

The characteristics of the 2022 Tonga volcanic tsunami in the Pacific Ocean

Gui Hu¹, Linlin Li^{1,2}, Zhiyuan Ren³, Kan Zhang¹

1. Guangdong Provincial Key Laboratory of Geodynamics and Geohazards, School of Earth Sciences and Engineering, Sun Yat-sen University, Guangzhou, China

2. Southern Marine Science and Engineering Guangdong Laboratory (Zhuhai), Zhuhai, China

3. Department of Civil and Environmental Engineering, National University of Singapore, Singapore.

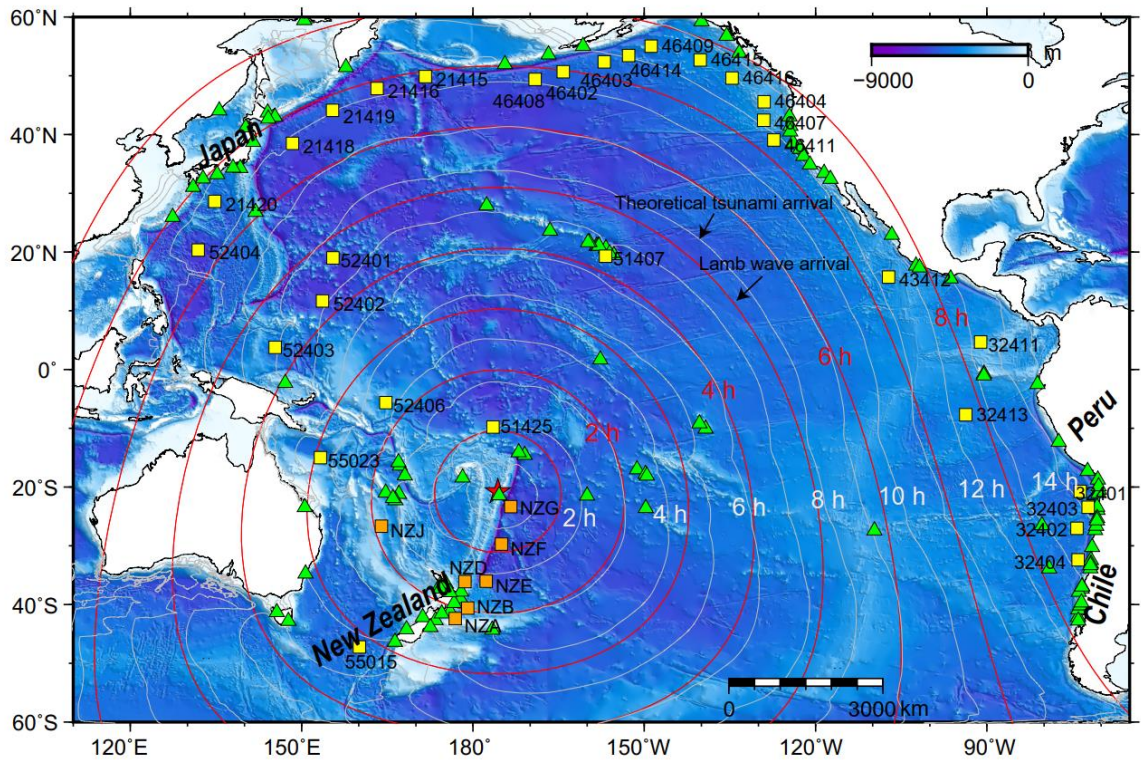
Correspondence to: Linlin Li (lilinlin3@mail.sysu.edu.cn)

Abstract. On 15th January 2022, an exceptional eruption of Hunga Tonga–Hunga Ha’apai volcano generated atmospheric and tsunami waves that were widely observed at oceans globally, gaining a remarkable attention to scientists in related fields. The tsunamigenic mechanism of this rare event remains an enigmatic due to its complexity and lacking of direct underwater observations. Here, to explore the tsunamigenic mechanisms of this volcanic tsunami event and its hydrodynamic processes in the Pacific Ocean, we conduct tsunami waveform and spectral analyses of the waveform recordings at 116 coastal gauges and 38 deep-ocean buoys across the Pacific Ocean. Combined with the constraints of some representative barometers, we obtain the plausible tsunamigenic origins during the volcano activity. We identify four distinct tsunami wave components generated by air-sea coupling and seafloor crustal deformation. Those tsunami components are differentiated by their different propagating speeds or period bands. The first-arriving tsunami component with ~80–100 min period was from shock waves spreading at a velocity of ~1000 m/s in vicinity of the eruption. The second component with extraordinary tsunami amplitude in deep sea was from Lamb waves. The Lamb wave with ~30–40 min period radically propagated outward from the eruption site with spatially decreasing propagation velocities from ~340 m/s to ~315m/s. The third component with ~10–30 min period was probably from some atmospheric gravity wave modes propagating faster than 200 m/s but slower than Lamb waves. The last component with ~3–5 min period originated from partial caldera collapse with dimension of ~0.8–1.8 km. Surprisingly, the 2022 Tonga volcanic tsunami produced long oscillation in the Pacific Ocean which is comparable with those of the 2011 Tohoku tsunami. We point out that the long oscillation is not only associated with the resonance effect with the atmospheric acoustic-gravity waves, but more importantly the interactions with local bathymetry. This rare event also calls for more attention to the tsunami hazards produced by atypical tsunamigenic source, e.g., volcanic eruption.

30 **1. Introduction**

31 On 15 January 2022 at 04:14:45 (UTC), a submarine volcano erupted violently at the uninhabited Hunga
32 Tonga-Hunga Ha'apai (HTHH) island at 20.546°S 175.390°W (USGS, 2022). The volcano is located ~67
33 km north of Nuku'alofa, the capital of Tonga (NASA, 2022) (Figure 1). The blasts launched plumes of
34 ash, steam, and gas ~58 km high into stratosphere (Yuen et al., 2022) which not only blanketed nearby
35 islands in ash (Duncombe, 2022; NASA, 2022), but caused various atmospheric acoustic-gravity wave
36 modes (AGWs) of various scales, e.g., Lamb waves from atmospheric surface pressure disturbance
37 associated with the eruption (Liu and Higuera, 2022; Adam, 2022; Kubota et al., 2022; Matoza et al.,
38 2022). Tsunami with conspicuous sea level changes were detected by coastal tide gauges and Deep-ocean
39 Assessment and Reporting of Tsunamis (DART) buoy stations in the Pacific (Figure 1), the Atlantic, and
40 Indian Oceans as well as the Caribbean and Mediterranean seas (Carvajal et al., 2022; Kubota et al., 2022;
41 Ramírez-Herrera et al., 2022), while the large waves were mainly concentrated in the Pacific Ocean, like
42 coastlines of New Zealand, Japan, California, and Chile (Carvajal et al., 2022). The event caused at least
43 3 fatalities in Tonga. Two people drowned in northern Peru when ~2 m destructive tsunami waves
44 inundated an island in the Lambayeque region, Chile (Edmonds, 2022).

45 Satellite images revealed that the elevation of HTHH island has gone through dramatic change before
46 and after the mid-January 2022 eruption. Previously, after the 2015 eruption, the two existing Hunga
47 Tonga and Hunga Ha'apai Islands were linked together. The volcanic island rose 1.8 km from the seafloor
48 where it stretched ~20 km across and topped a underwater caldera ~5 km in diameter (Garvin et al., 2018;
49 NASA, 2022). After the violent explosion on 15 January 2022, the newly formed island during 2015 was
50 completely gone, with only small tips left in far southwestern and northeastern HTHH island (NASA,
51 2022). HTHH volcano lies along the northern part of Tonga–Kermadec arc, where the Pacific Plate
52 subducts under the Indo-Australian Plate (Billen et al., 2003). The convergence rate (15~24 cm/year)
53 between the Tonga-Kermadec subduction system and the Pacific plate is among the fastest recorded plate
54 velocity on Earth, forming the second deepest trench around the globe (Satake, 2010; Bevis et al., 1995).
55 The fast convergence rate contributes to the frequent earthquakes, tsunamis and volcanic eruptions in
56 this region historically (Bevis et al., 1995). The 2022 HTHH volcano is part of a submarine-volcano
57 chain that extends all the way from New Zealand to Fiji (Plank et al., 2020). HTHH volcano had many
58 notable eruptions before 2022 since its first historically recorded eruption in 1912, i.e., in 1937, 1988,



60
 61 **Figure 1. The spatial distribution of the eruption site (red star), DART stations (squares), tide**
 62 **gauges (triangles) and the calculated tsunami arrival times. White contours indicate the modelled**
 63 **arrival times of conventional tsunami. Red contours indicate the estimated arrival times of Lamb**
 64 **waves (see how we derive these contours in section 3.1).**

65 The 2022 HTHH eruption is the first volcanic event which generates worldwide tsunami signatures since
 66 the 1883 Krakatau event (Matoza et al., 2022; Self and Rampino, 1981; Nomanbhoy and Satake, 1995).

67 The tsunamigenic mechanism of this rare volcanic eruption-induced tsunami is still poorly understood
 68 due to its complex nature and the deficiencies of near-field seafloor surveys. Various tsunami generation
 69 mechanisms have been proposed so far based on the observations of ground-based and spaceborne
 70 geophysical instrumentations (Kubota et al., 2022; Matoza et al., 2022; Carvajal et al., 2022). The
 71 mechanisms are closely associated with the air-sea coupling with atmospheric waves. Atmospheric
 72 waves propagating in the atmospheric fluid are generated by different physical mechanisms (Gossard
 73 and Hooke, 1975a). Lamb wave is a horizontally propagating acoustic waves in Lamb mode
 74 which is trapped at the earth's surface with group velocities close to the mean sound velocity of
 75 the lower atmosphere (e.g. Lamb, 1932). Atmospheric gravity wave is triggered when air
 76 molecules in the atmosphere are disturbed vertically other than horizontally (e.g. Le Pichon et
 77 al., 2010). Nonlinear propagation of atmospheric wave may cause period lengthening and the

78 formation of shock-wave (Matoza et al., 2022). The most-mentioned mechanism of the tsunami is
79 the fast-traveling atmospheric Lamb wave generated by the atmospheric pressure rise of ~ 2 hPa during
80 the eruption. The Lamb wave circled the Earth for several times with travelling speed close to that of the
81 sound wave in the lower atmosphere, leading to globally observed sea level fluctuations (Adam, 2022;
82 Duncombe, 2022; Kubota et al., 2022; Matoza et al., 2022) (Figure 1). The second mechanism is
83 suggested to be a variety of other acoustic-gravity wave modes (Adam, 2022; Matoza et al., 2022;
84 Themens et al., 2022; Zhang et al., 2022). The third mechanism may be related to the seafloor crustal
85 deformation induced by one or more volcanic activities in the vicinity of the eruption site (e.g.,
86 pyroclastic flows, partial collapse of the caldera) (Carvajal et al., 2022), which are more responsible for
87 the near-field tsunamis with theoretical tsunami speeds.

88 To investigate the possible tsunamigenic mechanisms and detailed hydrodynamic behaviors of this rare
89 volcanic tsunami event, in this study, we collect, process and analyze the sea level measurements from
90 116 tide gauge and 38 DART buoys in the Pacific Ocean (shown in Figures 1 and 2). We first do statistical
91 analysis of the tsunami waveforms to estimate the propagating speed of the Lamb wave and to understand
92 the tsunami wave characteristics in the Pacific Ocean through demonstrating the tsunami wave properties,
93 i.e., arrival times, wave heights and durations. We then conduct wavelet analysis for representative DART
94 buoys and tide gauges respectively to explore tsunamigenic mechanisms of the event and to better
95 understand its hydrodynamic processes in the Pacific Ocean. Aided by wavelet analysis of corresponding
96 barometers near the selected DART buoys and comparison with tsunami records of the 2011 Tohoku
97 tsunami, we are able to piece together all the analysis and demonstrate that the 2022 HTHH tsunami was
98 generated by air-sea coupling with a wide range of atmospheric waves with different propagating
99 velocities and period bands, and seafloor crustal deformation associated with the volcanic eruption. We
100 demonstrate as well that the tsunami was amplified at the far-field Pacific coastlines where the local
101 bathymetric effects play a dominant role in tsunami scale.

102 **2. Data and Methods**

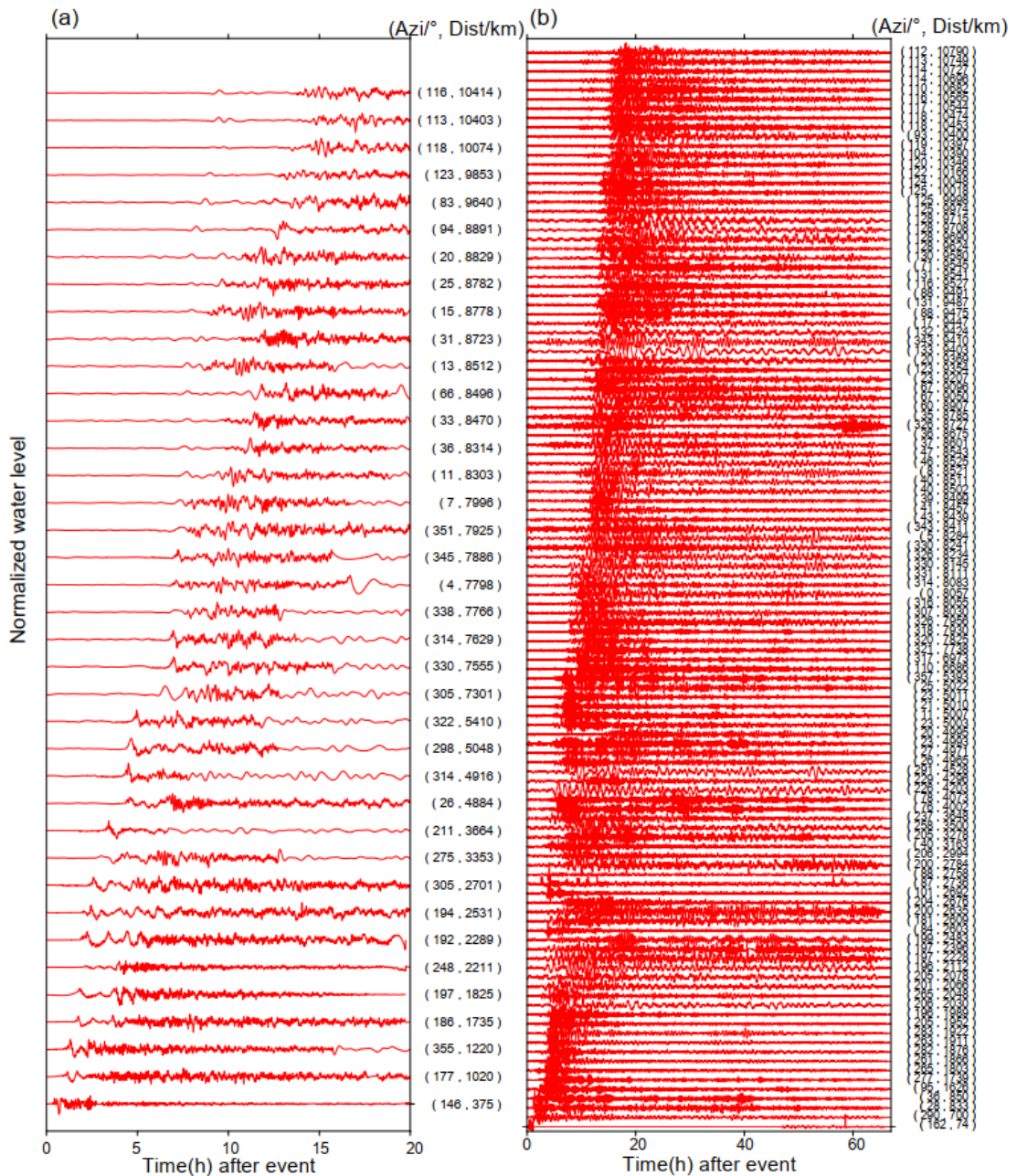
103 **2.1 Data**

104 We collected high-quality sea level records across the Pacific Ocean at 38 DART buoys (in which 31
105 stations from <https://nctr.pmel.noaa.gov/Dart/>, 7 stations from <https://tilde.geonet.org.nz/dashboard/>) and

106 116 tide gages from IOC (The Intergovernmental Oceanographic Commission, [http://www.ioc-](http://www.ioc-sealevelmonitoring.org)
107 [sealevelmonitoring.org](http://www.ioc-sealevelmonitoring.org)) (Figure 1). The epicentral distances of tide gauges and DART buoys range
108 between 74–10790 km and 375–10414 km, respectively. The sampling rates of DART buoys are
109 changing over time. Passing of tsunami event generally can trigger the DART system to enter its high
110 frequency sampling mode (15 seconds or 1 min) from normal frequency mode (15 min)
111 (www.ndbc.noaa.gov/dart). In contrast, sampling rates of normal tide gauges at coasts are uniform with
112 sampling interval of 1 min. The sampling interval of both DART and tide gauges is preprocessed to 15
113 seconds. Firstly, we eliminate abnormal spikes and fill gaps by linear interpolation. Secondly, we applied
114 a fourth-order Butterworth-Highpass filter with a cut-off frequency of 3.5×10^{-5} Hz (~ 8 hours) to remove
115 the tidal components (Figure 2) (Heidarzadeh and Satake, 2013). After the two steps, quality control step
116 is conducted to select high-quality data, in which we delete waveforms with spoiled data or massive data
117 loss due to equipment failure, or with the maximum tsunami heights of tide gauges less than 0.2 m, then
118 the selected data will be ready for further statistics and spectral analysis. We also collect and analyze the
119 atmospheric pressure disturbance data recorded by some representative barometers. The sampling rates
120 of the barometers is generally uniform with a sampling rate of 1 min except for some stations in New
121 Zealand with interval of 10 min. Considering the sample rate, we employ a fourth-order Butterworth-
122 Bandpass filter with period ranging between 2–150 min for wavelet analysis of the barometers with 1
123 min sample rate, while we apply the fourth-order Butterworth-Bandpass filter with range of 30–150 min
124 to long-period waveform display based on two reasons. (1) The barometer data we use for the analysis
125 include some in New Zealand with 10 min sample rate; (2) Filtering out the short-period waves helps
126 highlight long-period tsunami wave components.

127 The tsunami waveforms recorded by DART buoys which are installed offshore in the deep water are
128 expected to contain certain characteristics of the tsunami source (Wang et al., 2020, 2021). The
129 waveforms recorded by tide gauge distributed along coastlines are significantly influenced by local
130 bathymetry/topography which are used for investigating bathymetric effect on tsunami behaviors
131 (Rabinovich et al., 2017, 2006; Rabinovich, 2009). Therefore, we use the DART data for source-related
132 analysis and choose some tide gauge data to investigate the tsunami behaviors at the Pacific coastlines.

133



134
 135 **Figure 2. Detided tsunami waveforms at (a) DART buoys and (b) tide gauges. Waveforms in both**
 136 **subplots are shown in ascending distance. Azi stands for azimuth. The data are normalized with**
 137 **respect to the largest amplitude of each tide gauge.**

138 **2.2 Tsunami Modelling**

139 We use a numerical tsunami modelling package JAGURS (Baba et al. 2015) to simulate the tsunami
 140 propagation of the 2022 HTHH event and obtain the theoretical tsunami arrival time based on the shallow
 141 water wave speed (white contours in Figure 1). The code solves linear Boussinesq-type equations in a
 142 spherical coordinate system using a finite difference approximation with the leapfrog method. We specify
 143 a unit Gaussian-shaped vertical sea surface displacement at the volcanic base as the source of

144 conventional tsunami. For a unite source i with center at longitude φ_i and latitude θ_i , the
145 displacement distribution $Zi(\varphi, \theta)$ can be expressed as:

$$146 \quad Zi(\varphi, \theta) = \exp\left[-\frac{(\varphi-\varphi_i)^2+(\theta-\theta_i)^2}{2\sigma}\right] \quad (1)$$

147 Where we set characteristic length σ as 5 km (NASA, 2022). The bathymetric data is resampled from the
148 GEBCO 2019 with 15 arc-sec resolution (The General Bathymetric Chart of the Oceans, downloaded
149 from <https://www.gebco.net>).

150 **2.3 Spectral Analysis of Tsunami Waves**

151 To investigate the temporal changes of the dominant wave periods, we conduct continuous wavelet
152 transformation (frequency-time) analyses for some representative DART buoys, tide gauges and
153 barometers, in which wavelet Morlet mother function is implemented (Kristeková et al., 2006). The first
154 32-hour time series of DART buoys and barometers after the eruption (at 04:14:45 on 15 January 2022)
155 are used for source-related wavelet analysis. The first 48-hour time series of tide gauges after the eruption
156 are employed for hydrodynamics-related wavelet analysis at coastlines. We adopt the Averaged-Root-
157 Mean-Square (ARMS) method as a measure of absolute average tsunami amplitude with a moving time
158 window of 20 min to calculate the tsunami duration (Heidarzadeh and Satake, 2014). We define the time
159 durations as the time period where ARMS levels of tsunami waves are above those prior to the tsunami
160 arrivals.

161 **3. Results**

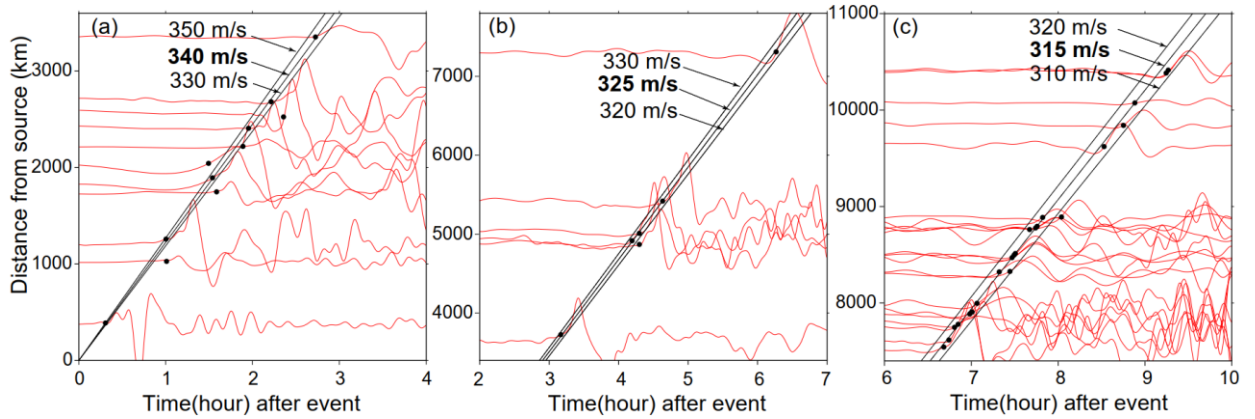
162 **3.1 The decreasing propagation velocities of the Lamb Wave**

163 Although many types of atmospheric waves were generated by the 2022 HTHH eruption, the most
164 prominent signature was the Lamb waves which were globally observed by ground-based and spaceborne
165 geophysical instrumentations (Kulichkov et al., 2022; Liu et al., 2022; Lin et al., 2022; Matoza et al.,
166 2022; Themens et al., 2022; Adam, 2022; Kubota et al., 2022). Interestingly, we notice that a wide range
167 of the velocities from 280 m/s to 340 m/s were proposed through observations and Lamb wave modelling
168 (e.g., Kubota et al., 2022; Lin et al., 2022; Matoza et al., 2022; Themens et al., 2022). The travelling
169 velocity of Lamb waves in real atmosphere is affected by temperature distributions, winds and dissipation
170 (Otsuka, 2022). To investigate whether the propagation speeds of the lamb wave change in space and

171 time, we analyze the waveforms recorded by the DART buoys in the Pacific Ocean. DART buoy with
172 pressure sensor deployed at the ocean's bottom records the sea level change that is transferred from
173 pressure records in Pascals, instead of direct water height. For the 2022 HTHH tsunami event, the
174 pressure fluctuation at DART buoy is a superposition of the pressure changes caused by tsunami and the
175 Lamb wave (Kubota et al., 2022). The Pacific DART buoys recorded the most discernible air-sea
176 coupling pulse in deep ocean with Lamb waves that arrived earlier than the theoretical tsunamis (Figure
177 1). The tsunami waveforms recorded by tide gauges did not clearly detect the tsunami signals associated
178 with the Lamb waves, therefore are not sufficient for further analysis (Figure 2). Thus, we estimate the
179 speed of Lamb waves using the waveforms recorded by the Pacific DART buoys. The Lamb wave
180 arrivals are limited within arrival time range from possible velocities of 280–340 m/s. The time points at
181 which the tsunami amplitudes first exceed 1 e-4 m above sea level are defined as Lamb wave arrivals.
182 By carefully fitting the arrivals with different constant velocities, we illustrate the velocities of Lamb
183 wave were generally uniform, but slightly decrease with the increase of propagation distance (Figure 3).
184 The Lamb waves initially propagated radially at speed of ~340 m/s before slowing to ~325 m/s after
185 reaching ~3400 km, and further decreasing to ~315 m/s at 7400 km. In an isothermal troposphere
186 assumption, the phase velocity of the Lamb wave (C_L) can be estimated with the following equation
187 (Gossard and Hooke, 1975b):

$$188 \quad C_L = \sqrt{\frac{\gamma \cdot R \cdot T}{M}} \quad (2)$$

189 Where $\gamma = 1.4$ (air specific heat ratio corresponding to atmospheric temperature), $R = 8314.36 \text{ J kmol}^{-1}$
190 K^{-1} (the universal gas constant), $M = 28.966 \text{ kg kmol}^{-1}$ (molecular mass for dry air) are constant for the
191 air, T is the absolute temperature in kelvin. Thus, Lamb wave velocity is mainly affected by the air
192 temperature, meaning the travelling velocity of lamb waves might decrease when propagating from
193 regions with high temperature towards those with low temperatures, e.g., the north pole. By assuming a
194 set of possible temperatures in January (Table 1), we calculated the velocities C_L could range between
195 312–343 m/s when temperatures vary between -30–20 °C. Therefore, the decreased velocity of the Lamb
196 waves could be a consequence of cooling of the air temperature.



197
 198 **Figure 3. Fitting the arrival times of normalized Lamb-induced tsunami waveforms with different**
 199 **velocities. Black dots mark the arrival times of the Lamb waves. Black lines represent velocities.**

200 **Table 1. Estimated Lamb wave velocities in an isothermal troposphere assumption**

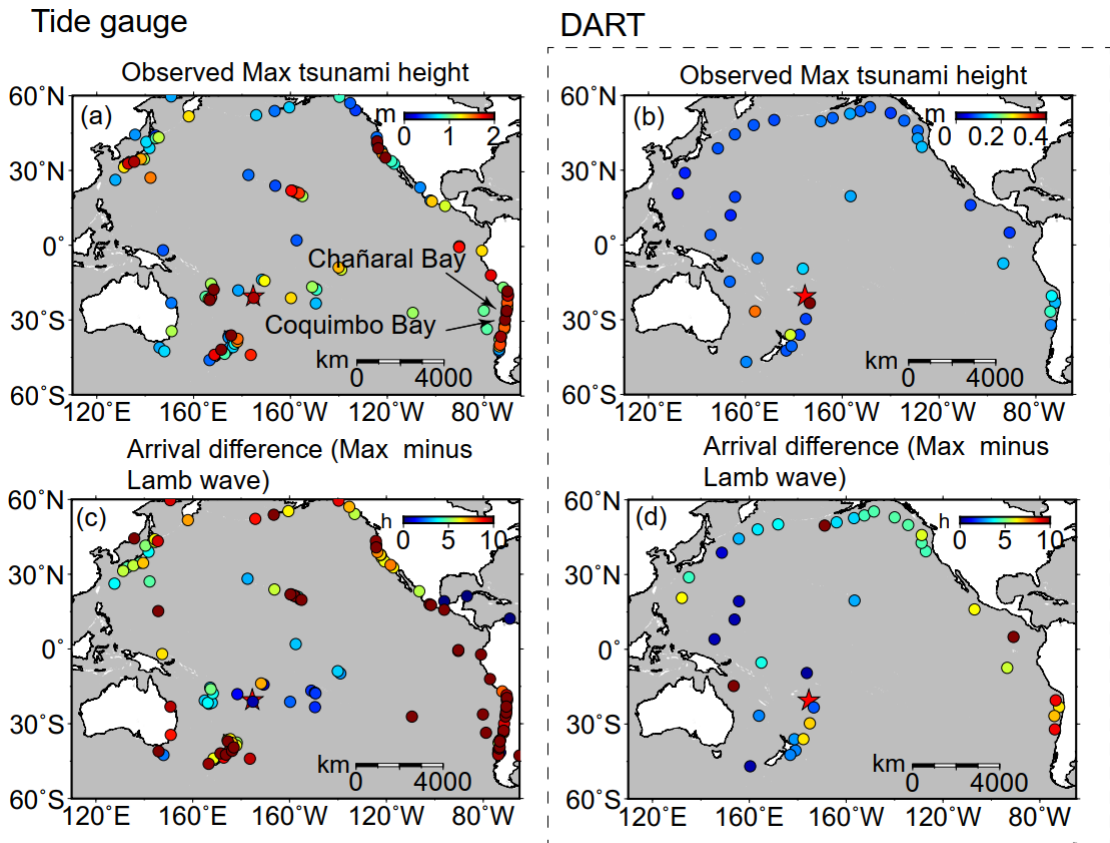
Celsius temperature (°C)	thermodynamic temperature (K)	C_L (m/s)
20	293.15	343.14
10	283.15	337.23
0	273.15	331.21
-10	263.15	325.19
-20	253.15	318.86
-30	243.15	312.49

201 **3.2 Tsunami features observed by DART buoys and Tide gauges**

202 The statistics of tsunami heights and arrival times recorded at 38 DART buoys and 116 tide gauges across
 203 the Pacific Ocean are used to interpret the tsunami characteristics. The comparison of the statistical
 204 characters between DART and tide gauge observations yields some useful information of the
 205 hydrodynamic process of tsunami propagation and help identify tsunami wave components with different
 206 traveling velocities.

207 The average value of the maximum tsunami wave height (trough-to-crest) for the 116 tide gauge stations
 208 is ~ 1.2 m. Figure 4a shows tide gauges with large tsunami heights exceeding 2 m are mainly distributed
 209 in coastlines with complex geometries (Figure S1a), such as gauges at New Zealand, Japan, and north
 210 and south America. For example, the largest tsunami height among tide gauges is 3.6 m at a bay-shaped
 211 coastal area Chañaral in Chile (Figure S1b). In sharp contrast to tide gauges, the maximum tsunami
 212 heights of most Pacific DART buoys are less than 0.2 m. The largest tsunami height in the DART buoys
 213 is only ~ 0.4 m recorded at the nearest one, 375 km from the volcano (Figure 4b). The comparison between
 214 DART buoys and tide gauges indicate that the direct contribution of air-sea coupling to the tsunami

215 heights is probably in the level of tens of centimeters (Kubota et al., 2022). The meter-scale tsunami
216 heights at the coastlines suggest the bathymetric effect could play a major role during tsunami
217 propagation. In respect to the arrival of maximum tsunami waves, the time lags between Lamb waves
218 and the maximum heights of tide gauges mainly range between ~0–10 h (Figure 4c). The delayed times
219 of ~10 h are observed in New Zealand, Hawaii, and west coast of America (Figure 4c), suggesting the
220 interaction between tsunami waves and local topography/bathymetry delays the arrival of the maximum
221 waves (e.g., Hu et al., 2022). For example, the delayed maximum tsunami height can be attributed to the
222 edge waves (Satake et al., 2020) and resonance effect (Wang et al., 2021) from tsunami interplays with
223 bays/harbors, islands, and continental shelves of various sizes. The significant regional dependence of
224 the coastal tsunami heights and the time lags of the maximum tsunami waves can be attributed to the
225 complexity of local bathymetry, such as continental shelves with different slopes, and harbor/bay with
226 different shapes and sizes (Satake et al., 2020). On the other hand, for tsunami events with earthquake
227 origins (e.g. Heidarzadeh and Satake, 2013), the first waves recorded by DART buoys are normally
228 observed as the largest wave since DART buoys are located in the deep sea and less influenced by
229 bathymetric variation. In the case of Tonga tsunami event, we observe the inconsistency between the
230 arrivals of the Lamb wave-induced tsunami waves and the maximum tsunami heights (Figure 4d). The
231 time lags of the maximum waves of DART buoys present a coarsely increasing tendency with the
232 increasing distance from the volcano, which indicates the contribution of other tsunami generation
233 mechanism propagating with a uniform but lower speed than Lamb wave.



234
 235 **Figure 4. The spatiotemporal signatures of the 2022 HTHH tsunami across the Pacific Ocean. (a)**
 236 **Observed the maximum tsunami height (trough-to-crest height) of tide gauges. (c) Arrival**
 237 **differences between the maximum tsunami height of tide gauges and Lamb waves. (b) and (d) are**
 238 **the same as (a) and (c) but for DART buoys.**

239 3.3 Tsunami components identified from wavelet analysis

240 The statistical analysis of tsunami waveforms at tide gauges and DART buoys suggest the tsunami waves
 241 likely contain several components with different source origins. To further identify these tsunami
 242 components, we conduct wavelet analysis for tsunami waveforms recorded by representative DART
 243 buoys and air pressure waveforms recorded by selected barometers. We demonstrate the analysis result
 244 through the frequency-time (f-t) plot of wavelet which shows how energy and period vary at frequency
 245 and time bands (Figure 5 and Figure 6). Tsunami components have clear signatures in all f-t plots as the
 246 energy levels are quite large when they arrive. Figure 5 shows the wavelet analysis of six DART buoys
 247 located in the vicinity of the eruption site (<3664 km). Figure 6 show the wavelet analysis of ten DART
 248 buoys located in the Pacific rim which are far away from the source location. We observe three interesting
 249 phenomena: 1) most of the tsunami wave energy is concentrated in four major period bands, i.e., 3–5
 250 min, ~10–30 min, ~30–40 min, and ~80–100 min; 2) The significant tsunami component with period

251 band of 3-5 mins are recorded by stations between the eruption site and the north tip of the New Zealand;
 252 3) There exists one exceptional tsunami component with longer wave period of ~80–100 min mainly
 253 recorded in the Tonga, the New Zealand and Hawaii, which travels even faster than the Lamb waves.
 254 To further explore the source mechanism of these tsunami components, we take advantage of the
 255 published information related to different propagating velocities of atmospheric gravity waves (Kubota
 256 et al., 2022) and add four kinds of propagating velocities as criteria to differentiate the tsunami arrivals
 257 from different sources (Figure 5 and Figure 6). The first reference speed is 1000 m/s related to the
 258 radically propagating atmospheric shock waves near the source region (Matoza et al., 2022; Themens et
 259 al., 2022). The second one is the velocities of Lamb wave ranging between 315–340 m/s derived from
 260 the aforementioned analysis in section 3.1 (Figure 3). The third one is 200 m/s corresponding to the lower
 261 limit of atmospheric gravity wave modes other than Lamb waves which were also excited by the volcanic
 262 eruptions (Kubota et al., 2022). The last is the arrival time of conventional tsunami given by tsunami
 263 modelling (Figure 1). The theoretical velocity of conventional tsunami is significantly nonuniform
 264 spatially as compared with those of the atmospheric waves. The conventional tsunami propagation speed
 265 is determined by the water depth along the propagation route. The velocity of non-dispersion shallow-
 266 water waves (C_H) in the ocean is given by:

$$267 \quad C_H = \sqrt{g \cdot H} \quad (3)$$

268 Where g is gravity acceleration (9.81m/s^2), H is the water depth. The propagation velocities of tsunami
 269 are ~296–328 m/s in the deepest trenches on earth (i.e., ~11 km in Mariana Trench and ~9 km in Tonga
 270 Trench). The velocities decrease quickly to only ~44 m/s at ~200 m depth along the edge of continental
 271 shelf. With the average depth of ~4–5 km, the average velocities in the Pacific Ocean range between
 272 ~200–224 m/s. Thus, theoretical tsunami velocities present significant slowness and variability. We
 273 delineate the arrival times of the four reference speeds in Figures 5 and 6.

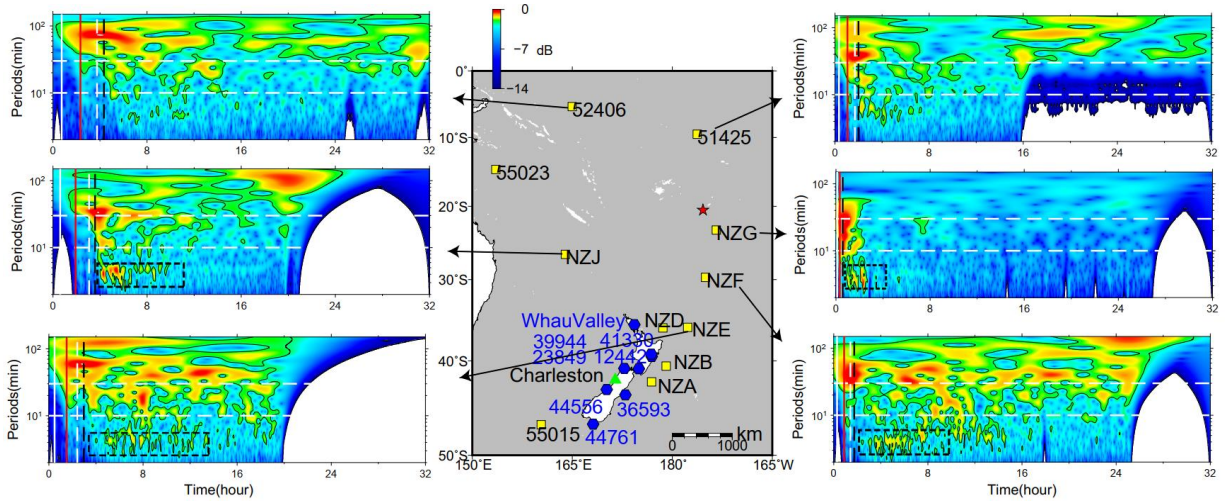
274 One particularly remarkable phenomenon is that the wave component with period of ~80–100 min
 275 propagated at a very fast speed of ~1000 m/s in the vicinity of the HTHH site, i.e., New Zealand and
 276 Hawaii (e.g., stations 52406, NZJ, NZE, 51425 in Figure 5, and 51407 in Fig. 6). We infer that the
 277 tsunami component within ~80–100 min period band was likely produced by the atmospheric shock
 278 waves during the initial stage of the volcanic eruption and spatially only cover the near-source region.
 279 To verify this observation, we select 16 representative barometers located in the near-source region and

280 far-field area for wavelet analysis (see the locations in Figure 5 and Figure 6). Figure 7 shows the
281 waveforms of atmospheric pressure at selected locations and Figure 8 provides the frequency-time (f-t)
282 plot of wavelet analysis of some representative barometers. Interestingly, we are able to discern the air
283 pressure pulses prior to Lamb waves at barometers in New Zealand (the two columns on the left in Figure
284 7), although such signals are not detectable in waveforms recorded by barometers far from the source
285 (the two columns on the right in Figure 7). The spatial distribution of such unusual pressure changes
286 suggest that the fast travelling shock waves were only limited in the near-source region, as reflected in
287 the travelling ionospheric disturbances (Matoza et al., 2022; Themens et al., 2022). Additionally, we also
288 see that the long period signals of ~80–100 min appear in DART buoys far away from the eruption site.
289 Such signals may be related with the long-period gravity waves (Matoza et al., 2022).

290 The tsunami components at period band of ~30–40 min can be readily associated with Lamb waves
291 because the arrival times of the tsunami waves and Lamb waves have excellent match, as shown in the
292 tsunami data recorded by DART buoys (e.g., NZJ and 51425 in Figure 5; 51407, 32401 and 32413 in
293 Figure 6) and pressure data by barometers (Figure 8).

294 For the tsunami components with the period band of ~10–30 min, although the arrivals of ~10–30 min
295 tsunami components cover some theoretical tsunami arrival times, they do not consistently match. The
296 tsunami components occurring within the time period between Lamb waves and the lower gravity waves'
297 velocities has a good agreement with the velocity range of several atmospheric gravity wave modes
298 (Matoza et al., 2022; Themens et al., 2022; Kubota et al., 2022). Similarly, the air pressure data also show
299 energy peaks at ~10–30 min period band, which is consistent with the tsunami data (Figure 8). Such
300 consistency further verifies the contribution of atmospheric gravity waves to the volcanic tsunami.

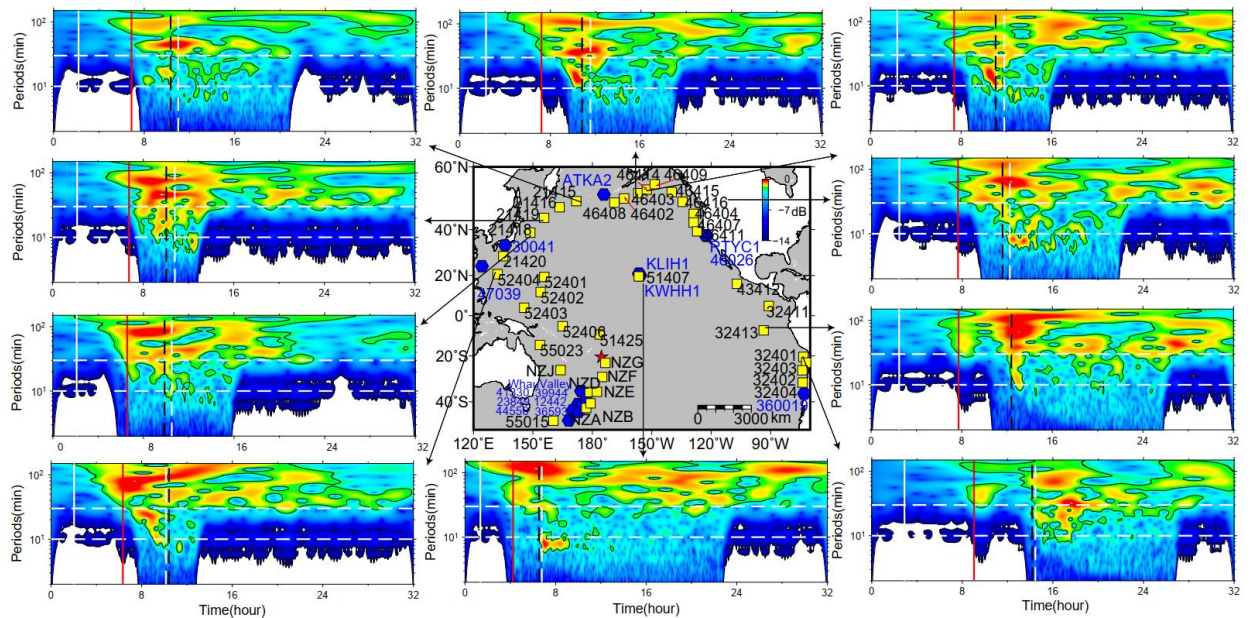
301 The tsunami components with the shortest period of ~3–5 min (stations NZE, NZF, NZG and NZJ;
302 marked with black dashed squares in Figure 5) are only observed at DART records near the eruption
303 location. Meanwhile, the arrival times of these components agree well with the modelled arrivals of
304 conventional tsunami. Thus, we believe the observed shortest period band should originate from the
305 seafloor crustal deformation. We further infer that this component could be generated by the partial
306 underwater caldera collapse and/or subaerial/submarine landslide failures associated with 2022 HTHH
307 volcanic eruption.



308

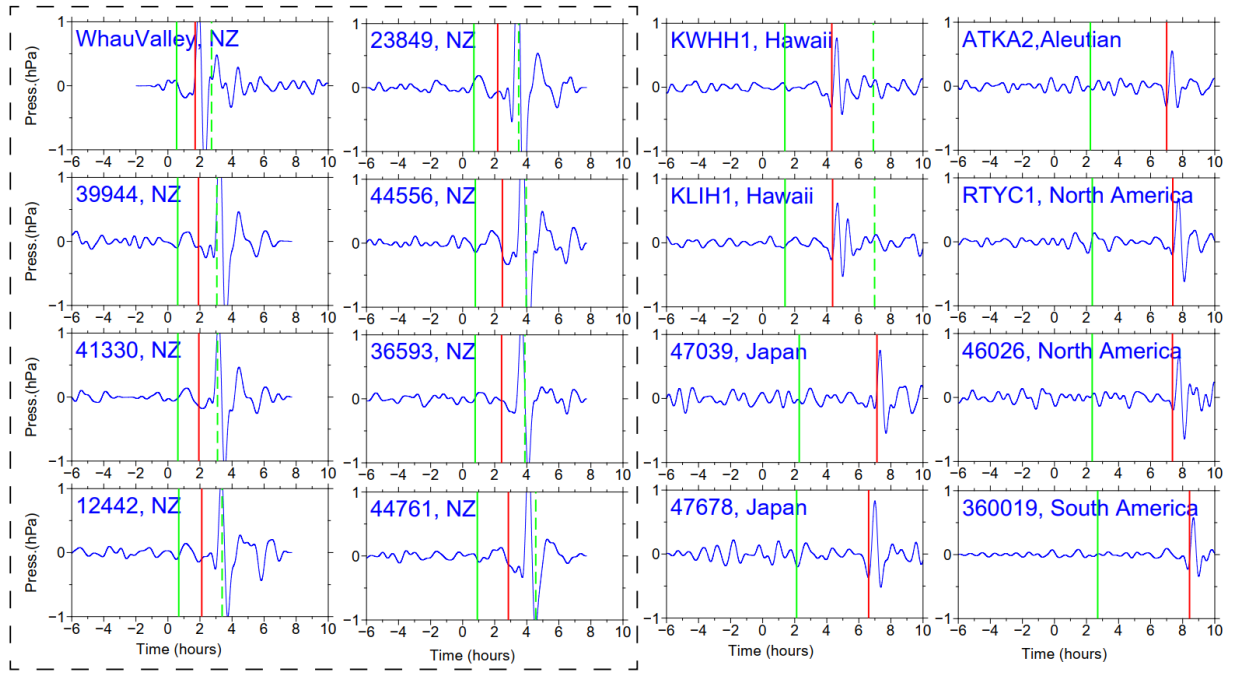
309 **Figure 5. Wavelet analysis of representative DART buoys in the vicinity of the HTHH volcano. In**
 310 **each sub-plot, the solid vertical white lines mark the arrival time with travelling velocity of 1000**
 311 **m/s. The solid vertical red lines mark the arrivals of Lamb waves. The dashed vertical white lines**
 312 **mark lower limit of AGWs' velocity of 200 m/s (Kubota et al., 2022). The dashed vertical black**
 313 **lines represent the theoretical tsunami arrivals. The dashed horizontal white lines mark two**
 314 **reference wave periods of 10 min and 30 min. The blue hexagons represent the locations of**
 315 **barometers. Green triangle makes the location of the tide gauges at Charleston. Decibel (dB) is**
 316 **calculated from: $dB = 10 \log(A/A_0)$, where A is wavelet power, A_0 is a reference wavelet power of**
 317 **the maximum one (Thomson and Emery, 2014).**

318



319

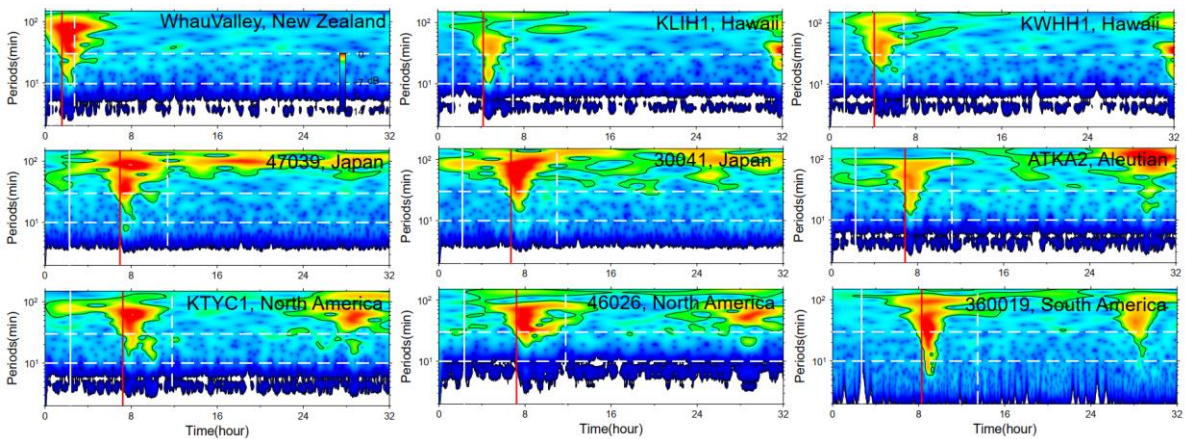
320 **Figure 6. Wavelet analysis of representative DART buoys far away from the HTHH volcano. In**
 321 **each sub-plot, the solid vertical white lines mark the arrival time with travelling velocity of 1000**
 322 **m/s. The solid vertical red lines mark the arrivals of Lamb waves. The dashed vertical white lines**
 323 **mark lower limit of AGWs' velocity of 200 m/s. The dashed vertical black lines represent the**
 324 **theoretical tsunami arrivals. The dashed horizontal white lines mark two reference wave periods**
 325 **of 10 min and 30 min. The blue hexagons represent the locations of barometers.**



326

327 **Figure 7. Shockwave-related atmospheric pressure waveforms of selected barometers in the Pacific**
 328 **Ocean. All traces have been filtered between 30 min and 150 min. In each sub-plot, the solid vertical**
 329 **green lines mark the arrival time with travelling velocity of 1000 m/s. The solid vertical red lines**
 330 **mark the arrivals of Lamb waves. The dashed vertical green lines mark lower limit of AGWs'**
 331 **velocity of 200 m/s.**

332



333

334 **Figure 8. Wavelet analysis of some representative barometers. In each sub-plot, the solid vertical**
 335 **white lines mark the arrival time with travelling velocity of 1000 m/s. The solid vertical red lines**
 336 **mark the arrivals of Lamb waves. The dashed vertical white lines mark lower limit of AGWs'**
 337 **velocity 200 m/s. The dashed horizontal white lines mark three reference periods of 10 min and 30**
 338 **min.**

339 **4. Discussion**

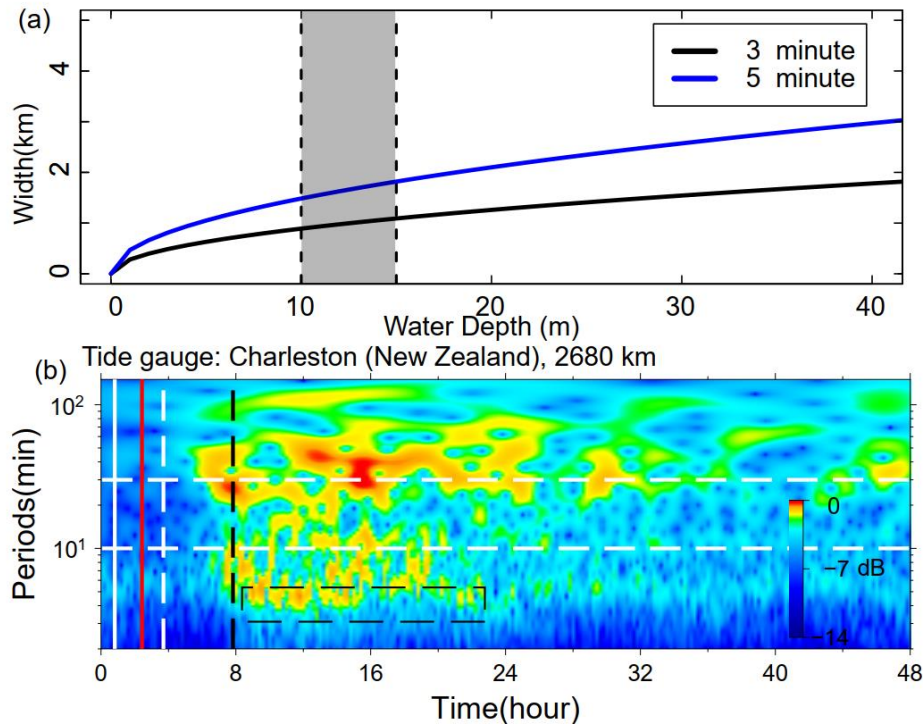
340 **4.1 Tsunami from Caldera Collapse and Its Long-distance Traveling Capability**

341 The tsunami wave energy distributed in different period bands is identified with reference arrival times.
342 The tsunami component with 3–5 min period is most likely generated by seafloor crustal deformation in
343 the volcanic site, but specific mechanism is not determined. A variety of possible scenarios associated
344 with the eruption could be responsible for the near-field tsunami waves, such as volcanic earthquakes,
345 pyroclastic flows entering the sea, underwater caldera flank collapse, and subaerial/submarine failures
346 (Self and Rampino, 1981; Pelinovsky et al., 2005). To further investigate the source mechanism, we
347 apply a simplified model (Rabinovich, 1997) to estimate the probable dimension of tsunami source:

$$348 \quad L = \frac{T\sqrt{gH}}{2} \quad (4)$$

349 Where L is the typical dimension (length or width) of the tsunami source, H is average water depth in the
350 source area, g is the gravity acceleration, and T is primary tsunami period. By comparing with the post-
351 2015 morphology of the HTHH caldera which was obtained through drone photogrammetry and
352 multibeam sounder surveys, Stern et al. (2022) estimate that much of the newly-formed Hunga Tonga
353 Island and the 2014/2015 cone were destroyed by the 2022 eruption, and the vertical deformation of
354 Hunga Ha’apai Island is ~10–15 m (Stern et al., 2022). With no more quantitative constraint of the
355 seafloor deformation, we tentatively assume H as 10–15 m, then the possible dimension of seafloor
356 crustal deformation responsible for the small-scale tsunami could be in the scale of 0.8–1.8 km (Figure
357 9a). The estimated size is very likely from partial caldera collapse that usually has limited scale in
358 volcanic site (Ramalho et al., 2015; Omira et al., 2022). If it is the case, the partial flank collapse could
359 be located between Hunga Tonga and Hunga Ha’apai Islands.

360



361
 362 **Figure 9. Mechanism of tsunami component with 3–5min period. (a) The source dimension**
 363 **estimated by equation 4. (b) Wavelet analysis of tide gauge at Charleston, New Zealand, 2680 km**
 364 **away from the eruption site. The solid vertical white line marks the arrival time with travelling**
 365 **velocity of 1000 m/s. The solid vertical red line marks the arrival of Lamb wave. The dashed**
 366 **vertical white line marks lower limit of AGWs' velocity 200 m/s. The dashed vertical black line**
 367 **marks the theoretical tsunami arrivals.**

368 An interesting phenomenon is that the tsunami component with 3–5 min period can still be observed in
 369 a bay-shaped coastal area at Charleston in New Zealand (see the location in Figure 5) which is 2680 km
 370 away from the eruption site and maintains a high energy level lasting up to 14 h (Figure 9b). The long-
 371 traveling capability could be associated with the ~ 10000 m deep water depth of the Tonga Trench that
 372 keeps the source signals from substantial attenuation. In deep open ocean, the wavelength of a tsunami
 373 can reach two hundred kilometers, but the height of the tsunami may be only a few centimeters. Tsunami
 374 waves in the deep ocean can travel thousands of kilometers at high speeds, meanwhile losing very little
 375 energy in the process. The long oscillation can be attributed to the multiple reflections of the incoming
 376 waves trapped in the shallow-water bay at Charleston.

377 Generally, devastating tsunamis with long-distance travelling capability are mostly generated by
 378 megathrust earthquakes (Titov et al., 2005). Caldera collapses or submarine landslides with limited scale
 379 normally only generate local tsunamis, e.g., the 1998 PNG (Papua New Guinea) tsunami event (Kawata
 380 et al., 1999) and the 1930 Cabo Girão tsunami event (Ramalho et al., 2015). Therefore, it's exceptional

381 that the tsunami component from scale-limited failure could travel at-least 2680 km away from the
382 eruption site. It demonstrates that tsunamis from small-scale tsunamigenic source have the capability to
383 travel long distance and cause long oscillation at favored condition, e.g., deep trench, ocean ridge and
384 bay-shaped coasts.

385 **4.2 The Possible Mechanisms of Long Tsunami Oscillation**

386 An important tsunami behavior of the 2022 HTHH tsunami is the long-lasting oscillation ~ 3 days in the
387 Pacific Ocean (Figure 10a), which is comparable to that of the 2011 Tohoku tsunami, ~ 4 days
388 (Heidarzadeh and Satake, 2013). We demonstrate the duration time of the tsunami oscillation through
389 ARMS (Averaged-Root-Mean-Square) approach that is a measure of absolute average tsunami amplitude
390 in a time period. The long-lasting tsunami energy can be observed at many regions, such as the coasts of
391 New Zealand, Japan, Aleutian, Chile, Hawaii, and west coasts of America. Several mechanisms could
392 account for the long-lasting tsunami, including (1) Lamb waves circling the Earth multiple times
393 (Amores et al., 2022; Matoza et al., 2022), (2) resonance effect between ocean waves and atmospheric
394 waves (Kubota et al., 2022), and (3) bathymetric effect. We discuss the contribution of each mechanism
395 in the following section.

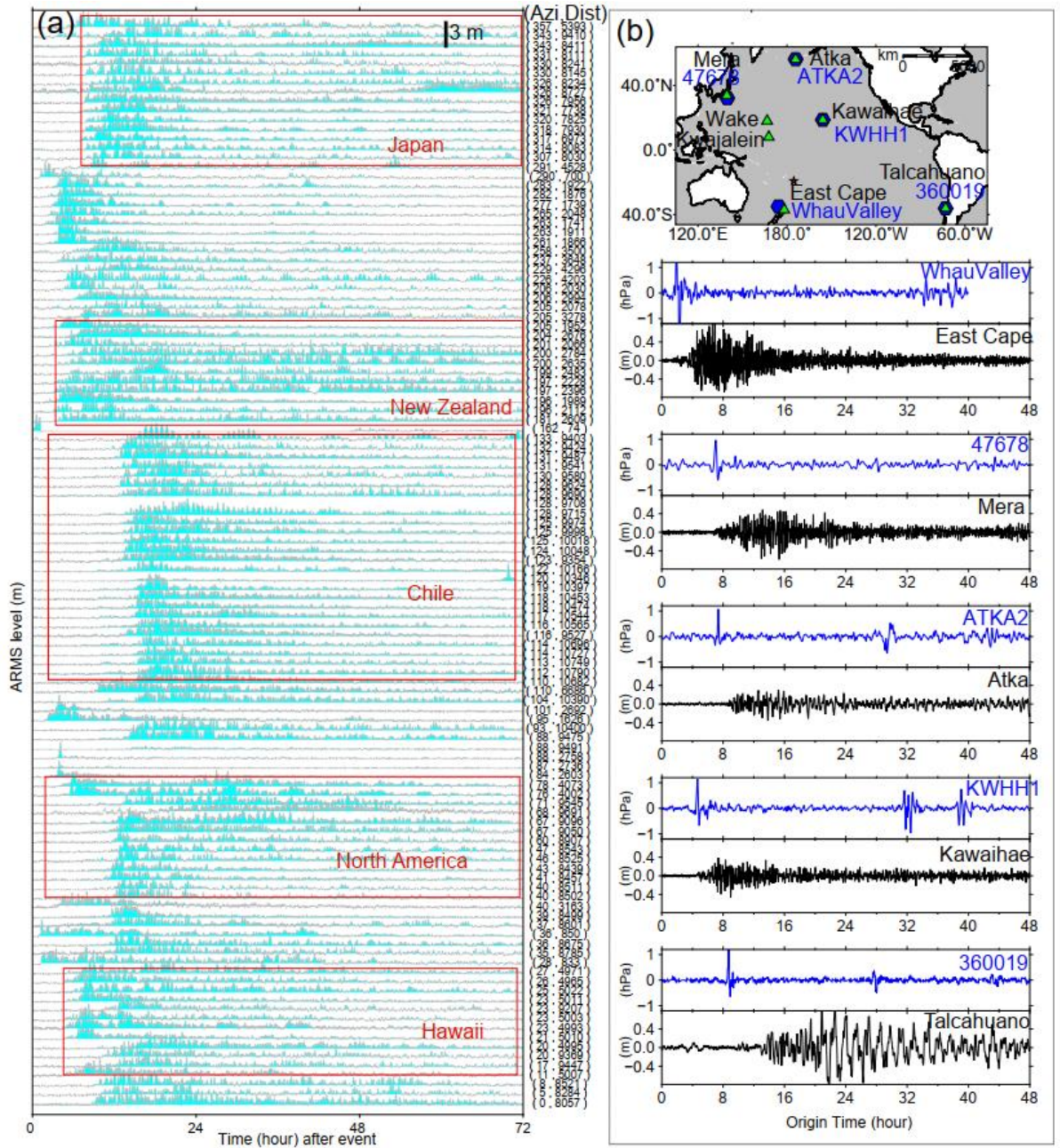
396 To investigate the contribution of Lamb wave to the long-lasting tsunami, we compare the air pressure
397 disturbances recorded by selected barometers together with the tsunami waveforms of nearby tide gauges
398 (Figure 10b). While the barometers present discernible wave pulses at each Lamb wave's arrival, only
399 the first Lamb wave triggered clear tsunami signal and no detectable tsunami signatures correspond to
400 the following passage, suggesting the Lamb waves do not directly contribute to the long oscillation.

401 The resonance effects between ocean waves and atmospheric waves could contribute to the long
402 oscillation on coastlines. Besides the Lamb wave, Watanabe et al., 2022 detected internal Pekeris wave
403 which propagate with a slower horizontal phase speed of ~ 245 m/s and gravity waves with even slower
404 propagation speed by analyzing radiance observations taken from the Himawari-8 geostationary satellite.
405 Atmospheric waves with such speeds are more likely to resonant with the conventional tsunami waves
406 and provide continuous energy supply (Kubota et al., 2022).

407 To examine the role of local bathymetry in the long-lasting tsunami, we choose a well-studied and well-
408 recorded event: the 2011 Mw 9.0 Tohoku tsunami as a reference event and compare the tsunami records
409 of these two events at the same coastal stations. Although the two tsunami events were generated by

410 completely different mechanisms, i.e., large-scale seafloor deformation for the Mw 9.0 megathrust
411 earthquake (Mori et al., 2011) and fast-moving atmospheric waves for the Mw 5.8 volcanic eruption
412 (Matoza et al., 2022), they both produced widespread transoceanic tsunamis which were well recorded
413 in the Pacific DART buoys and tide gauges. In the near-field, the 2011 Tohoku earthquake produced
414 runup up to 40 m at Miyako in the Iwate Prefecture in Japan's Tohoku region (Mori et al., 2011). The
415 epicenter is approximately 70 km east coast of the Oshika Peninsula of Tohoku region. However, the
416 2022 HTHH tsunami produced only ~13 m runup in the near field from eyewitness accounts in
417 Kanokupolu, 60 km from the volcano (Lynett et al., 2022). However, in the far-field (>1000 km), we
418 observe comparable tsunami wave heights in certain coastal regions. Based on the tsunami records at 21
419 tide gauges surrounding the Pacific Ocean, Heidarzadeh & Satake (2013) calculated the average value
420 of the maximum tsunami heights (trough-to-crest) of the 2011 Tohoku tsunami is 1.6 m with the largest
421 height of 3.9 m at the Coquimbo Bay in Chile (Heidarzadeh and Satake, 2013). Coincidentally, the statistics
422 of 116 tide gauges in this study also suggest the average tsunami heights of the 2022 HTHH tsunami is
423 around the same order, ~1.2 m, among which, the largest height is 3.6 m at Chañaral Bay in Chile.
424 Interestingly, in the coastal region of South America, the locations of the largest tsunami heights of both
425 events are adjacent (Figure 4a), i.e., Coquimbo (the 2011 Tohoku) and Chañaral (The 2022 HTHH).
426 To further compare the far-field hydrodynamic processes between these two events quantitatively, we
427 conduct wavelet analysis for four representative tide gauges distributed across the Pacific Ocean, i.e.
428 coastal gauges at East Cape in New Zealand, Kwajalein Island, Wake Island, and Talcahuaho in Chile
429 (see their locations in Figures 10b). The temporal changes of tsunami energy of both events can be seen
430 in Figure 11. At each tide gauge, the tsunami energy of the 2011 HTHH (Figure 11a) and the 2022 Tohoku
431 tsunamis (Figure 11b) for the first few hours after the arrivals is nonuniform with different significant
432 peaks distributed within a wide period band of ~3–100 min. Then, the following long-lasting energy of
433 the both at each station presents similar pattern and is concentrated at identical and fairly narrower period
434 channel, i.e., ~20–30 min at East Cape in New Zealand, ~40–60 min at Kwajalein Island, ~10 min at
435 Wake Island, and ~100 min at Talcahuaho in Chile, which reflects the local bathymetric effects of natural
436 permanent oscillations (Hu et al., 2022; Satake et al., 2020). Specifically, many bathymetric effects can
437 contribute to the long-lasting tsunami, such as multiple reflections across the basins, or the continental
438 shelves, and the excited tsunami resonance in bays/harbors with variable shapes and sizes (Aranguiz et

439 al., 2019; Satake et al., 2020). For example, tide gauges around New Zealand are primarily distributed in
440 harbors/ports with major natural oscillation modes of ~20–30 min (De Lange and Healy, 1986; Lynett et
441 al., 2022). The first oscillation mode of central Chile is centered around ~100 min (Aranguiz et al., 2019).
442 Consequently, Figure 11 illustrates that the long-lasting tsunami energy of the two events is respectively
443 distributed in 20–30 min period at East Cape in New Zealand and in ~100 min period at Talcahuaho in
444 central Chile. The coupling of bathymetric oscillation mode with tsunami containing similar-period wave
445 results in the excitement of tsunami resonance, which amplifies tsunami waves and prolongs the tsunami
446 oscillation at the two stations (Heidarzadeh et al., 2019, 2021; Hu et al., 2022; Wang et al., 2022).
447 Simply put, we do not have clear evidence that atmospheric acoustic-gravity waves from the 2022 HTHH
448 eruption directly contribute to the long-lasting tsunami, but the resonance effect associated with ocean
449 waves could a possible source of increased wave energy and amplification. However, the similarity of
450 far-field hydrodynamic behaviors between the 2022 HTHH volcanic tsunami and the 2011 Tohoku
451 seismogenic tsunami well demonstrates the both went through similar hydrodynamic processes after their
452 arrivals. The consistency favors that the long-lasting tsunami of 2022 HTHH tsunami event can very
453 likely be attributed by the interplays between local bathymetry and conventional tsunami left after each
454 passage of atmospheric waves, which can well explain why the two completely distinct tsunami events
455 possess a comparable duration time.

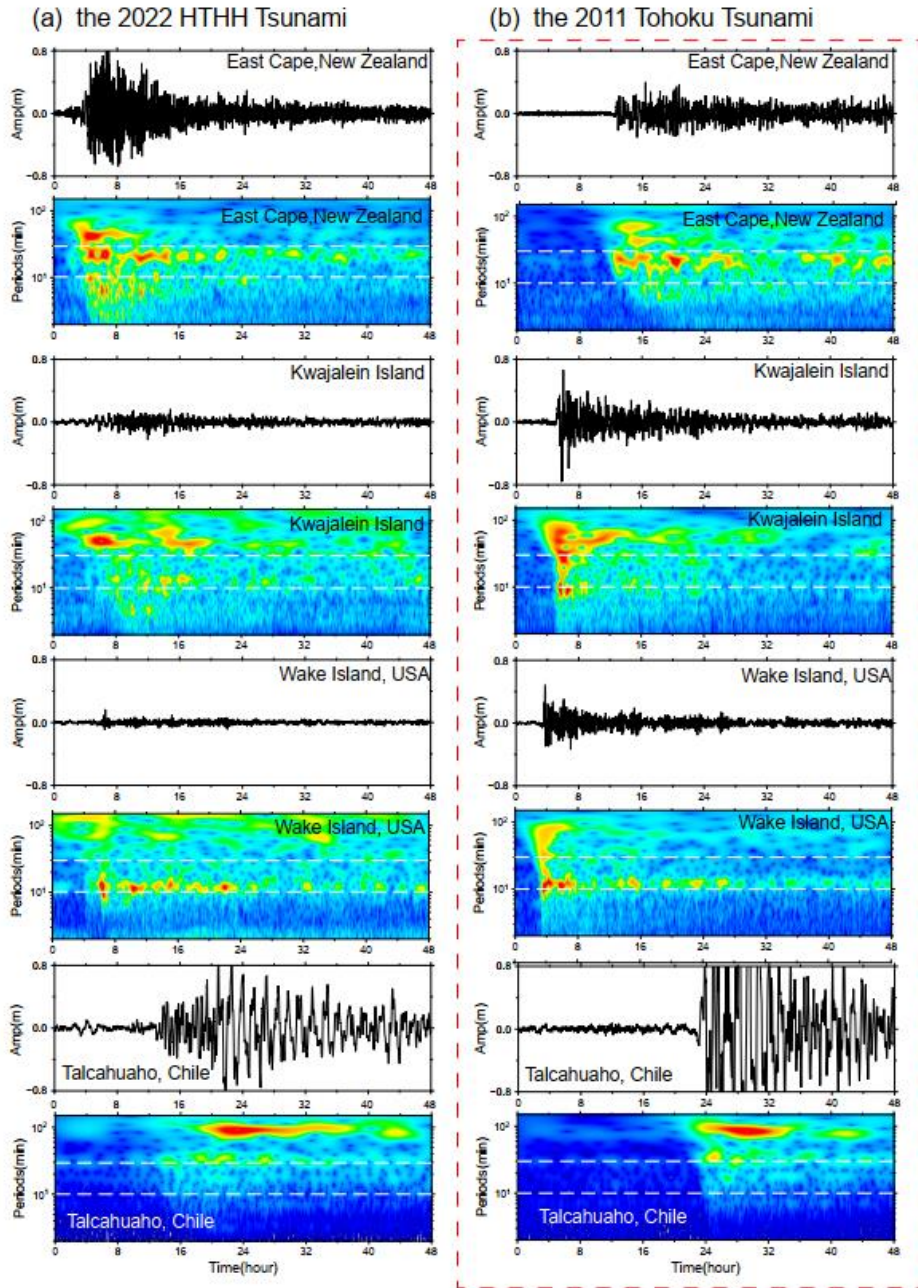


456

457 **Figure 10. Tsunami duration. (a) Tsunami durations at Pacific 116 tide gauges through ARMS level**
 458 **approach. (b) the location of barographs (blue curves) and nearby tide gauges (green curves), as**
 459 **well as their waveforms.**

460

461



462
 463 **Figure 11. Wavelet analysis of tsunami waveforms recorded by 4 tide gauges during (a) the 2022**
 464 **HTHH tsunami event, and (b) the 2011 Tohoku tsunami event. Horizontal white dashed lines**
 465 **respectively mark reference periods of 10 min and 30 min.**

466 **4.3 Challenges for Tsunami Warning**

467 The generation mechanisms and hydrodynamic characteristics of the 2022 HTHH volcanic tsunami are
 468 more complicated than pure seismogenic tsunami, which challenge the traditional tsunami warning
 469 approach.

470 The first challenge is posed by the tsunami components with propagating velocities faster than the

471 conventional tsunami. The Tonga volcanic tsunami event provides an excellent example which highlights
472 that the tsunamigenic mechanisms are not limited to tectonic activities related with the sudden seafloor
473 displacements, but also include a variety of atmospheric waves with distinct propagation velocities. The
474 tsunami components in 2022 HTHH event generated by the air-sea coupling possess a wide range of
475 velocities from 1000 m/s to 200 m/s. The Lamb waves recorded in both the 2022 HTHH event and the
476 1833 Krakatoa volcanic event traveled along the Earth's surface globally for several times (Carvajal et
477 al., 2022). The tsunami waves produced by Lamb waves, the wave components associated with resonance
478 of the air-sea coupling and their superimposition increase the difficulty of tsunami warning.
479 Another critical challenge is associated with the interplays between tsunami waves and local bathymetry.
480 The tsunami waves left by each passage of the atmospheric waves can interact with local bathymetry at
481 coastlines, such as continental shelves with different slopes, and harbor/bay with different shapes and
482 sizes. The interaction can intensify the tsunami impact and excite a variety of natural oscillation periods.
483 The 2022 HTHH tsunami with an extremely wide period range of ~2–100 min have a great potential to
484 couple with the excited natural oscillations and form extensive tsunami resonance phenomena. The
485 resonance effects result in long-lasting oscillation and delayed tsunami wave peaks. The uncertain
486 arrivals of the maximum tsunami waves pose an extra challenge to tsunami warning.

4875. **Conclusion**

488 In the study, we explore the tsunamigenic mechanisms and the hydrodynamic characteristics of the 2022
489 HTHH volcanic tsunami event. Through extensive analysis of waveforms recorded by the DART buoys,
490 tide gauges and barometers in the Pacific Ocean, we reach the main findings as follows:

491 (1) We identify four distinct tsunami wave components based on their distinct propagation velocities or
492 period bands (~80–100 min, 10–30 min, 30–40 min, and 3–5 min). The generation mechanisms of these
493 tsunami components range from air-sea coupling to seafloor crustal deformation during the volcanic
494 eruption.

495 (2) The first-arriving tsunami component with 80–100 min period was most likely from shock wave
496 spreading at a velocity of ~1000 m/s in the vicinity of the eruption. This tsunami component was not
497 clearly identified by currently available publication and it's not easy to be visually observed through time
498 series of the waveforms. The physical mechanism is yet to be understood. The second tsunami component

499 with 30–40 min period was from Lamb waves, and was the most discussed tsunami source of this event
500 so far. A thorough analysis of DART measurements indicates that the Lamb waves traveled at the speed
501 of ~340 m/s in the vicinity of the eruption and decreased to ~315 m/s when traveling away due to cooling
502 of the air temperature. The third tsunami component was from some atmospheric gravity wave modes
503 with propagation velocity faster than 200 m/s but slower than Lamb waves. The last tsunami component
504 with the shortest periods 3-5 min was probably produced by partial caldera collapse with estimated
505 dimension of ~0.8–1.8 km.

506 (3) Although the resonance effect with the atmospheric acoustic-gravity waves could be a source of
507 increased wave energy, its direct contribution to the long-lasting oscillation is not demonstrated yet.
508 However, the comparison of hydrodynamical characteristics between the 2022 HTHH tsunami event and
509 the 2011 Tohoku tsunami event well demonstrated that the interactions between the ocean waves left by
510 atmospheric waves and local bathymetry contribute to the long-lasting Pacific oscillation of the 2022
511 tsunami event.

512 (4) The extraordinary features of this rare volcanic tsunami event challenge the current tsunami warning
513 system which is mainly designed for seismogenic tsunamis. It is necessary to improve the awareness of
514 people at risks about the potential tsunami hazards associated with volcanic eruptions. New approaches
515 are expected to be developed for tsunami hazard assessments with these unusual sources: various
516 atmospheric waves radiated by volcanic eruptions besides those traditionally recognized, e.g.
517 earthquakes, landslides, caldera collapses and pyroclastic flows etc.

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524 The JAGURS tsunami simulation code is employed for tsunami modelling (Baba et al., 2015;
525 <https://zenodo.org/record/6118212#.Yk98qdtBxPY>). Bathymetry data are obtained from GEBCO at
526 <http://www.gebco.net>. The sea level records in deep ocean are available from the Deep Ocean Assessment

527 and Reporting of Tsunamis (DART) buoy network in the Pacific (<https://nctr.pmel.noaa.gov/Dart/>), and
528 GeoNet New Zealand DART network (<https://tilde.geonet.org.nz>). The sea level records of tide gauges
529 are downloaded from UNESCO/ IOC (<http://www.ioc-sealevelmonitoring.org/>). Barometer data are
530 provided by the following providers: Direcció'n Meteoroló'gica de Chile
531 (<https://climatologia.meteochile.gob.cl>), NOAA National Weather Service
532 (<https://www.weather.gov/ilm/observations>), Japan Meteorological Agency (<https://www.jma.go.jp>),
533 The UK Met Office Weather Observation (<https://wow.metoffice.gov.uk/observations>), and Fiji
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