

1 **The characteristics of the 2022 Tonga volcanic tsunami in the Pacific Ocean**

2 Gui Hu¹, Linlin Li^{1,2}, Zhiyuan Ren³, Kan Zhang¹

3 1. Guangdong Provincial Key Laboratory of Geodynamics and Geohazards, School of Earth Sciences
4 and Engineering, Sun Yat-sen University, Guangzhou, China

5 2. Southern Marine Science and Engineering Guangdong Laboratory (Zhuhai), Zhuhai, China

6 3. Department of Civil and Environmental Engineering, National University of Singapore, Singapore.

7 *Correspondence to:* Linlin Li (lilinlin3@mail.sysu.edu.cn)

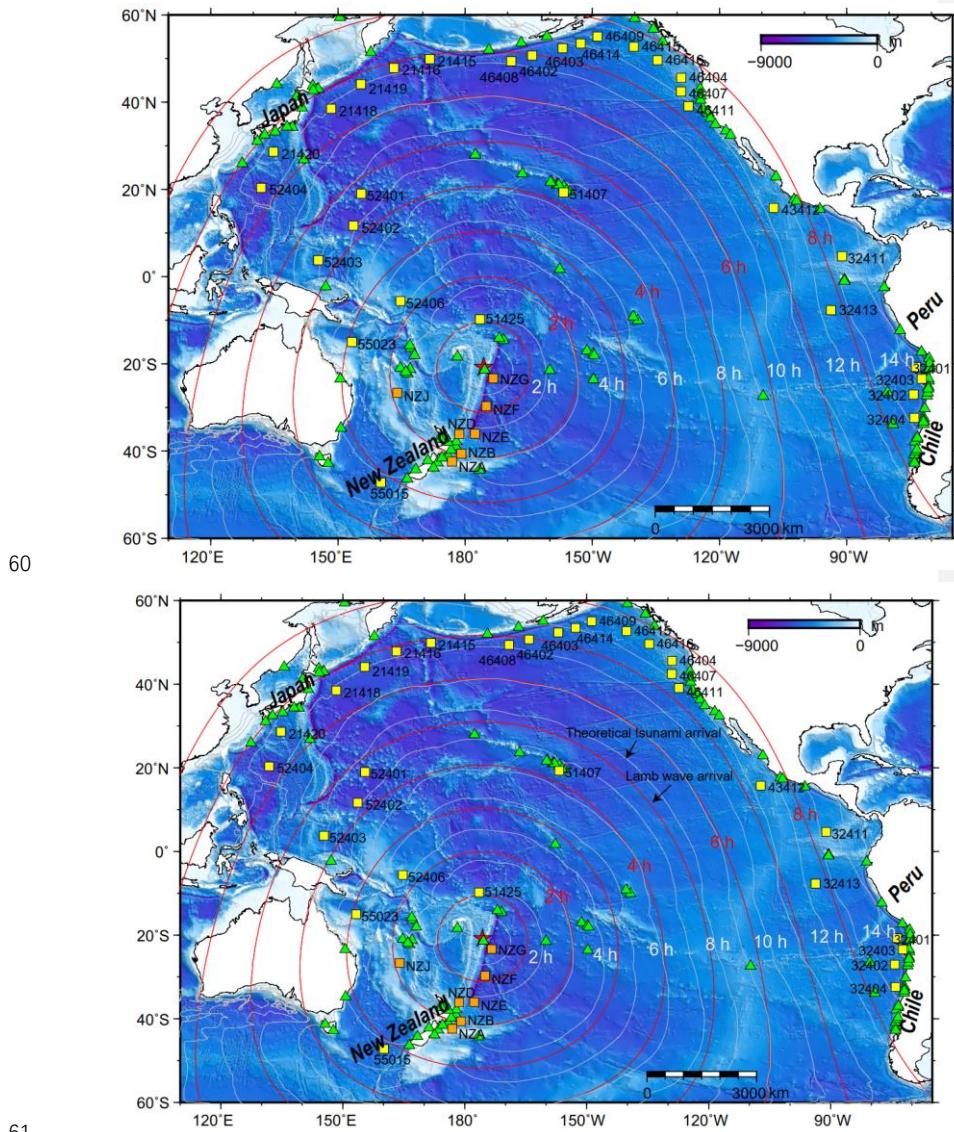
8 **Abstract.** On 15th January 2022, an exceptional eruption of Hunga Tonga–Hunga Ha’apai volcano
9 generated atmospheric and tsunami waves that were widely observed at oceans globally, gaining a
10 remarkable attention to scientists in related fields. The tsunamigenic mechanism of this rare event
11 remains an enigmatic due to its complexity and lacking of direct underwater observations. Here, to
12 explore the tsunamigenic mechanisms of this volcanic tsunami event and its hydrodynamic processes in
13 the Pacific Ocean, we conduct tsunami waveform and spectral analyses of the waveform recordings at
14 116 coastal gauges and 38 deep-ocean buoys across the Pacific Ocean. Combined with the constraints of
15 some representative barometers, we obtain the plausible tsunamigenic origins during the volcano activity.
16 We identify four distinct tsunami wave components generated by air-sea coupling and seafloor crustal
17 deformation. Those tsunami components are differentiated by their different propagating speeds or period
18 bands. The first-arriving tsunami component with ~80–100 min period was from shock waves spreading
19 at a velocity of ~1000 m/s in the vicinity of the eruption. The second component with extraordinary
20 tsunami amplitude in deep sea was from Lamb waves. The Lamb wave with ~30–40 min period radically
21 propagated outward from the eruption site with spatially decreasing propagation velocities from ~340
22 m/s to ~315 m/s. The third component with ~10–30 min period was probably from some atmospheric
23 gravity wave modes propagating faster than 200 m/s but slower than Lamb waves. The last component
24 with ~3–5 min period originated from partial caldera collapse with dimension of ~0.8–1.8 km.
25 Surprisingly, the 2022 Tonga volcanic tsunami produced long oscillation in the Pacific Ocean which is
26 comparable with those of the 2011 Tohoku tsunami. We point out that the long oscillation is not only
27 associated with the resonance effect with the atmospheric acoustic-gravity waves, but more importantly
28 the interactions with local bathymetry. This rare event also calls for more attention to the tsunami hazards
29 produced by atypical tsunamigenic source, e.g., volcanic eruption.

30 **1. Introduction**

31 On 15 January 2022 at 04:14:45 (UTC), a submarine volcano erupted violently at the uninhabited Hunga
32 Tonga-Hunga Ha'apai (HTHH) island at 20.546°S 175.390°W (USGS, 2022). The volcano is located ~67
33 km north of Nuku'alofa, the capital of Tonga (NASA, 2022) (Figure 1). The blasts launched plumes of
34 ash, steam, and gas ~58 km high into stratosphere (Yuen et al., 2022) which not only blanketed nearby
35 islands in ash (Duncombe, 2022; NASA, 2022), but caused various atmospheric acoustic-gravity wave
36 modes (AGWs) of various scales, e.g., Lamb waves from atmospheric surface pressure disturbance
37 associated with the eruption (Liu and Higuera, 2022; Adam, 2022; Kubota et al., 2022; Matoza et al.,
38 2022). Tsunami with conspicuous sea level changes were detected by coastal tide gauges and Deep-ocean
39 Assessment and Reporting of Tsunamis (DART) buoy stations in the Pacific (Figure 1), the Atlantic, and
40 Indian Oceans as well as the Caribbean and Mediterranean seas (Carvajal et al., 2022; Kubota et al., 2022;
41 Ramírez-Herrera et al., 2022), while the large waves were mainly concentrated in the Pacific Ocean, like
42 coastlines of New Zealand, Japan, California, and Chile (Carvajal et al., 2022). The event caused at least
43 3 fatalities in Tonga. Two people drowned in northern Peru when ~2 m destructive tsunami waves
44 inundated an island in the Lambayeque region, Chile (Edmonds, 2022).

45 Satellite images revealed that the elevation of HTHH island has gone through dramatic change before
46 and after the mid-January 2022 eruption. Previously, after the 2015 eruption, the two existing Hunga
47 Tonga and Hunga Ha'apai Islands were linked together. The volcanic island rose 1.8 km from the seafloor
48 where it stretched ~20 km across and topped a underwater caldera ~5 km in diameter (Garvin et al., 2018;
49 NASA, 2022). After the violent explosion on 15 January 2022, the newly formed island during 2015 was
50 completely gone, with only small tips left in far southwestern and northeastern HTHH island (NASA,
51 2022). HTHH volcano lies along the northern part of Tonga-Kermadec arc, where the Pacific Plate
52 subducts under the Indo-Australian Plate (Billen et al., 2003). The convergence rate (15~24 cm/year)
53 between the Tonga-Kermadec subduction system and the Pacific plate is among the fastest recorded plate
54 velocity on Earth, forming the second deepest trench around the globe (Satake, 2010; Bevis et al., 1995).
55 The fast convergence rate contributes to the frequent earthquakes, tsunamis and volcanic eruptions in
56 this region historically (Bevis et al., 1995). The 2022 HTHH volcano is part of a submarine-volcano
57 chain that extends all the way from New Zealand to Fiji (Plank et al., 2020). HTHH volcano had many
58 notable eruptions before 2022 since its first historically recorded eruption in 1912, i.e., in 1937, 1988,

59 2009, 2014-2015 (Global Volcanism Program, <https://volcano.si.edu>).



60
61
62 **Figure 1.** The spatial distribution of the eruption site (red star), DART stations (squares), tide
63 gauges (triangles) and the calculated tsunami arrival times. White contours indicate the modelled
64 arrival times of conventional tsunami. Red contours indicate the estimated arrival times of Lamb
65 waves (see how we derive these contours in section 3.1).

66 The 2022 HTHH eruption is the first volcanic event which generates worldwide tsunami signatures since
67 the 1883 Krakatau event (Matoza et al., 2022; Self and Rampino, 1981; Nomanbhoy and Satake, 1995).

68 The tsunamigenic mechanism of this rare volcanic eruption-induced tsunami is still poorly understood
69 due to its complex nature and the deficiencies of near-field seafloor surveys. Various tsunami generation
70 mechanisms have been proposed so far based on the observations of ground-based and spaceborne
71 geophysical instrumentations (Kubota et al., 2022; Matoza et al., 2022; Carvajal et al., 2022). The
72 mechanisms are closely associated with the air-sea coupling with atmospheric waves. Atmospheric
73 waves propagating in the atmospheric fluid with different speeds are generated by different physical
74 mechanisms (E. E. Gossard and W. H. Hooke, 1975). Nonlinearities in the process may lead to the
75 formation of shock-wave and period lengthening. The balance between gravity and buoyancy causes
76 gravity waves. The acoustic wave propagate by atmospheric fluid compression and rarefaction. (Matoza
77 et al., 2022). The most-mentioned mechanism of the tsunami is the fast-traveling atmospheric Lamb wave
78 generated by the atmospheric pressure rise of ~ 2 hPa during the eruption. The Lamb wave circled the
79 Earth for several times with travelling speed close to that of the sound wave in the lower atmosphere,
80 leading to globally observed sea level fluctuations (Adam, 2022; Duncombe, 2022; Kubota et al., 2022;
81 Matoza et al., 2022) (Figure 1). The second mechanism is suggested to be a variety of other acoustic-
82 gravity wave modes (Adam, 2022; Matoza et al., 2022; Themens et al., 2022; Zhang et al., 2022). The
83 third mechanism may be related to the seafloor crustal deformation induced by one or more volcanic
84 activities in the vicinity of the eruption site (e.g., pyroclastic flows, partial collapse of the caldera)
85 (Carvajal et al., 2022), which are more responsible for the near-field tsunamis with theoretical tsunami
86 speeds.

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87 To investigate the possible tsunamigenic mechanisms and detailed hydrodynamic behaviors of this rare
88 volcanic tsunami event, in this study, we collect, process and analyze the sea level measurements from
89 116 tide gauge and 38 DART buoys in the Pacific Ocean (shown in Figures 1 and 2). We first do statistical
90 analysis of the tsunami waveforms to estimate the propagating speed of the Lamb wave and to understand
91 the tsunami wave characteristics in the Pacific Ocean through demonstrating the tsunami wave properties,
92 i.e., arrival times, wave heights and durations. We then conduct wavelet analysis for representative DART
93 buoys and tide gauges respectively to explore tsunamigenic mechanisms of the event and to better
94 understand its hydrodynamic processes in the Pacific Ocean. Aided by wavelet analysis of corresponding
95 barometers near the selected DART buoys and comparison with tsunami records of the 2011 Tohoku
96 tsunami, we are able to piece together all the analysis and demonstrate that the 2022 HTHH tsunami was

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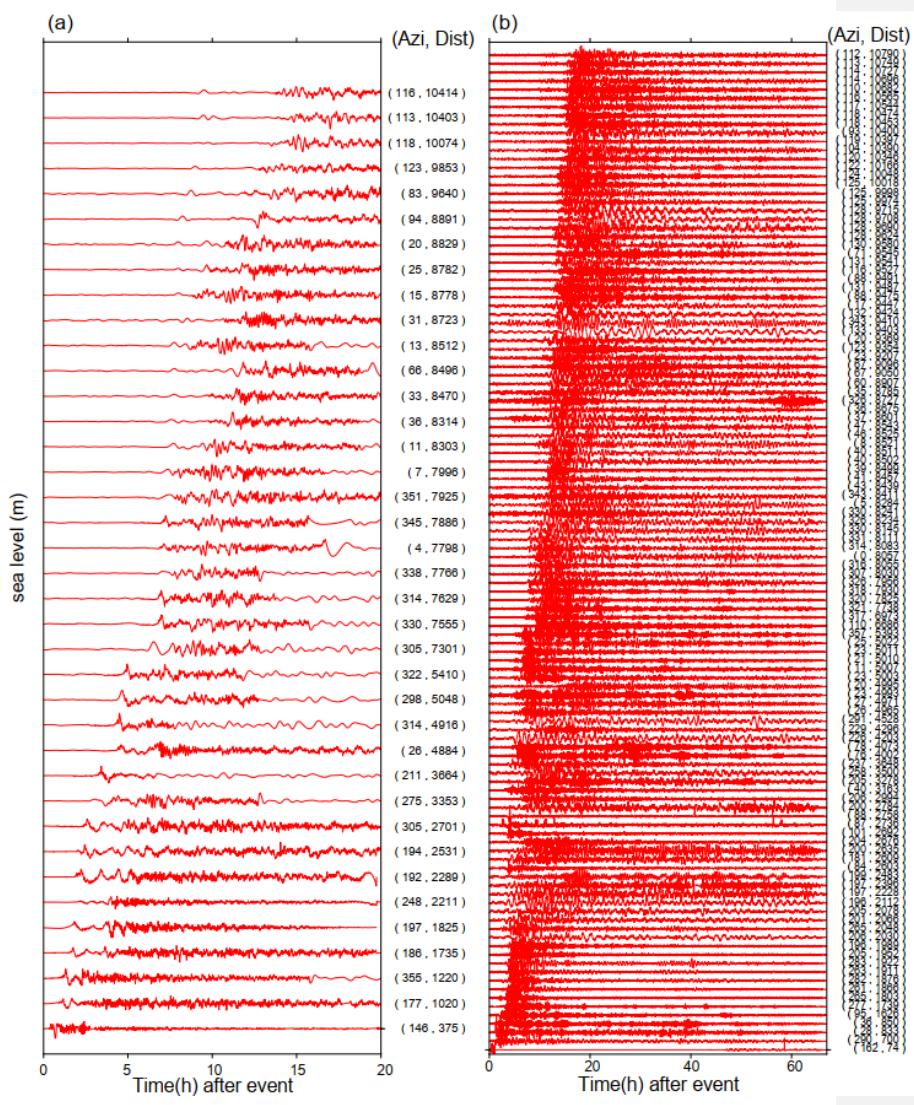
97 generated by air-sea coupling with a wide range of atmospheric waves with different propagating
98 velocities and period bands, and seafloor crustal deformation associated with the volcanic eruption. We
99 demonstrate as well that the tsunami was amplified at the far-field Pacific coastlines where the local
100 bathymetric effects play a dominant role in tsunami scale.

101 **2. Data and Methods**

102 **2.1 Data**

103 We collected high-quality sea level records across the Pacific Ocean at 38 DART buoys (in which 31
104 stations from <https://nctr.pmel.noaa.gov/Dart/>, 7 stations from <https://tilde.geonet.org.nz/dashboard/>) and
105 116 tide gages from IOC (The Intergovernmental Oceanographic Commission, <http://www.ioc-sealevelmonitoring.org>) (Figure 1). The epicentral distances of tide gauges and DART buoys range
106 between 74–10790 km and 375–10414 km, respectively. The sampling rates of DART buoys are
107 changing over time. Passing of tsunami event generally can trigger the DART system to enter its high
108 frequency sampling mode (15 seconds or 1 min) from normal frequency mode (15 min)
109 (www.ndbc.noaa.gov/dart). In contrast, sampling rates of normal tide gauges at coasts are uniform with
110 sampling interval of 1 min. The sampling interval of both DART and tide gauges is preprocessed to 15
111 seconds. Firstly, we eliminate abnormal spikes and fill gaps by linear interpolation. Secondly, we applied
112 a fourth-order Butterworth-Highpass filter with a cut-off frequency of 3.5 e-5 Hz (~ 8 hours) to remove
113 the tidal components (Figure 2) (Heidarzadeh and Satake, 2013). After the two steps, quality control step
114 is conducted to select high-quality data, in which we delete waveforms with spoiled data or massive data
115 loss due to equipment failure, or with the maximum tsunami heights of tide gauges less than 0.2 m, then
116 the selected data will be ready for further statistics and spectral analysis. We also collect and analyze the
117 atmospheric pressure disturbance data recorded by some representative barometers. The sampling rates
118 of the barometers is generally uniform with a sampling rate of 1 min except for some stations in New
119 Zealand with interval of 10 min. Considering the sample rate, we employ a fourth-order Butterworth-
120 Bandpass filter with period ranging between 2–150 min for wavelet analysis of the barometers with 1
121 min sample rate, while we apply the fourth-order Butterworth-Bandpass filter with range of 30–150 min
122 to long-period waveform display based on two reasons. (1) The barometer data we use for the analysis
123 include some in New Zealand with 10 min sample rate; (2) Filtering out the short-period waves helps

125 highlight long-period tsunami wave components.
126 The tsunami waveforms recorded by DART buoys which are installed offshore in the deep water are
127 expected to contain certain characteristics of the tsunami source (Wang et al., 2020, 2021). The
128 waveforms recorded by tide gauge distributed along coastlines are significantly influenced by local
129 bathymetry/topography which are used for investigating bathymetric effect on tsunami behaviors
130 (Rabinovich et al., 2017, 2006; Rabinovich, 2009). Therefore, we use the DART data for source-related
131 analysis and choose some tide gauge data to investigate the tsunami behaviors at the Pacific coastlines.



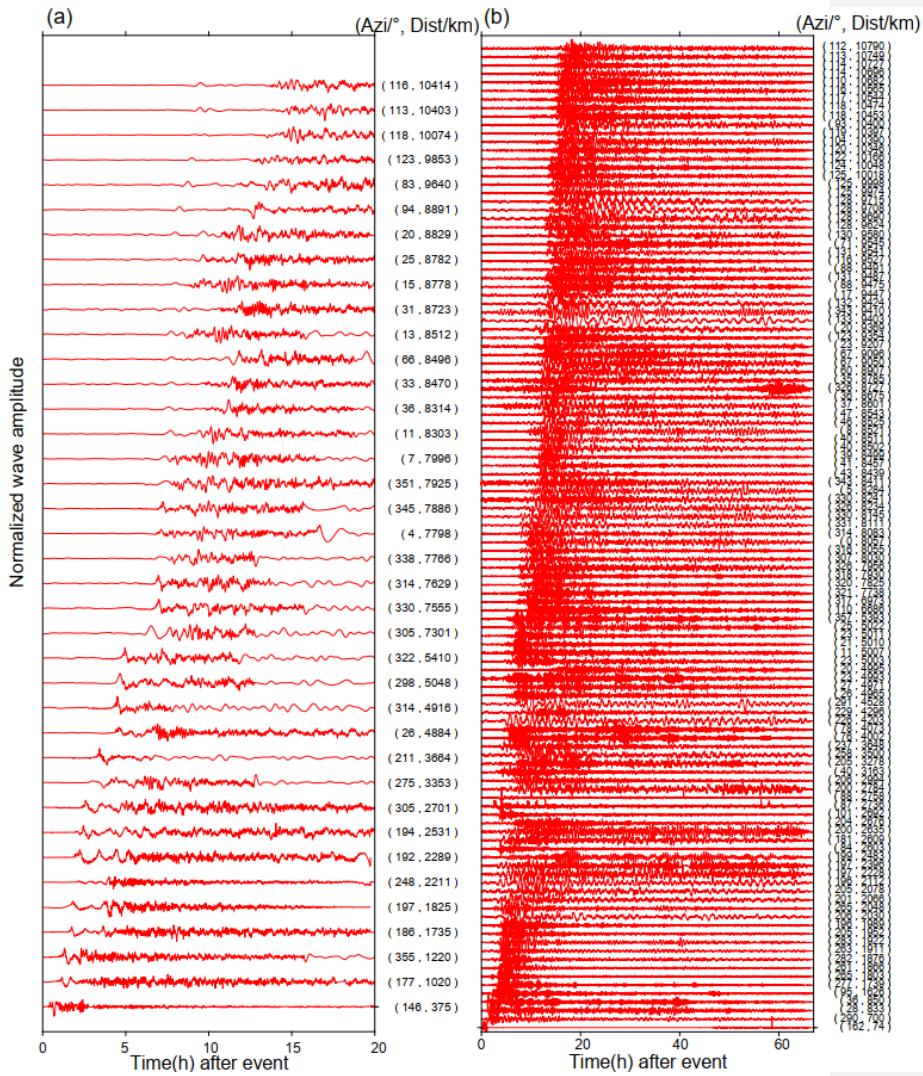


Figure 2. Detidied tsunami waveforms at (a) DART buoys and (b) tide gauges. Waveforms in both subplots are shown in ascending distance. Azi stands for azimuth.

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136 2.2 Tsunami Modelling

137 We use a numerical tsunami modelling package JAGURS (Baba et al. 2015) to simulate the tsunami
138 propagation of the 2022 HTHH event and obtain the theoretical tsunami arrival time based on the shallow
139 water wave speed (white contours in Figure 1). The code solves linear Boussinesq-type equations in a
140 spherical coordinate system using a finite difference approximation with the leapfrog method. We specify

141 a unit Gaussian-shaped vertical sea surface displacement at the volcanic base as the source of
142 conventional tsunami. For a unite source i with center at longitude φ_i and latitude θ_i , the
143 displacement distribution $Zi(\varphi, \theta)$ can be expressed as:

$$144 Zi(\varphi, \theta) = \exp\left[-\frac{(\varphi-\varphi_i)^2 + (\theta-\theta_i)^2}{2\sigma^2}\right] \quad (1)$$

145 Where we set characteristic length σ as 5 km (NASA, 2022). The bathymetric data is resampled from the
146 GEBCO 2019 with 15 arc-sec resolution (The General Bathymetric Chart of the Oceans, downloaded
147 from <https://www.gebco.net>).

148 2.3 Spectral Analysis of Tsunami Waves

149 To investigate the temporal changes of the dominant wave periods, we conduct continuous wavelet
150 transformation (frequency-time) analyses for some representative DART buoys, tide gauges and
151 barometers, in which wavelet Morlet mother function is implemented (Krsteková et al., 2006). The first
152 32-hour time series of DART buoys and barometers after the eruption (at 04:14:45 on 15 January 2022)
153 are used for source-related wavelet analysis. The first 48-hour time series of tide gauges after the eruption
154 are employed for hydrodynamics-related wavelet analysis at coastlines. We adopt the Averaged-Root-
155 Mean-Square (ARMS) method as a measure of absolute average tsunami amplitude with a moving time
156 window of 20 min to calculate the tsunami duration (Heidarzadeh and Satake, 2014). We define the time
157 durations as the time period where ARMS levels of tsunami waves are above those prior to the tsunami
158 arrivals.

159 3. Results

160 3.1 The decreasing propagation velocities of the Lamb Wave

161 Although many types of atmospheric waves were generated by the 2022 HTHH eruption, the most
162 prominent signature was the Lamb waves which were globally observed by ground-based and spaceborne
163 geophysical instrumentations (Kulichkov et al., 2022; Liu et al., 2022; Lin et al., 2022; Matoza et al.,
164 2022; Themens et al., 2022; Adam, 2022; Kubota et al., 2022). Interestingly, we notice that a wide range
165 of the velocities from 280 m/s to 340 m/s were proposed through observations and Lamb wave modelling
166 (e.g., Kubota et al., 2022; Lin et al., 2022; Matoza et al., 2022; Themens et al., 2022). The travelling
167 velocity of Lamb waves in real atmosphere is affected by temperature distributions, winds and dissipation

168 (Otsuka, 2022). To investigate whether the propagation speeds of the lamb wave change in space and
169 time, we analyze the waveforms recorded by the DART buoys in the Pacific Ocean. DART buoy with
170 pressure sensor deployed at the ocean's bottom records the sea level change that is transferred from
171 pressure records in Pascals, instead of direct water height. For the 2022 HTHH tsunami event, the
172 pressure fluctuation at DART buoy is a superposition of the pressure changes caused by tsunami and the
173 Lamb wave (Kubota et al., 2022). The Pacific DART buoys recorded the most discernible air-sea
174 coupling pulse in deep ocean with Lamb waves that arrived earlier than the theoretical tsunamis (Figure
175 1). The tsunami waveforms recorded by tide gauges did not clearly detect the tsunami signals associated
176 with the Lamb waves, therefore not sufficient for further analysis (Figure 2). Thus, we estimate the
177 speed of Lamb waves using the waveforms recorded by the Pacific DART buoys. The Lamb wave
178 arrivals are limited within arrival time range from possible velocities of 280–340 m/s. The time points at
179 which the tsunami amplitudes first exceed 1 e-4 m above sea level are defined as Lamb wave arrivals.
180 By carefully fitting the arrivals with different constant velocities, we illustrate the velocities of Lamb
181 wave were generally uniform, but slightly decrease with the increase of propagation distance (Figure 3).
182 The Lamb waves initially propagated radially at speed of ~340 m/s before slowing to ~325 m/s after
183 reaching ~3400 km, and further decreasing to ~315 m/s at 7400 km. In an isothermal troposphere
184 assumption, the phase velocity of the Lamb wave (C_L) can be estimated with the following equation
185 (Gossard and Hooke, 1975):

$$C_L = \sqrt{\frac{\gamma R T}{M}} \quad (2)$$

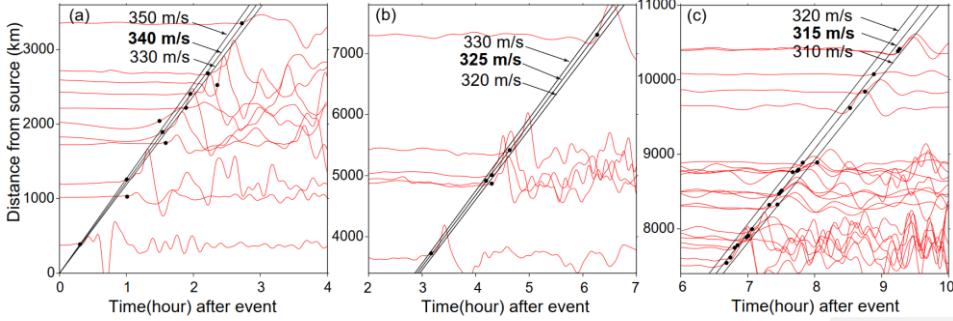
186 Where $\gamma = 1.4$ (air specific heat ratio corresponding to atmospheric temperature), $R = 8314.36 \text{ J kmol}^{-1}$
187 K^{-1} (the universal gas constant), $M = 28.966 \text{ kg kmol}^{-1}$ (molecular mass for dry air) are constant for the
188 air, T is the absolute temperature in kelvin. Thus, Lamb wave velocity is mainly affected by the air
189 temperature, meaning the travelling velocity of lamb waves might decrease when propagating from
190 regions with high temperature towards those with low temperatures, e.g., the north pole. By assuming a
191 set of possible temperatures in January (Table 1), we calculated the velocities C_L could range between
192 312–343 m/s when temperatures vary between -30–20 °C. Therefore, the decreased velocity of the Lamb
193 waves could be a consequence of cooling of the air temperature.

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196 **Figure 3. Fitting the arrival times of normalized Lamb waveforms with different velocities. Black**
197 **dots mark the arrival times of the Lamb waves. Black lines represent velocities.**

198 **Table 1. Estimated Lamb wave velocities in an isothermal troposphere assumption**

Celsius temperature (°C)	thermodynamic temperature (K)	C_L (m/s)
20	293.15	343.14
10	283.15	337.23
0	273.15	331.21
-10	263.15	325.19
-20	253.15	318.86
-30	243.15	312.49

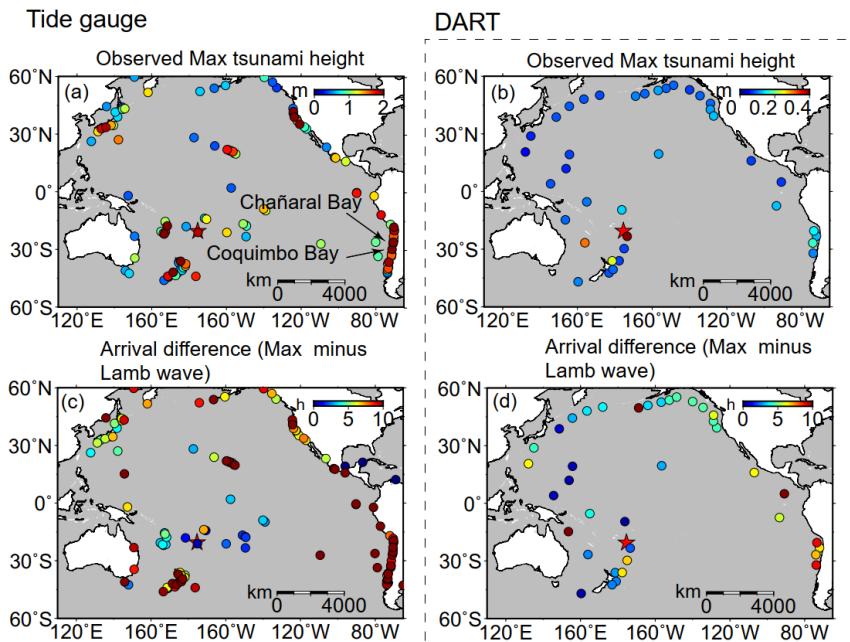
199 **3.2 Tsunami features observed by DART buoys and Tide gauges**

200 The statistics of tsunami heights and arrival times recorded at 38 DART buoys and 116 tide gauges across
201 the Pacific Ocean are used to interpret the tsunami characteristics. The comparison of the statistical
202 characters between DART and tide gauge observations yields some useful information of the
203 hydrodynamic process of tsunami propagation and help identify tsunami wave components with different
204 traveling velocities.

205 The average value of the maximum tsunami wave height (trough-to-crest) for the 116 tide gauge stations
206 is ~ 1.2 m. Figure 4a shows tide gauges with large tsunami heights exceeding 2 m are mainly distributed
207 in coastlines with complex geometries (Figure S1a), such as gauges at New Zealand, Japan, and north
208 and south America. For example, the largest tsunami height among tide gauges is 3.6 m at a bay-shaped
209 coastal area Chañaral in Chile (Figure S1b). In sharp contrast to tide gauges, the maximum tsunami
210 heights of most Pacific DART buoys are less than 0.2 m. The largest tsunami height in the DART buoys
211 is only ~ 0.4 m recorded at the nearest one, 375 km from the volcano (Figure 4b). The comparison between
212 DART buoys and tide gauges indicate that the direct contribution of air-sea coupling to the tsunami

213 heights is probably in the level of tens of centimeters (Kubota et al., 2022). The meter-scale tsunami
214 heights at the coastlines suggest the bathymetric effect could play a major role during tsunami
215 propagation. In respect to the arrival of maximum tsunami waves, the time lags between Lamb waves
216 and the maximum heights of tide gauges mainly range between ~0–10 h (Figure 4c). The delayed times
217 of ~10 h are observed in New Zealand, Hawaii, and west coast of America (Figure 4c), suggesting the
218 interaction between tsunami waves and local topography/bathymetry delays the arrival of the maximum
219 waves (e.g., Hu et al., 2022). For example, the delayed maximum tsunami height can be attributed to the
220 edge waves (Satake et al., 2020) and resonance effect (Wang et al., 2021) from tsunami interplays with
221 bays/harbors, islands, and continental shelves of various sizes. The significant regional dependence of
222 the coastal tsunami heights and the time lags of the maximum tsunami waves can be attributed to the
223 complexity of local bathymetry, such as continental shelves with different slopes, and harbor/bay with
224 different shapes and sizes (Satake et al., 2020). On the other hand, since the DART records are less
225 influenced by bathymetric variation in space, the first waves in DART buoys are supposed to be the
226 maximum tsunami waves as observed in the 2011 Tohoku tsunami event (Heidarzadeh and Satake, 2013).
227 However, we observe the inconsistency between the arrivals of the Lamb waves and the maximum
228 tsunami heights (Figure 4d). The time lags of the maximum waves of DART buoys present a coarsely
229 increasing tendency with the increasing distance from the volcano, which indicates the contribution of
230 other tsunami generation mechanism propagating with a uniform but lower speed than Lamb wave.

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232 **Figure 4.** The spatiotemporal signatures of the 2022 HTHH tsunami across the Pacific Ocean. (a)
233 Observed the maximum tsunami height (trough-to-crest height) of tide gauges. (c) Arrival
234 differences between the maximum tsunami height of tide gauges and Lamb waves. (b) and (d) are
235 the same as (a) and (c) but for DART buoys.

236 **3.3 Tsunami components identified from wavelet analysis**

237 The statistical analysis of tsunami waveforms at tide gauges and DART buoys suggest the tsunami waves
238 likely contain several components with different source origins. To further identify these tsunami
239 components, we conduct wavelet analysis for tsunami waveforms recorded by representative DART
240 buoys and air pressure waveforms recorded by selected barometers. We demonstrate the analysis result
241 through the frequency-time (f-t) plot of wavelet which shows how energy and period vary at frequency
242 and time bands (Figure 5 and Figure 6). Tsunami components have clear signatures in all f-t plots as the
243 energy levels are quite large when they arrive. Figure 5 shows the wavelet analysis of six DART buoys
244 located in the vicinity of the eruption site (<3664 km). Figure 6 show the wavelet analysis of ten DART
245 buoys located in the Pacific rim which are far away from the source location. We observe three interesting
246 phenomena: 1) most of the tsunami wave energy is concentrated in four major period bands, i.e., 3-5
247 min, ~10-30 min, ~30-40 min, and ~80-100 min; 2) min, ~80-100 min, 10-30 min, 30-40 min, and 3-5 min; 2)

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248 The significant tsunami component with period band of 3-5 mins are recorded by stations between the
249 eruption site and the north tip of the New ZealandThe stations with 3-5 mins wave periods are mainly
250 located in the vicinity of the volcano site; 3) There exists one exceptional tsunami component with longer
251 wave period of ~80–100 min, mainly recorded in the Tonga, the New Zealand and Hawaii, in the near
252 source region which travels even faster than the Lamb Lamb waves.

253 To further explore the source mechanism of these tsunami components, we take advantage of the
254 published information related to different propagating velocities of atmospheric gravity waves (Kubota
255 et al., 2022) and add four kinds of propagating velocities as criteria to differentiate the tsunami arrivals
256 from different sources (Figure 5 and Figure 6). The first reference speed is 1000 m/s related to the
257 radically propagating atmospheric shock waves near the source region (Matoza et al., 2022; Themens et
258 al., 2022). The second one is the velocities of Lamb wave ranging between 315–340 m/s derived from
259 the aforementioned analysis in section 3.1 (Figure 3). The third one is 200 m/s corresponding to the lower
260 limit of atmospheric gravity wave modes other than Lamb waves which were also excited by the volcanic
261 eruptions (Kubota et al., 2022). The last is the arrival time of conventional tsunami given by tsunami
262 modelling (Figure 1). The theoretical velocity of conventional tsunami is significantly nonuniform
263 spatially as compared with those of the atmospheric waves. The conventional tsunami propagation speed
264 is determined by the water depth along the propagation route. The velocity of non-dispersion shallow-
265 water waves (C_H) in the ocean is given by:

$$C_H = \sqrt{g \cdot H} \quad (3)$$

266 Where g is gravity acceleration (9.81 m/s^2), H is the water depth. The propagation velocities of tsunami
267 are ~296–328 m/s in the deepest trenches on earth (i.e., ~11 km in Mariana Trench and ~9 km in Tonga
268 Trench). The velocities decrease quickly to only ~44 m/s at ~200 m depth along the edge of continental
269 shelf. With the average depth of ~4–5 km, the average velocities in the Pacific Ocean range between
270 ~198–221 m/s. Thus, theoretical tsunami velocities present significant slowness and variability. We
271 delineate the arrival times of the four reference speeds in Figures 5 and 6. For each panel of the figures,
272 from left to right, the solid vertical white lines mark velocity of 1000 m/s. The solid vertical red lines
273 mark the arrival of Lamb waves. The dashed vertical white lines mark lower limit of gravity waves'
274 velocity of 200 m/s. The dashed vertical black lines represent the calculated theoretical tsunami arrivals.
275 Horizontal white dashed lines mark two reference periods of 10 min and 30 min.

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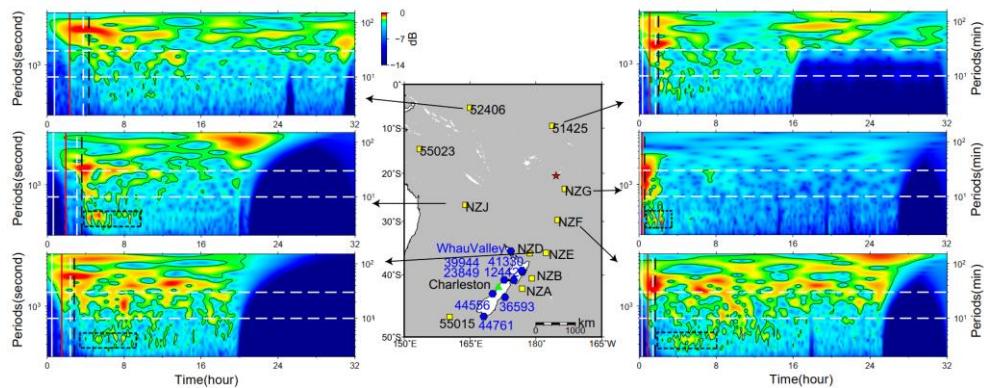
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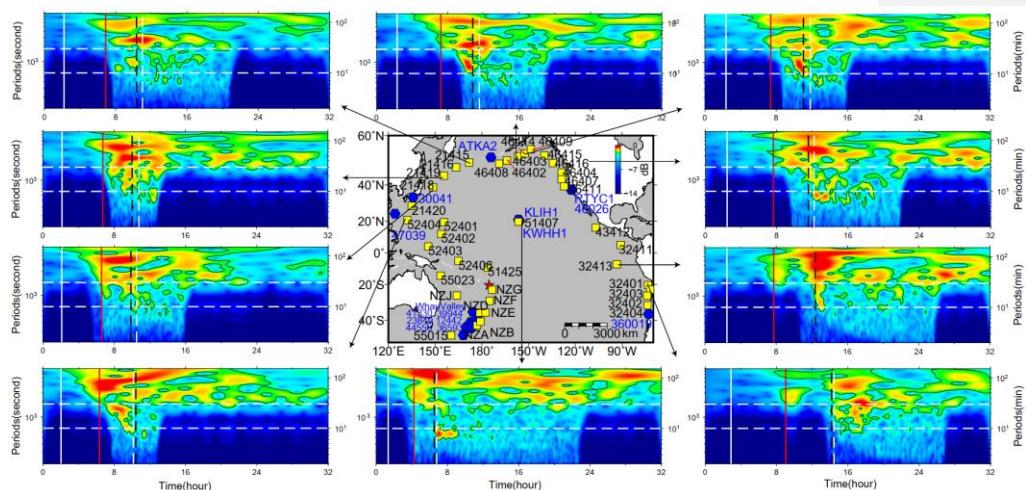
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277 One particularly remarkable phenomenon is that the wave component with period of ~80–100 min
278 propagated at a very fast speed of ~1000 m/s in the vicinity of the HTHH site, i.e., New Zealand and
279 Hawaii (e.g., stations 52406, NZJ, NZE, 51425 in Figure 5, and 51407 in Fig. 6). We infer that the
280 tsunami component within ~80–100 min period band was likely produced by the atmospheric shock
281 waves during the initial stage of the volcanic eruption and spatially only cover the near-source region.
282 To verify this observation, we select 16 representative barometers located in the near-source region and
283 far-field area for wavelet analysis (see the locations in Figure 5 and Figure 6). Figure 7 shows the
284 waveforms of atmospheric pressure at selected locations and Figure 8 provides the frequency-time (f-t)
285 plot of wavelet analysis of some representative barometers. Interestingly, we are able to discern the air
286 pressure pulses prior to Lamb waves at barometers in New Zealand (the two columns on the left in Figure
287 7), although such signals are not detectable in waveforms recorded by barometers far from the source
288 (the two columns on the right in Figure 7). The spatial distribution of such unusual pressure changes
289 suggest that the fast travelling shock waves were only limited in the near-source region, as reflected in
290 the travelling ionospheric disturbances (Matoza et al., 2022; Themens et al., 2022). Additionally, we also
291 see that the long period signals of ~80–100 min appear in DART buoys far away from the eruption site.
292 Such signals may be related with the long-period gravity waves (Matoza et al., 2022).
293 The tsunami components at period band of ~30–40 min can be readily associated with Lamb waves
294 because the arrival times of tsunami waves and Lamb waves have excellent match, as shown in the
295 tsunami data recorded by DART buoys (e.g., NZJ and 51425 in Figure 5; 51407, 32401 and 32413 in
296 Figure 6) and pressure data by barometers (Figure 8).
297 For the tsunami components with the period band of ~10–30 min, although the arrivals of ~10–30 min
298 tsunami components cover some theoretical tsunami arrival times, they do not consistently match. The
299 tsunami components occurring within the time period between Lamb waves and the lower gravity waves'
300 velocities has a good agreement with the velocity range of several atmospheric gravity wave modes
301 (Matoza et al., 2022; Themens et al., 2022; Kubota et al., 2022). Similarly, the air pressure data also show
302 energy peaks at ~10–30 min period band, which is consistent with the tsunami data (Figure 8). Such
303 consistency further verifies the contribution of atmospheric gravity waves to the volcanic tsunami.
304 The tsunami components with the shortest period of ~3–5 min (stations NZE, NZF, NZG and NZJ;
305 marked with black dashed squares in Figure 5) are only observed at DART records near the eruption

306 location. Meanwhile, the arrival times of these components agree well with the modelled arrivals of
 307 conventional tsunami. Thus, we believe the observed shortest period band should originate from the
 308 seafloor crustal deformation. We further infer that this component could be generated by the partial
 309 underwater caldera collapse and/or subaerial/submarine landslide failures associated with 2022 HTHH
 310 volcanic eruption.

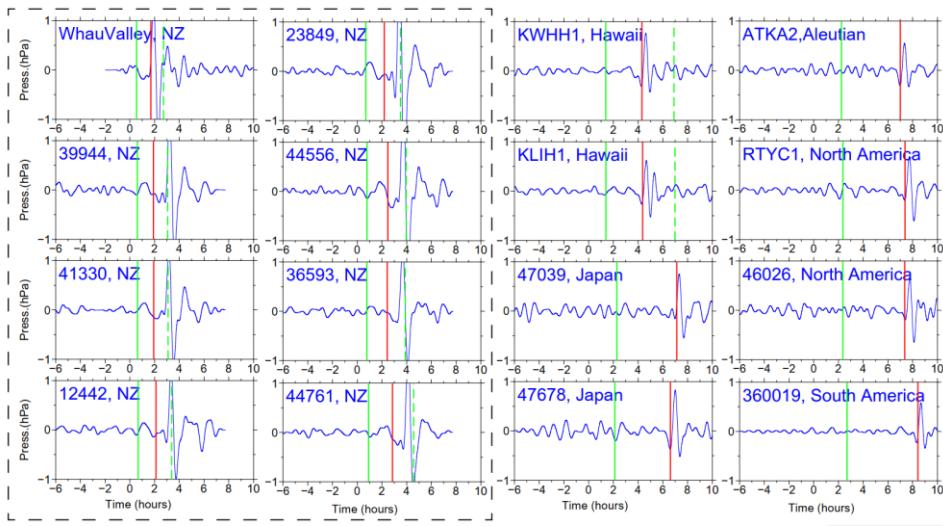


311
 312 **Figure 5. Wavelet analysis of representative DART buoys in the vicinity of the HTHH volcano. In**
 313 **each sub-plot, the solid vertical white lines mark the arrival time with travelling velocity of 1000**
 314 **m/s. The solid vertical red lines mark the arrivals of Lamb waves. The dashed vertical white lines**
 315 **mark lower limit of AGWs' velocity of 200 m/s (Kubota et al., 2022). The dashed vertical black**
 316 **lines represent the theoretical tsunami arrivals. The dashed horizontal white lines mark two**
 317 **reference wave periods of 10 min and 30 min. The blue hexagons represent the locations of**
 318 **barometers. Green triangle makes the location of the tide gauges at Charleston.**

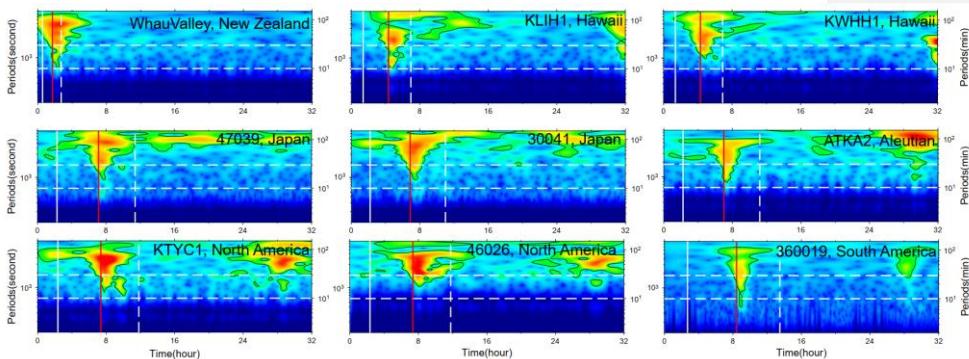


319

320 **Figure 6.** Wavelet analysis of representative DART buoys far away from the HTHH volcano. In
 321 each sub-plot, the solid vertical white lines mark the arrival time with travelling velocity of 1000
 322 m/s. The solid vertical red lines mark the arrivals of Lamb waves. The dashed vertical white lines
 323 mark lower limit of AGWs' velocity of 200 m/s. The dashed vertical black lines represent the
 324 theoretical tsunami arrivals. The dashed horizontal white lines mark two reference wave periods
 325 of 10 min and 30 min. The blue hexagons represent the locations of barometers.



326 **Figure 7.** Shockwave-related atmospheric pressure waveforms of selected barometers in the Pacific
 327 Ocean. All traces have been filtered between 30 min and 150 min. In each sub-plot, the solid vertical
 328 green lines mark the arrival time with travelling velocity of 1000 m/s. The solid vertical red lines
 329 mark the arrivals of Lamb waves. The dashed vertical green lines mark lower limit of AGWs'
 330 velocity of 200 m/s.
 331



332 **Figure 8.** Wavelet analysis of some representative barometers. In each sub-plot, the solid vertical
 333 white lines mark the arrival time with travelling velocity of 1000 m/s. The solid vertical red lines
 334 mark the arrivals of Lamb waves. The dashed vertical white lines mark lower limit of AGWs'
 335

336 **velocity 200 m/s. The dashed horizontal white lines mark three reference periods of 10 min and 30**
337 **min.**

338 **4. Discussion**

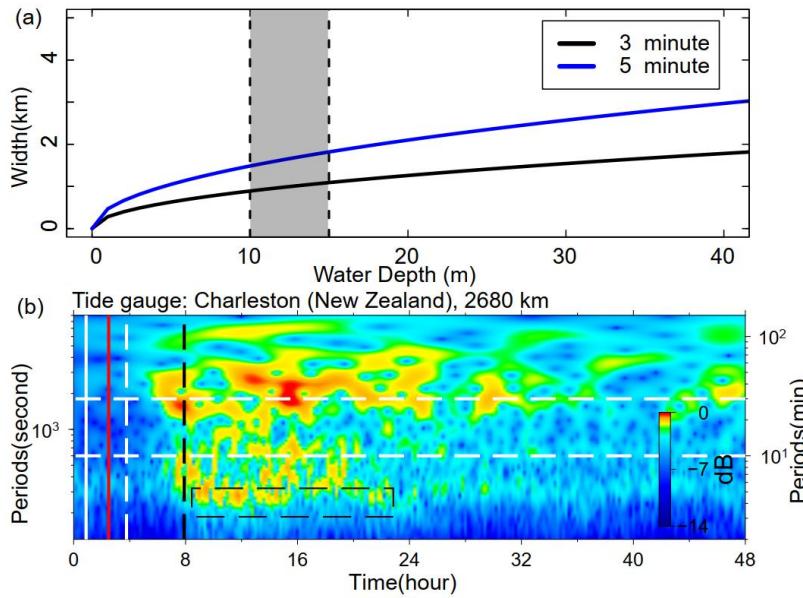
339 **4.1 Tsunami from Caldera Collapse and Its Long-distance Traveling Capability**

340 The tsunami wave energy distributed in different period bands is identified with reference arrival times.
341 The tsunami component with 3–5 min period is most likely generated by seafloor crustal deformation in
342 the volcanic site, but specific mechanism is not determined. A variety of possible scenarios associated
343 with the eruption could be responsible for the near-field tsunami waves, such as volcanic earthquakes,
344 pyroclastic flows entering the sea, underwater caldera flank collapse, and subaerial/submarine failures
345 (Self and Rampino, 1981; Pelinovsky et al., 2005). To further investigate the source mechanism, we
346 apply a simplified model (Rabinovich, 1997) to estimate the probable dimension of tsunami source:

$$347 L = \frac{T\sqrt{gH}}{2} \quad (4)$$

348 Where L is the typical dimension (length or width) of the tsunami source, H is average water depth in the
349 source area, g is the gravity acceleration, and T is primary tsunami period. By comparing with the post-
350 2015 morphology of the HTHH caldera which was obtained through drone photogrammetry and
351 multibeam sounder surveys, Stern et al. (2022) estimate that much of the newly-formed Hunga Tonga
352 Island and the 2014/2015 cone were destroyed by the 2022 eruption, and the vertical deformation of
353 Hunga Ha'apai Island is ~10–15 m (Stern et al., 2022). With no more quantitative constraint of the
354 seafloor deformation, we tentatively assume H as 10–15 m, then the possible dimension of seafloor
355 crustal deformation responsible for the small-scale tsunami could be in the scale of 0.8–1.8 km (Figure
356 9a). The estimated size is very likely from partial caldera collapse that usually has limited scale in
357 volcanic site (Ramalho et al., 2015; Omira et al., 2022). If it is the case, the partial flank collapse could
358 be located between Hunga Tonga and Hunga Ha'apai Islands.

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359
360 **Figure 9. Mechanism of tsunami component with 3–5 min period.** (a) The source dimension
361 estimated by equation 4. (b) Wavelet analysis of tide gauge at Charleston, New Zealand, 2680 km
362 away from the eruption site. The solid vertical white line marks the arrival time with travelling
363 velocity of 1000 m/s. The solid vertical red line marks the arrival of Lamb wave. The dashed
364 vertical white line marks lower limit of AGWs' velocity 200 m/s. The dashed vertical black line
365 marks the theoretical tsunami arrivals.

366 An interesting phenomenon is that the tsunami component with 3–5 min period can still be observed in
367 a bay-shaped coastal area at Charleston in New Zealand (see the location in Figure 5) which is 2680 km
368 away from the eruption site and maintains a high energy level lasting up to 14 h (Figure 9b). The long-
369 traveling capability could be associated with the ~10000 m deep water depth of the Tonga Trench that
370 keeps the source signals from substantial attenuation. In deep open ocean, the wavelength of a tsunami
371 can reach two hundred kilometers, but the height of the tsunami may be only a few centimeters. Tsunami
372 waves in the deep ocean can travel thousands of kilometers at high speeds, meanwhile losing very little
373 energy in the process. The long oscillation can be attributed to the multiple reflections of the incoming
374 waves trapped in the shallow-water bay at Charleston.

375 Generally, devastating tsunamis with long-distance travelling capability are mostly generated by
376 megathrust earthquakes (Titov et al., 2005). Caldera collapses or submarine landslides with limited scale
377 normally only generate local tsunamis, e.g., the 1998 PNG (Papua New Guinea) tsunami event (Kawata

378 et al., 1999) and the 1930 Cabo Girão tsunami event (Ramalho et al., 2015). Therefore, it's exceptional
379 that the tsunami component from scale-limited failure could travel at-least 2680 km away from the
380 eruption site. It demonstrates that tsunamis from small-scale tsunamigenic source have the capability to
381 travel long distance and cause long oscillation at favored condition, e.g., deep trench, ocean ridge and
382 bay-shaped coasts.

383 **4.2 The Possible Mechanisms of Long Tsunami Oscillation**

384 An important tsunami behavior of the 2022 HTHH tsunami is the long-lasting oscillation ~ 3 days in the
385 Pacific Ocean (Figure 10a), which is comparable to that of the 2011 Tohoku tsunami, ~4 days
386 (Heidarzadeh and Satake, 2013). We demonstrate the duration time of the tsunami oscillation through
387 ARMS (Averaged-Root-Mean-Square) approach that is a measure of absolute average tsunami amplitude
388 in a time period. The long-lasting tsunami energy can be observed at many regions, such as the coasts of
389 New Zealand, Japan, Aleutian, Chile, Hawaii, and west coasts of America. Several mechanisms could
390 account for the long-lasting tsunami, including (1) Lamb waves circling the Earth multiple times
391 (Amores et al., 2022; Matoza et al., 2022), (2) resonance effect between ocean waves and atmospheric
392 waves (Kubota et al., 2022), and (3) bathymetric effect. We discuss the contribution of each mechanism
393 in the following section.

394 To investigate the contribution of Lamb wave to the long-lasting tsunami, we compare the air pressure
395 disturbances recorded by selected barometers together with the tsunami waveforms of nearby tide gauges
396 (Figure 10b). While the barometers present discernible wave pulses at each Lamb wave's arrival, only
397 the first Lamb wave triggered clear tsunami signal and no detectable tsunami signatures correspond to
398 the following passage, suggesting the Lamb waves do not directly contribute to the long oscillation.

399 Theoretically, the resonance effects between ocean waves and atmospheric waves could contribute to the
400 long oscillation on coastlines based on the following reasons. First, part of the atmospheric gravity waves
401 propagated at velocities close to averaged velocities of conventional tsunami in the Pacific Ocean (198–
402 221 m/s) which resulted in the resonance with ocean waves (Kubota et al., 2022). Second, when Lamb
403 wave speed approaches the tsunami speed, Proudman resonance gradually increase tsunami heights,
404 wherein Proudman resonance optimally maximizes tsunami heights when they match well. In deep
405 oceanic trenches, such as Mariana and Tonga-Kermadec trench (10000–11000 m), tsunami velocities
406 range between ~314–330m/s which are comparable with those of the observed Lamb waves 315–340

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407 ~~m/s. Second, in deep oceanic trenches, such as Mariana and Tonga-Kermadec trench (10000–11000 m),~~
408 ~~tsunami velocities range between ~314–330 m/s which are comparable with those of the observed Lamb~~
409 ~~waves 315–340 m/s. When Lamb wave speed approaches the tsunami speed, Proudman resonance~~
410 ~~gradually increase tsunami heights, wherein Proudman resonance optimally maximizes tsunami heights~~
411 ~~when they match well (Tanioka et al., 2022; Lynett et al., 2022). Therefore, the resonance effect, could~~
412 ~~be a possible source of increased wave energy continuously supplied wave energy to the ocean, especially~~
413 in the deep trenches (Lynett et al., 2022; Tanioka et al., 2022).

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414 To examine the role of local bathymetry in the long-lasting tsunami, we choose a well-studied and well-
415 recorded event: the 2011 Mw 9.0 Tohoku tsunami as a reference event and compare the tsunami records
416 of these two events at the same coastal stations. Although the two tsunami events were generated by
417 completely different mechanisms, i.e., large-scale seafloor deformation for the Mw 9.0 megathrust
418 earthquake (Mori et al., 2011) and fast-moving atmospheric waves for the Mw 5.8 volcanic eruption
419 (Matoza et al., 2022) (Titov et al., 2005), they both produced widespread transoceanic tsunamis which
420 were well recorded in the Pacific DART buoys and tide gauges. In the near-field, the 2011 Tohoku
421 earthquake produced runup up to 40 m at Miyako in the Iwate Prefecture, ~70 km from the source (Mori
422 et al., 2011) (Tanioka et al., 2022), while the 2022 HTHH tsunami produced only ~13 m runup in the
423 near field from eyewitness accounts in Kanokupolu, 60 km from the volcano (Lynett et al., 2022).
424 However, in the far-field (>1000 km), we observe comparable tsunami wave heights in certain coastal
425 regions. Based on the tsunami records at 21 tide gauges surrounding the Pacific Ocean, Heidarzadeh &
426 Satake (2013) calculated the average value of the maximum tsunami heights (trough-to-crest) of the 2011
427 Tohoku tsunami is 1.6 m with the largest height of 3.9 m at the Coquimbo Bay in Chile (Heidarzadeh
428 and Satake, 2013). Coincidentally, the statistics of 116 tide gauges in this study also suggest the average
429 tsunami heights of the 2022 HTHH tsunami is around the same order, ~1.2 m, among which, the largest
430 height is 3.6 m at Chañaral Bay in Chile. Interestingly, in the coastal region of South America, the
431 locations of the largest tsunami heights of both events are adjacent (Figure 4a), i.e., Coquimbo (the 2011
432 Tohoku) and Chañaral (The 2022 HTHH).
433 To further compare the far-field hydrodynamic processes between these two events quantitatively, we
434 conduct wavelet analysis for four representative tide gauges distributed across the Pacific Ocean, i.e.
435 coastal gauges at East Cape in New Zealand, Kwajalein Island, Wake Island, and Talcahuano in Chile

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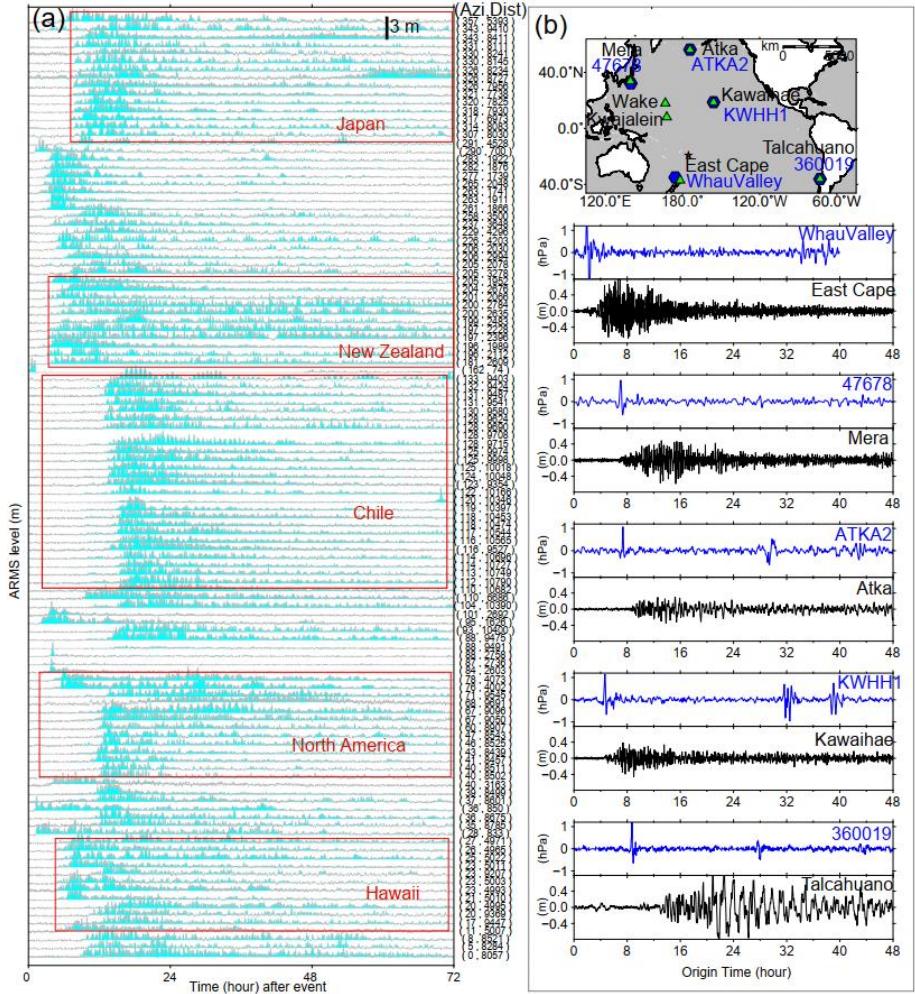
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436 (see their locations in Figures 10b). The temporal changes of tsunami energy of both events can be seen
437 in Figure 11. At each tide gauge, the tsunami energy of the 2011 HTHH (Figure 11a) and the 2022 Tohoku
438 tsunamis (Figure 11b) for the first few hours after the arrivals is nonuniform with different significant
439 peaks distributed within a wide period band of ~3–100 min. Then, the following long-lasting energy of
440 the both at each station presents similar pattern and is concentrated at identical and fairly narrower period
441 channel, i.e., ~20–30 min at East Cape in New Zealand, ~40–60 min at Kwajalein Island, ~10 min at
442 Wake Island, and ~100 min at Talcahuano in Chile, which reflects the local bathymetric effects of natural
443 permanent oscillations (Hu et al., 2022; Satake et al., 2020). Specifically, many bathymetric effects can
444 contribute to the long-lasting tsunami, such as multiple reflections across the basins, or the continental
445 shelves, and the excited tsunami resonance in bays/harbors with variable shapes and sizes (Aranguiz et
446 al., 2019; Satake et al., 2020). For example, tide gauges around New Zealand are primarily distributed in
447 harbors/ports with major natural oscillation modes of ~20–30 min (De Lange and Healy, 1986; Lynett et
448 al., 2022). The first oscillation mode of central Chile is centered around ~100 min (Aranguiz et al., 2019).
449 Consequently, Figure 11 illustrates that the long-lasting tsunami energy of the two events is respectively
450 distributed in 20–30 min period at East Cape in New Zealand and in ~100 min period at Talcahuano in
451 central Chile. The coupling of bathymetric oscillation mode with tsunami containing similar-period wave
452 results in the excitement of tsunami resonance, which amplifies tsunami waves and prolongs the tsunami
453 oscillation at the two stations (Heidarzadeh et al., 2019, 2021; Hu et al., 2022; Wang et al., 2022).

454 Simply put, we do not have clear evidence that atmospheric acoustic-gravity waves from the 2022 HTHH
455 eruption directly contribute to the long-lasting tsunami, but the resonance effect associated with ocean
456 waves could a possible source of increased wave energy and amplification~~atmospheric acoustic-gravity~~
457 ~~waves from the 2022 HTHH eruption do not directly contribute to the long-lasting tsunami, but the~~
458 ~~resonance effect associated with ocean waves theoretically could contribute to it.~~ However, the similarity
459 of far-field hydrodynamic behaviors between the 2022 HTHH volcanic tsunami and the 2011 Tohoku
460 seismogenic tsunami well demonstrates the both went through similar hydrodynamic processes after their
461 arrivals. The consistency favors that the long-lasting tsunami of 2022 HTHH tsunami event can very
462 likely be attributed by the interplays between local bathymetry and conventional tsunami left after each
463 passage of atmospheric waves, which can well explain why the two completely distinct tsunami events
464 possess a comparable duration time.

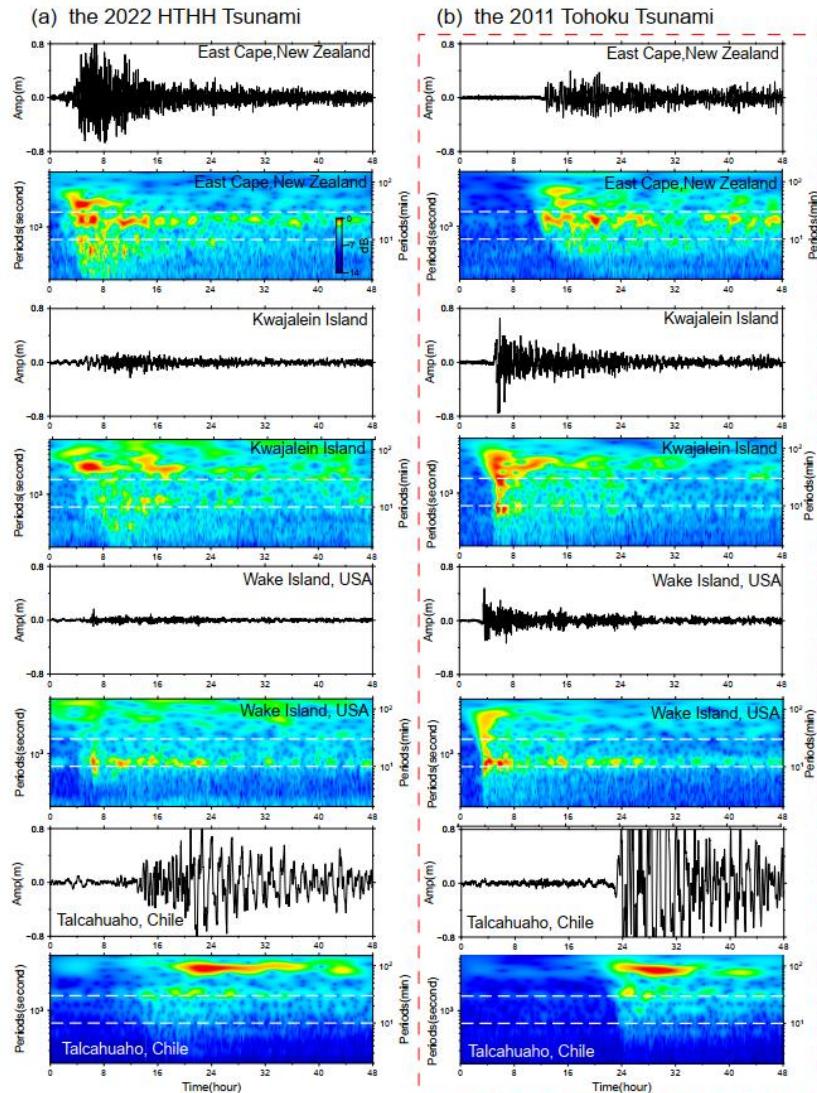
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465
466 **Figure 10. Tsunami duration.** (a) Tsunami durations at Pacific 116 tide gauges through ARMS level
467 approach. (b) the location of barographs (blue curves) and nearby tide gauges (green curves), as
468 well as their waveforms.

469



470
 471 **Figure 11. Wavelet analysis of tsunami waveforms recorded by 4 tide gauges during (a) the 2022**
 472 **HTHH tsunami event, and (b) the 2011 Tohoku tsunami event. Horizontal white dashed lines**
 473 **respectively mark reference periods of 10 min and 30 min.**

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4.3 Challenges for Tsunami Warning

475 The generation mechanisms and hydrodynamic characteristics of the 2022 HTHH volcanic tsunami are
 476 more complicated than pure seismogenic tsunami, which challenge the traditional tsunami warning

477 approach.

478 The first challenge is posed by the tsunami components with propagating velocities faster than the
479 conventional tsunami. The Tonga volcanic tsunami event provides an excellent example which highlights
480 that the tsunamigenic mechanisms are not limited to tectonic activities related with the sudden seafloor
481 displacements, but also include a variety of atmospheric waves with distinct propagation velocities. The
482 tsunami components in 2022 HTHH event generated by the air-sea coupling possess a wide range of
483 velocities from 1000 m/s to 200 m/s. The Lamb waves recorded in both the 2022 HTHH event and the
484 1833 Krakatoa volcanic event traveled along the Earth's surface globally for several times (Carvajal et
485 al., 2022). The tsunami waves produced by Lamb waves, the wave components associated with resonance
486 of the air-sea coupling and their superimposition increase the difficulty of tsunami warning.

487 Another critical challenge is associated with the interplays between tsunami waves and local bathymetry.
488 The tsunami waves left by each passage of the atmospheric waves can interact with local bathymetry at
489 coastlines, such as continental shelves with different slopes, and harbor/bay with different shapes and
490 sizes. The interaction can intensify the tsunami impact and excite a variety of natural oscillation periods.
491 The 2022 HTHH tsunami with an extremely wide period range of ~2–100 min have a great potential to
492 couple with the excited natural oscillations and form extensive tsunami resonance phenomena. The
493 resonance effects result in long-lasting oscillation and delayed tsunami wave peaks. The uncertain
494 arrivals of the maximum tsunami waves pose an extra challenge to tsunami warning.

4955. Conclusion

496 In the study, we explore the tsunamigenic mechanisms and the hydrodynamic characteristics of the 2022
497 HTHH volcanic tsunami event. Through extensive analysis of waveforms recorded by the DART buoys,
498 tide gauges and barometers in the Pacific Ocean, we reach the main findings as follows:

499 (1) We identify four distinct tsunami wave components based on their distinct propagation velocities or
500 period bands (~80–100 min, 10–30 min, 30–40 min, and 3–5 min). The generation mechanisms of these
501 tsunami components range from air-sea coupling to seafloor crustal deformation during the volcanic
502 eruption.

503 (2) The first-arriving tsunami component with 80–100 min period was most likely from shock wave
504 spreading at a velocity of ~1000 m/s in the vicinity of the eruption. This tsunami component was not

505 clearly identified by currently available publication and it's not easy to be visually observed through time
506 series of the waveforms. The physical mechanism is yet to be understood. The second tsunami component
507 with 30–40 min period was from Lamb waves, and was the most discussed tsunami source of this event
508 so far. A thorough analysis of DART measurements indicates that the Lamb waves traveled at the speed
509 of ~340 m/s in the vicinity of the eruption and decreased to ~315 m/s when traveling away due to cooling
510 of the air temperature. The third tsunami component was from some atmospheric gravity wave modes
511 with propagation velocity faster than 200 m/s but slower than Lamb waves. The last tsunami component
512 with the shortest periods 3–5 min was probably produced by partial caldera collapse with estimated
513 dimension of ~0.8–1.8 km.

514 (3) Although the resonance effect with the atmospheric acoustic-gravity waves theoretically could be a
515 source of increased wave energy, its direct contribution to the long-lasting oscillation is not demonstrated
516 yet. However, the comparison of hydrodynamical characteristics between the 2022 HTHH tsunami event
517 and the 2011 Tohoku tsunami event well demonstrated that the interactions between the ocean waves left
518 by atmospheric waves and local bathymetry contribute to the long-lasting Pacific oscillation of the 2022
519 tsunami event. The long-lasting Pacific oscillation of this tsunami event was not only associated with the
520 resonance effect with the atmospheric acoustic-gravity waves, but more importantly the interactions with
521 local bathymetry. The velocities of tsunami waves in deep ocean (especially at Mariana and Tonga-
522 Kermadec trenches) close to those of acoustic Lamb waves and some gravity wave modes produced
523 resonance effects, which supplied energy to the ocean. The comparison of hydrodynamical characteristics
524 between the 2022 HTHH tsunami event and the 2011 Tohoku tsunami event suggests the volcanic
525 tsunami oscillation was prolonged by their interplays with local bathymetry.

526 (4) The extraordinary features of this rare volcanic tsunami event challenge the current tsunami warning
527 system which is mainly designed for seismogenic tsunamis. It is necessary to improve the awareness of
528 people at risks about the potential tsunami hazards associated with volcanic eruptions. New approaches
529 are expected to be developed for tsunami hazard assessments with these unusual sources: various
530 atmospheric waves radiated by volcanic eruptions besides those traditionally recognized, e.g.
531 earthquakes, landslides, caldera collapses and pyroclastic flows etc.

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538 The JAGURS tsunami simulation code is employed for tsunami modelling (Baba et al., 2015;
539 <https://zenodo.org/record/6118212#.Yk98qdtBxPY>). Bathymetry data are obtained from GEBCO at
540 <http://www.gebco.net>. The sea level records in deep ocean are available from the Deep Ocean Assessment
541 and Reporting of Tsunamis (DART) buoy network in the Pacific (<https://nctr.pmel.noaa.gov/Dart/>), and
542 GeoNet New Zealand DART network (<https://tilde.geonet.org.nz>). The sea level records of tide gauges
543 are downloaded from UNESCO/ IOC (<http://www.ioc-sealevelmonitoring.org/>). Barometer data are
544 provided by the following providers: Dirección Meteorológica de Chile
545 (<https://climatologia.meteochile.gob.cl>), NOAA National Weather Service
546 (<https://www.weather.gov/ilm/observations>), Japan Meteorological Agency (<https://www.jma.go.jp>),
547 The UK Met Office Weather Observation (<https://www.metoffice.gov.uk/observations>), and Fiji
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