NZE, NZF, NZG and NZJ;
303 marked with black dashed squares in Figure 5) are only observed at DART records near the eruption
304 location. Meanwhile, the arrival times of these components agree well with the modelled arrivals of
305 conventional tsunami. Thus, we believe the observed shortest period band should originate from the
306 seafloor crustal deformation. We further infer that this component could be generated by the partial
307 underwater caldera collapse and/or subaerial/submarine landslide failures associated with 2022 HTHH
308 volcanic eruption.



309

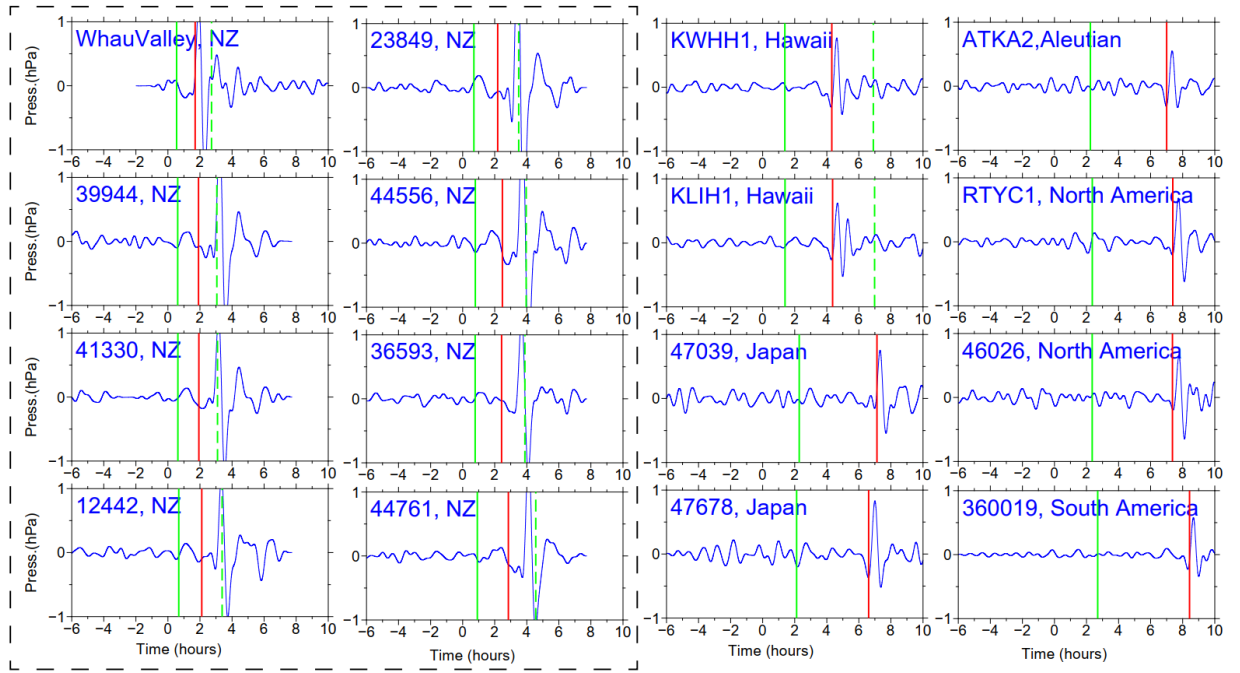
310 **Figure 5. Wavelet analysis of representative DART buoys in the vicinity of the HTHH volcano. In**
 311 **each sub-plot, the solid vertical white lines mark the arrival time with travelling velocity of 1000**
 312 **m/s. The solid vertical red lines mark the arrivals of Lamb waves. The dashed vertical white lines**
 313 **mark lower limit of AGWs' velocity of 200 m/s (Kubota et al., 2022). The dashed vertical black**
 314 **lines represent the theoretical tsunami arrivals. The dashed horizontal white lines mark two**
 315 **reference wave periods of 10 min and 30 min. The blue hexagons represent the locations of**
 316 **barometers. Green triangle makes the location of the tide gauges at Charleston. Decibel (dB) is**
 317 **calculated from: $dB = 10 \log(A/A_0)$, where A is wavelet power, A_0 is a reference wavelet power of**
 318 **the maximum one (Thomson and Emery, 2014).**

319



320

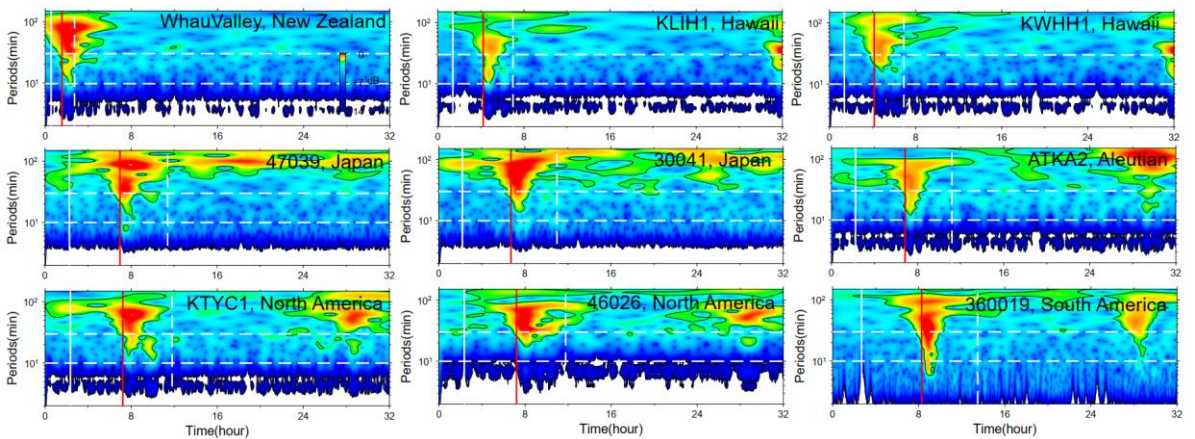
321 **Figure 6. Wavelet analysis of representative DART buoys far away from the HTHH volcano. In**
 322 **each sub-plot, the solid vertical white lines mark the arrival time with travelling velocity of 1000**
 323 **m/s. The solid vertical red lines mark the arrivals of Lamb waves. The dashed vertical white lines**
 324 **mark lower limit of AGWs' velocity of 200 m/s. The dashed vertical black lines represent the**
 325 **theoretical tsunami arrivals. The dashed horizontal white lines mark two reference wave periods**
 326 **of 10 min and 30 min. The blue hexagons represent the locations of barometers.**



327

328 **Figure 7. Shockwave-related atmospheric pressure waveforms of selected barometers in the Pacific**
 329 **Ocean. All traces have been filtered between 30 min and 150 min. In each sub-plot, the solid vertical**
 330 **green lines mark the arrival time with travelling velocity of 1000 m/s. The solid vertical red lines**
 331 **mark the arrivals of Lamb waves. The dashed vertical green lines mark lower limit of AGWs'**
 332 **velocity of 200 m/s.**

333



334

335 **Figure 8. Wavelet analysis of some representative barometers. In each sub-plot, the solid vertical**
 336 **white lines mark the arrival time with travelling velocity of 1000 m/s. The solid vertical red lines**
 337 **mark the arrivals of Lamb waves. The dashed vertical white lines mark lower limit of AGWs'**
 338 **velocity 200 m/s. The dashed horizontal white lines mark three reference periods of 10 min and 30**
 339 **min.**

340 **4. Discussion**

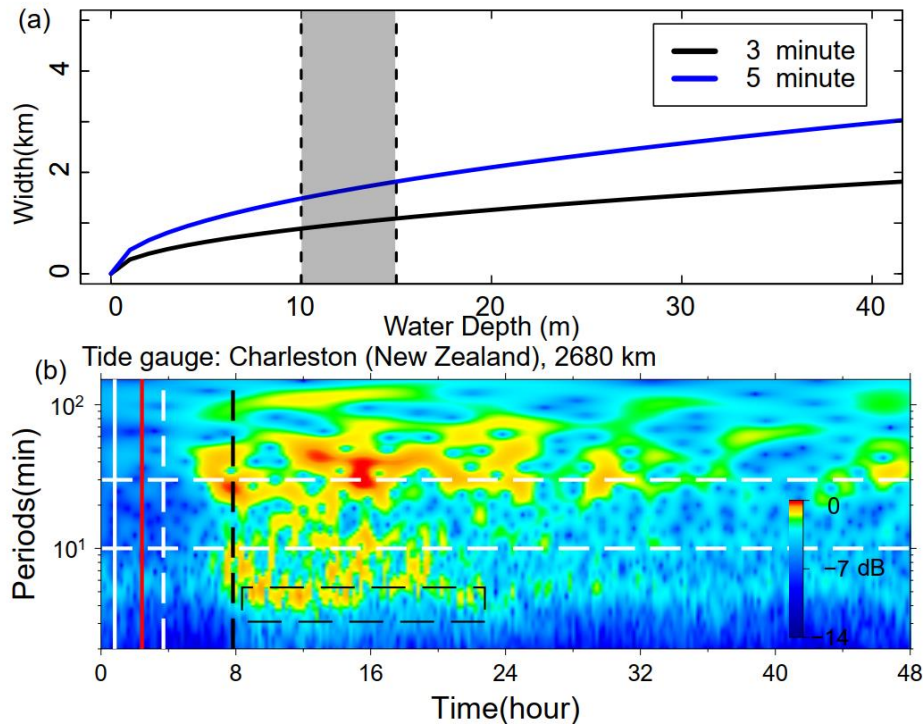
341 **4.1 Tsunami from Caldera Collapse and Its Long-distance Traveling Capability**

342 The tsunami wave energy distributed in different period bands is identified with reference arrival times.
343 The tsunami component with 3–5 min period is most likely generated by seafloor crustal deformation in
344 the volcanic site, but specific mechanism is not determined. A variety of possible scenarios associated
345 with the eruption could be responsible for the near-field tsunami waves, such as volcanic earthquakes,
346 pyroclastic flows entering the sea, underwater caldera flank collapse, and subaerial/submarine failures
347 (Self and Rampino, 1981; Pelinovsky et al., 2005). To further investigate the source mechanism, we
348 apply a simplified model (Rabinovich, 1997) to estimate the probable dimension of tsunami source:

349
$$L = \frac{T\sqrt{gH}}{2} \quad (4)$$

350 Where L is the typical dimension (length or width) of the tsunami source, H is average water depth in the
351 source area, g is the gravity acceleration, and T is primary tsunami period. By comparing with the post-
352 2015 morphology of the HTHH caldera which was obtained through drone photogrammetry and
353 multibeam sounder surveys, Stern et al. (2022) estimate that much of the newly-formed Hunga Tonga
354 Island and the 2014/2015 cone were destroyed by the 2022 eruption, and the vertical deformation of
355 Hunga Ha’apai Island is ~10–15 m (Stern et al., 2022). With no more quantitative constraint of the
356 seafloor deformation, we tentatively assume H as 10–15 m, then the possible dimension of seafloor
357 crustal deformation responsible for the small-scale tsunami could be in the scale of 0.8–1.8 km (Figure
358 9a). The estimated size is very likely from partial caldera collapse that usually has limited scale in
359 volcanic site (Ramalho et al., 2015; Omira et al., 2022). If it is the case, the partial flank collapse could
360 be located between Hunga Tonga and Hunga Ha’apai Islands.

361



362
 363 **Figure 9. Mechanism of tsunami component with 3–5min period. (a) The source dimension**
 364 **estimated by equation 4. (b) Wavelet analysis of tide gauge at Charleston, New Zealand, 2680 km**
 365 **away from the eruption site. The solid vertical white line marks the arrival time with travelling**
 366 **velocity of 1000 m/s. The solid vertical red line marks the arrival of Lamb wave. The dashed**
 367 **vertical white line marks lower limit of AGWs' velocity 200 m/s. The dashed vertical black line**
 368 **marks the theoretical tsunami arrivals.**

369 An interesting phenomenon is that the tsunami component with 3–5 min period can still be observed in
 370 a bay-shaped coastal area at Charleston in New Zealand (see the location in Figure 5) which is 2680 km
 371 away from the eruption site and maintains a high energy level lasting up to 14 h (Figure 9b). The long-
 372 traveling capability could be associated with the ~ 10000 m deep water depth of the Tonga Trench that
 373 keeps the source signals from substantial attenuation. In deep open ocean, the wavelength of a tsunami
 374 can reach two hundred kilometers, but the height of the tsunami may be only a few centimeters. Tsunami
 375 waves in the deep ocean can travel thousands of kilometers at high speeds, meanwhile losing very little
 376 energy in the process. The long oscillation can be attributed to the multiple reflections of the incoming
 377 waves trapped in the shallow-water bay at Charleston.

378 Generally, devastating tsunamis with long-distance travelling capability are mostly generated by
 379 megathrust earthquakes (Titov et al., 2005). Caldera collapses or submarine landslides with limited scale
 380 normally only generate local tsunamis, e.g., the 1998 PNG (Papua New Guinea) tsunami event (Kawata
 381 et al., 1999) and the 1930 Cabo Girão tsunami event (Ramalho et al., 2015). Therefore, it's exceptional

382 that the tsunami component from scale-limited failure could travel at-least 2680 km away from the
383 eruption site. It demonstrates that tsunamis from small-scale tsunamigenic source have the capability to
384 travel long distance and cause long oscillation at favored condition, e.g., deep trench, ocean ridge and
385 bay-shaped coasts.

386 **4.2 The Possible Mechanisms of Long Tsunami Oscillation**

387 An important tsunami behavior of the 2022 HTHH tsunami is the long-lasting oscillation ~ 3 days in the
388 Pacific Ocean (Figure 10a), which is comparable to that of the 2011 Tohoku tsunami, ~ 4 days
389 (Heidarzadeh and Satake, 2013). We demonstrate the duration time of the tsunami oscillation through
390 ARMS (Averaged-Root-Mean-Square) approach that is a measure of absolute average tsunami amplitude
391 in a time period. The long-lasting tsunami energy can be observed at many regions, such as the coasts of
392 New Zealand, Japan, Aleutian, Chile, Hawaii, and west coasts of America. Several mechanisms could
393 account for the long-lasting tsunami, including (1) Lamb waves circling the Earth multiple times
394 (Amores et al., 2022; Matoza et al., 2022), (2) resonance effect between ocean waves and atmospheric
395 waves (Kubota et al., 2022), and (3) bathymetric effect. We discuss the contribution of each mechanism
396 in the following section.

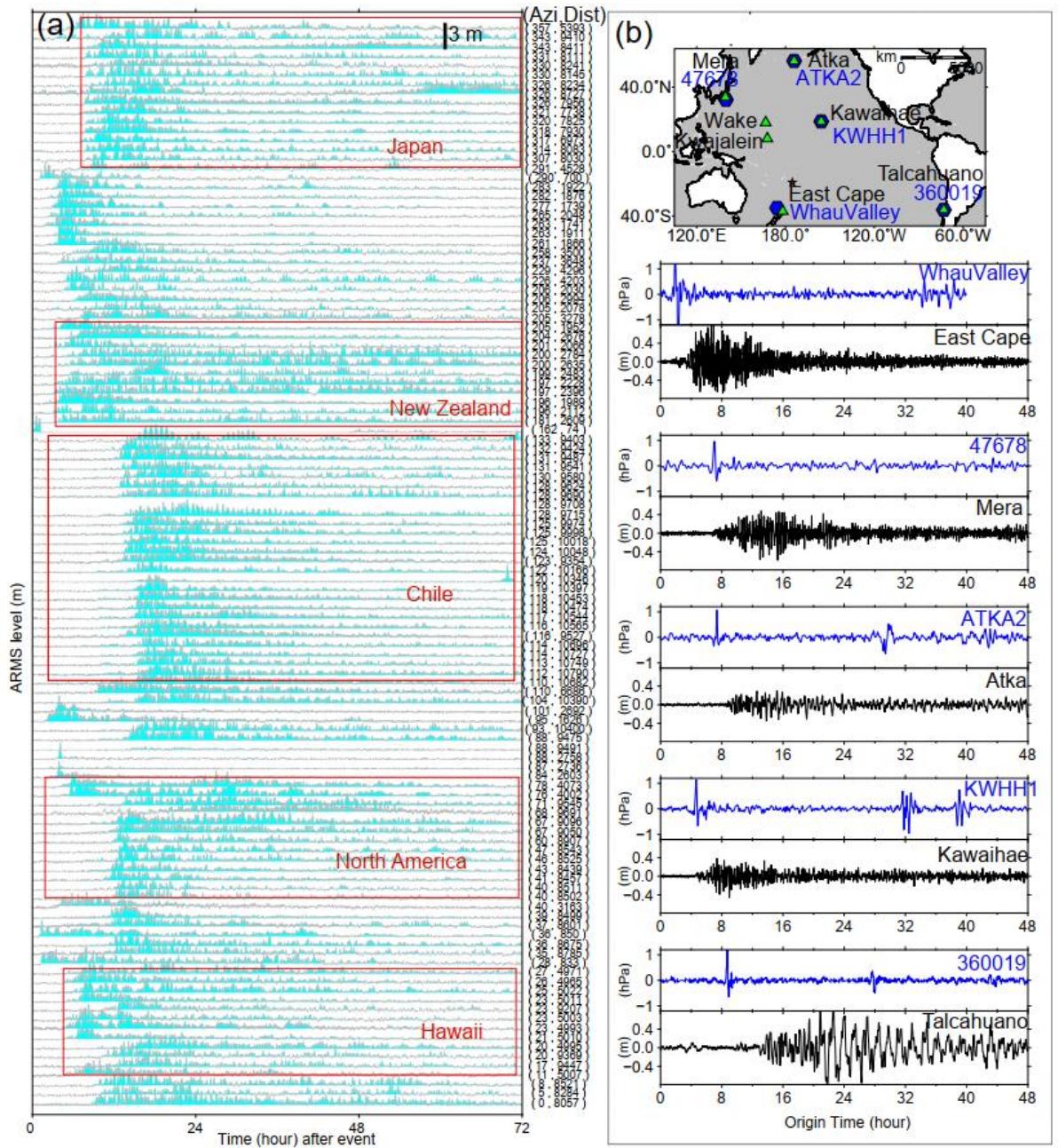
397 To investigate the contribution of Lamb wave to the long-lasting tsunami, we compare the air pressure
398 disturbances recorded by selected barometers together with the tsunami waveforms of nearby tide gauges
399 (Figure 10b). While the barometers present discernible wave pulses at each Lamb wave's arrival, only
400 the first Lamb wave triggered clear tsunami signal and no detectable tsunami signatures correspond to
401 the following passage, suggesting the Lamb waves do not directly contribute to the long oscillation.

402 The resonance effects between ocean waves and atmospheric waves could contribute to the long
403 oscillation on coastlines. Besides the Lamb wave, Watanabe et al., 2022 detected internal Pekeris wave
404 which propagate with a slower horizontal phase speed of ~ 245 m/s and gravity waves with even slower
405 propagation speed by analyzing radiance observations taken from the Himawari-8 geostationary satellite.
406 Atmospheric waves with such speeds are more likely to resonant with the conventional tsunami waves
407 and provide continuous energy supply (Kubota et al., 2022).

408 To examine the role of local bathymetry in the long-lasting tsunami, we choose a well-studied and well-
409 recorded event: the 2011 Mw 9.0 Tohoku tsunami as a reference event and compare the tsunami records
410 of these two events at the same coastal stations. Although the two tsunami events were generated by

411 completely different mechanisms, i.e., large-scale seafloor deformation for the Mw 9.0 megathrust
412 earthquake (Mori et al., 2011) and fast-moving atmospheric waves for the Mw 5.8 volcanic eruption
413 (Matoza et al., 2022), they both produced widespread transoceanic tsunamis which were well recorded
414 in the Pacific DART buoys and tide gauges. In the near-field, the 2011 Tohoku earthquake produced
415 runup up to 40 m at Miyako in the Iwate Prefecture in Japan's Tohoku region (Mori et al., 2011). The
416 epicenter is approximately 70 km east coast of the Oshika Peninsula of Tohoku region. However, the
417 2022 HTHH tsunami produced only ~13 m runup in the near field from eyewitness accounts in
418 Kanokupolu, 60 km from the volcano (Lynett et al., 2022). However, in the far-field (>1000 km), we
419 observe comparable tsunami wave heights in certain coastal regions. Based on the tsunami records at 21
420 tide gauges surrounding the Pacific Ocean, Heidarzadeh & Satake (2013) calculated the average value
421 of the maximum tsunami heights (trough-to-crest) of the 2011 Tohoku tsunami is 1.6 m with the largest
422 height of 3.9 m at the Coquimbo Bay in Chile (Heidarzadeh and Satake, 2013). Coincidentally, the statistics
423 of 116 tide gauges in this study also suggest the average tsunami heights of the 2022 HTHH tsunami is
424 around the same order, ~1.2 m, among which, the largest height is 3.6 m at Chañaral Bay in Chile.
425 Interestingly, in the coastal region of South America, the locations of the largest tsunami heights of both
426 events are adjacent (Figure 4a), i.e., Coquimbo (the 2011 Tohoku) and Chañaral (The 2022 HTHH).
427 To further compare the far-field hydrodynamic processes between these two events quantitatively, we
428 conduct wavelet analysis for four representative tide gauges distributed across the Pacific Ocean, i.e.
429 coastal gauges at East Cape in New Zealand, Kwajalein Island, Wake Island, and Talcahuaho in Chile
430 (see their locations in Figures 10b). The temporal changes of tsunami energy of both events can be seen
431 in Figure 11. At each tide gauge, the tsunami energy of the 2011 HTHH (Figure 11a) and the 2022 Tohoku
432 tsunamis (Figure 11b) for the first few hours after the arrivals is nonuniform with different significant
433 peaks distributed within a wide period band of ~3–100 min. Then, the following long-lasting energy of
434 the both at each station presents similar pattern and is concentrated at identical and fairly narrower period
435 channel, i.e., ~20–30 min at East Cape in New Zealand, ~40–60 min at Kwajalein Island, ~10 min at
436 Wake Island, and ~100 min at Talcahuaho in Chile, which reflects the local bathymetric effects of natural
437 permanent oscillations (Hu et al., 2022; Satake et al., 2020). Specifically, many bathymetric effects can
438 contribute to the long-lasting tsunami, such as multiple reflections across the basins, or the continental
439 shelves, and the excited tsunami resonance in bays/harbors with variable shapes and sizes (Aranguiz et

440 al., 2019; Satake et al., 2020). For example, tide gauges around New Zealand are primarily distributed in
441 harbors/ports with major natural oscillation modes of ~20–30 min (De Lange and Healy, 1986; Lynett et
442 al., 2022). The first oscillation mode of central Chile is centered around ~100 min (Aranguiz et al., 2019).
443 Consequently, Figure 11 illustrates that the long-lasting tsunami energy of the two events is respectively
444 distributed in 20–30 min period at East Cape in New Zealand and in ~100 min period at Talcahuaho in
445 central Chile. The coupling of bathymetric oscillation mode with tsunami containing similar-period wave
446 results in the excitement of tsunami resonance, which amplifies tsunami waves and prolongs the tsunami
447 oscillation at the two stations (Heidarzadeh et al., 2019, 2021; Hu et al., 2022; Wang et al., 2022).
448 Simply put, we do not have clear evidence that atmospheric acoustic-gravity waves from the 2022 HTHH
449 eruption directly contribute to the long-lasting tsunami, but the resonance effect associated with ocean
450 waves could a possible source of increased wave energy and amplification. However, the similarity of
451 far-field hydrodynamic behaviors between the 2022 HTHH volcanic tsunami and the 2011 Tohoku
452 seismogenic tsunami well demonstrates the both went through similar hydrodynamic processes after their
453 arrivals. The consistency favors that the long-lasting tsunami of 2022 HTHH tsunami event can very
454 likely be attributed by the interplays between local bathymetry and conventional tsunami left after each
455 passage of atmospheric waves, which can well explain why the two completely distinct tsunami events
456 possess a comparable duration time.

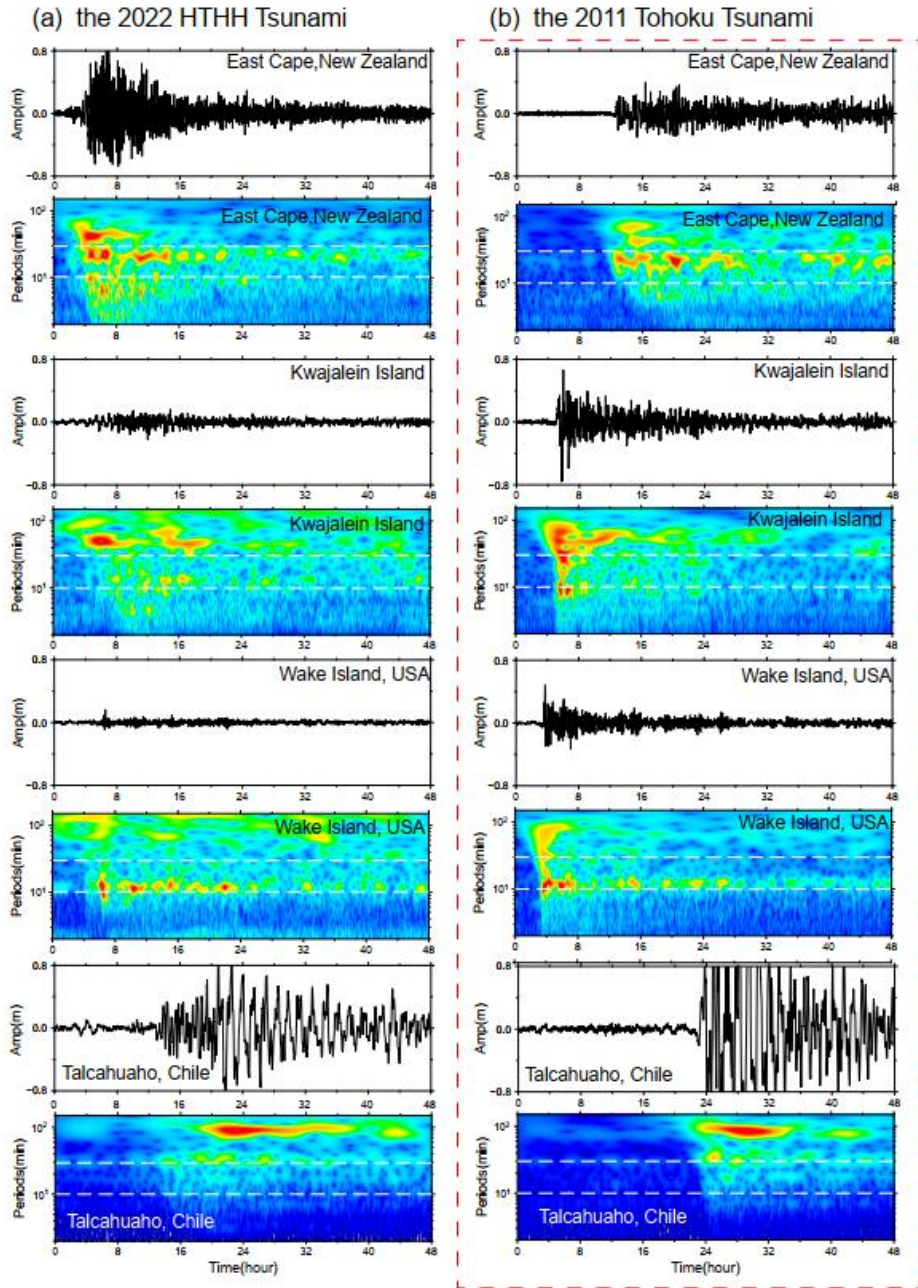


457

458 **Figure 10. Tsunami duration. (a) Tsunami durations at Pacific 116 tide gauges through ARMS level**
 459 **approach. (b) the location of barographs (blue curves) and nearby tide gauges (green curves), as**
 460 **well as their waveforms.**

461

462



463
 464 **Figure 11. Wavelet analysis of tsunami waveforms recorded by 4 tide gauges during (a) the 2022**
 465 **HTHH tsunami event, and (b) the 2011 Tohoku tsunami event. Horizontal white dashed lines**
 466 **respectively mark reference periods of 10 min and 30 min.**

467 **4.3 Challenges for Tsunami Warning**

468 The generation mechanisms and hydrodynamic characteristics of the 2022 HTHH volcanic tsunami are
 469 more complicated than pure seismogenic tsunami, which challenge the traditional tsunami warning
 470 approach.

471 The first challenge is posed by the tsunami components with propagating velocities faster than the

472 conventional tsunami. The Tonga volcanic tsunami event provides an excellent example which highlights
473 that the tsunamigenic mechanisms are not limited to tectonic activities related with the sudden seafloor
474 displacements, but also include a variety of atmospheric waves with distinct propagation velocities. The
475 tsunami components in 2022 HTHH event generated by the air-sea coupling possess a wide range of
476 velocities from 1000 m/s to 200 m/s. The Lamb waves recorded in both the 2022 HTHH event and the
477 1833 Krakatoa volcanic event traveled along the Earth's surface globally for several times (Carvajal et
478 al., 2022). The tsunami waves produced by Lamb waves, the wave components associated with resonance
479 of the air-sea coupling and their superimposition increase the difficulty of tsunami warning.
480 Another critical challenge is associated with the interplays between tsunami waves and local bathymetry.
481 The tsunami waves left by each passage of the atmospheric waves can interact with local bathymetry at
482 coastlines, such as continental shelves with different slopes, and harbor/bay with different shapes and
483 sizes. The interaction can intensify the tsunami impact and excite a variety of natural oscillation periods.
484 The 2022 HTHH tsunami with an extremely wide period range of ~2–100 min have a great potential to
485 couple with the excited natural oscillations and form extensive tsunami resonance phenomena. The
486 resonance effects result in long-lasting oscillation and delayed tsunami wave peaks. The uncertain
487 arrivals of the maximum tsunami waves pose an extra challenge to tsunami warning.

4885. **Conclusion**

489 In the study, we explore the tsunamigenic mechanisms and the hydrodynamic characteristics of the 2022
490 HTHH volcanic tsunami event. Through extensive analysis of waveforms recorded by the DART buoys,
491 tide gauges and barometers in the Pacific Ocean, we reach the main findings as follows:

492 (1) We identify four distinct tsunami wave components based on their distinct propagation velocities or
493 period bands (~80–100 min, 10–30 min, 30–40 min, and 3–5 min). The generation mechanisms of these
494 tsunami components range from air-sea coupling to seafloor crustal deformation during the volcanic
495 eruption.

496 (2) The first-arriving tsunami component with 80–100 min period was most likely from shock wave
497 spreading at a velocity of ~1000 m/s in the vicinity of the eruption. This tsunami component was not
498 clearly identified by currently available publication and it's not easy to be visually observed through time
499 series of the waveforms. The physical mechanism is yet to be understood. The second tsunami component

500 with 30–40 min period was from Lamb waves, and was the most discussed tsunami source of this event
501 so far. A thorough analysis of DART measurements indicates that the Lamb waves traveled at the speed
502 of ~340 m/s in the vicinity of the eruption and decreased to ~315 m/s when traveling away due to cooling
503 of the air temperature. The third tsunami component was from some atmospheric gravity wave modes
504 with propagation velocity faster than 200 m/s but slower than Lamb waves. The last tsunami component
505 with the shortest periods 3-5 min was probably produced by partial caldera collapse with estimated
506 dimension of ~0.8–1.8 km.

507 (3) Although the resonance effect with the atmospheric acoustic-gravity waves could be a source of
508 increased wave energy, its direct contribution to the long-lasting oscillation is not demonstrated yet.
509 However, the comparison of hydrodynamical characteristics between the 2022 HTHH tsunami event and
510 the 2011 Tohoku tsunami event well demonstrated that the interactions between the ocean waves left by
511 atmospheric waves and local bathymetry contribute to the long-lasting Pacific oscillation of the 2022
512 tsunami event.

513 (4) The extraordinary features of this rare volcanic tsunami event challenge the current tsunami warning
514 system which is mainly designed for seismogenic tsunamis. It is necessary to improve the awareness of
515 people at risks about the potential tsunami hazards associated with volcanic eruptions. New approaches
516 are expected to be developed for tsunami hazard assessments with these unusual sources: various
517 atmospheric waves radiated by volcanic eruptions besides those traditionally recognized, e.g.
518 earthquakes, landslides, caldera collapses and pyroclastic flows etc.

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525 The JAGURS tsunami simulation code is employed for tsunami modelling (Baba et al., 2015;
526 <https://zenodo.org/record/6118212#.Yk98qdtBxPY>). Bathymetry data are obtained from GEBCO at
527 <http://www.gebco.net>. The sea level records in deep ocean are available from the Deep Ocean Assessment

528 and Reporting of Tsunamis (DART) buoy network in the Pacific (<https://nctr.pmel.noaa.gov/Dart/>), and
529 GeoNet New Zealand DART network (<https://tilde.geonet.org.nz>). The sea level records of tide gauges
530 are downloaded from UNESCO/ IOC (<http://www.ioc-sealevelmonitoring.org/>). Barometer data are
531 provided by the following providers: Direcció'n Meteoroló'gica de Chile
532 (<https://climatologia.meteochile.gob.cl>), NOAA National Weather Service
533 (<https://www.weather.gov/ilm/observations>), Japan Meteorological Agency (<https://www.jma.go.jp>),
534 The UK Met Office Weather Observation (<https://wow.metoffice.gov.uk/observations>), and Fiji
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