The Influence of Infragravity Waves on the Safety of Coastal Defences: A Case Study of the Dutch Wadden Sea

Christopher H. Lashley, Sebastiaan N. Jonkman, Jentsje van der Meer, Jeremy D. Bricker, Vincent Vuik

1Dept. of Hydraulic Engineering, Delft University of Technology, Stevinweg 1, 2628 CN Delft, the Netherlands
2Van der Meer Consulting, P.O. Box 11, 8490 AA Akkrum, the Netherlands
3Coastal & Urban Risk & Resilience Dept., IHE Delft, Westvest 7, 2611 AX Delft, the Netherlands
4Dept. of Civil and Environmental Engineering, University of Michigan, 2350 Hayward St., Ann Arbor, MI 48109-2125.

Correspondence to: Jeremy D. Bricker (J.D.Bricker@tudelft.nl)

Abstract. Many coastlines around the world are protected by coastal dikes fronted by shallow foreshores (e.g. saltmarshes and mudflats) that attenuate storm waves and are expected to reduce the likelihood of waves overtopping the dikes behind them. However, most of the studies to-date that assessed their effectiveness have excluded the influence of infragravity (IG) waves, which often dominate in shallow water. Here, we propose a modular and adaptable framework to estimate the probability of coastal dike failure by overtopping waves ($P_f$). The influence of IG waves on wave overtopping is included using an empirical approach, which is first validated against observations made during two recent storms (2015 and 2017). The framework is then applied to compare the $P_f$ of the dikes along the Dutch Wadden Sea coast, with and without the influence of IG waves. Findings show that including IG waves results in 1.1 to 1.6 times higher $P_f$ values, suggesting that safety may be overestimated when they are neglected. This increase is attributed to the influence of the IG waves on the design wave period, and to a lesser extent the wave height, at the dike toe. The spatial variation in this effect, observed for the case considered, highlights its dependence on local conditions—with IG waves showing greater influence at locations with larger offshore waves and shallower water depths. Finally, the change in $P_f$ due to the IG waves varied significantly depending on the empirical wave overtopping model selected, emphasizing the importance of tools developed specifically shallow foreshore environments.

1 Introduction

1.1 Background

Coastal defences (e.g. dikes or seawalls), fronted by wide, shallow foreshores, protect many coastlines around the world. Examples include the sandy foreshores along the Belgian coast (Altomare et al., 2016), the wide shelves of the Mekong Delta, Vietnam (Nguyen et al., 2020) and the intertidal flats of the Wadden Sea along the north coast of the Netherlands (Vuik et al., 2016). These bodies of sediment reduce the water depth in front of the structure such that large incident waves are forced to
break. This reduced wave load at the structure is then expected to improve safety by reducing the likelihood of the structure failing. If vegetation is present, the drag forces exerted by stems, branches and leaves further enhance this attenuation effect.

Several studies sought to quantify the hazard mitigation potential of shallow foreshores, with and without vegetation, including physical model tests (Möller et al., 2014), numerical simulations (Vuik et al., 2016; Willemsen et al., 2020) and field measurements (Garzon et al., 2019). While these studies assessed the ability of the foreshore to attenuate the height of wind-sea and swell (hereafter, SS) waves—with frequencies typically greater than 0.05 Hz—they neglected the influence of infragravity (hereafter, IG) waves.

Under extreme conditions, with large offshore waves and very shallow local water depths, the shoaling and subsequent breaking of SS waves results in energy transfer to lower frequencies and the growth of IG waves, also referred to in literature as “surfbeat” (Bertin et al., 2018; Van Dongeren et al., 2016). IG waves are widely recognized as the driving force behind coastal erosion and flooding along shallow coastlines. Recent reports of their impact include: unexpectedly high wave run-up at the coast of Banneg island, France (Sheremet et al., 2014); extensive damage and casualties along the coral-reef lined coast in the Philippines during Typhoon Haiyan (Roebber and Bricker, 2015); the erosion and overwash of several dunes along the west coast of France (Baumann et al., 2017; Lashley et al., 2019); and damage along Seisho Coast of Japan during Typhoon Lan (Matsuba et al., 2020). Despite this knowledge, IG waves are often not considered in the risk assessment of coastal defences. This oversight is linked to the use of phase-averaged wave models—which inherently exclude IG-wave dynamics—to estimate the wave load at the structure. The impact of this neglect on the safety of coastal defences has yet to be thoroughly investigated.

In the Netherlands, coastal defences are typically designed to resist the volume of water expected to pass over the crest of (or overtop) the structure due to wave action during storms associated with a very high return period (2,000 to 10,000 years). This phenomenon, referred to as wave overtopping, is typically represented by a mean (time-averaged) discharge per meter width of structure, $q$ (m$^3$/s or l/s per m). The probability of failure due to wave overtopping is then determined by assessing the likelihood that the actual discharge ($q_a$) exceeds some critical value ($q_c$), which is dependent on the erosion resistance of the grass-covered landward slope. Following a recent policy revision, the safety standard for the coastal defences in the Netherlands is now defined by an (acceptable) probability of failure. For example, typical values for the Dutch Wadden Sea coast—a shallow, intertidal area in the north of the country—are failure probabilities of 1/1000 and 1/3000 per year. This approach usually considers multiple failure mechanisms (e.g. scour, armour unit instability and slope instability); however, here we limit the analysis to dike failure by wave overtopping, which typically governs the design process.

From a design perspective, the presence of IG waves typically results in higher characteristic values of the two main parameters used to estimate $q_a$: namely, the significant wave height and spectral wave period, both assessed at the structure toe (Lashley et al., 2020a). Vuik et al. (2018b) assessed the overtopping failure probability of an idealised the dike-foreshore system,
representative of the Dutch Wadden Sea coast, considering the effects of vegetation. This study considered the influence of IG waves on the wave period at the toe using the Hofland et al. (2017) empirical model, but neglected their influence on the wave height. Furthermore, Nguyen et al. (2020) later showed that the Hofland et al. (2017) formulae tend to underestimate the development of longer spectral wave periods on foreshore slopes milder than 1:250 (Nguyen et al., 2020)—which is a typical characteristic of the Wadden Sea. As a result, the true influence of IG waves along the Dutch Wadden Sea coast remains unknown.

Oosterlo et al. (2018) carried out a similar probabilistic assessment of a dike with a sandy foreshore, in the south of the Netherlands, but directly included the IG waves by using the XBeach Surfbeat numerical model to estimate the wave parameters at the toe. The authors found, for the considered case, that accounting for the IG waves resulted in $10^2$ times higher overtopping failure probabilities compared to methods that neglected them. This rather striking finding requires further investigation; particularly, to determine if the large IG-wave influence reported by Oosterlo et al. (2018) holds for other cases as well or if it was merely an artefact of the method used.

Lashley et al. (2020b) demonstrated that the influence of IG waves on wave parameters at the toe could be accurately estimated using a combined numerical and empirical approach. In this approach, the phase-averaged wave model (SWAN) is used to simulate the dissipation of SS waves in shallow water; while the IG component is estimated using empirical formulae. Since this approach allows for the accurate representation of IG waves at the dike toe but maintains the utility and speed of phase-averaged wave modelling, it can be applied on a large scale with little computational effort. In the present study, this approach is extended and used as a key component to assess the influence of IG waves on the probability of dike failure along the Dutch Wadden Sea coast.

1.2 Objective and Approach

Previous studies either neglected the influence of IG waves on the probability of failure by wave overtopping or yielded inconclusive results. Consequently, the influence of IG waves on the safety of coastal defences remains unknown. To remedy this, it is the primary aim of the current paper to investigate the influence of IG waves on the probability of failure due to wave overtopping for coastal defences (dikes) with shallow foreshores. This is achieved by first augmenting the probabilistic framework developed by Vuik et al. (2018b) by incorporating newly validated empirical formulae that capture the influence of IG waves on design parameters, following the approach of Lashley et al. (2020b). The modified framework is then used to estimate the probability of dike failure by wave overtopping along the Dutch Wadden Sea coast.

1.3 Outline

This paper is organized as follows: Section 2.1 describes the geographic area that will be the focus of this study. Section 2.2 provides a detailed description of the dike-foreshore system under consideration and the probabilistic framework applied,
including descriptions of the numerical and empirical models therein. Section 2.3 describes the field dataset used to validate the empirical approach for the inclusion of IG waves. In Section 3, the results of the validation are presented, followed by the application of the framework to the wider Wadden Sea coast. Section 4 discusses the results and their implications for practice; and Section 5 concludes the paper by addressing the overall research objective, stating limitations and identifying areas for future work.

2 Materials and Methods

2.1 Study Area

The Dutch Wadden Sea is a shallow, mildly sloping intertidal zone situated between the Netherlands (mainland) and several barrier islands, which shield the area from large North Sea waves (Figure 1). Along the Wadden Sea coast, a system of dikes fronted by saltmarshes and mudflats exists. In the present study, we consider the stretch of dikes with shallow foreshores—that is, with bed levels at the toe either just below or a few meters higher than mean sea level (NAP)—that are typically impacted by North-westerly waves during storms (Figure 1). Further information on storm conditions in the study area is presented in Section 2.3 and Table 1.

Figure 1: Location of the Dutch Wadden Sea with reference to wider Europe (inset). Circle markers indicate dikes considered, while ‘x’ markers highlight those that were excluded from the analysis. Star marker indicates the location of the field site at Uithuizerwad.
The analysis includes the dikes from the city of Harlingen, in the province of Friesland to those west of Eemshaven, in the province of Groningen; but it excludes the flood defences in front of harbours and areas referred to as summer polders (Figure 1). Summer polders are low-lying, embanked areas situated in front of the dike and are usually dry in the summer months but may flood during winter storms. As these polders extend for several kilometres, the 1 km transect approach taken here would not be representative.

### 2.2 Model Framework and System Description

The model framework used to compute the probability of flooding due to wave overtopping is presented in Figure 2, and is modified after Vuik et al. (2018b) to include the effect of IG waves. Boundary conditions of offshore wave heights, periods and water levels are transformed over the foreshore to the structure toe using SWAN. These SWAN estimates at the toe are then modified using empirical formulae to account for IG waves. These estimates are then used as input to calculate the actual overtopping discharge, $q_a$. The probability of failure by wave overtopping, $P(Z < 0)$ is then obtained using the open-source implementation of the First Order Reliability Method (FORM (Hasofer and Lind, 1974)).

![Figure 2: Model framework.](https://doi.org/10.5194/nhess-2021-211)
While the above framework follows that of Vuik et al. (2018b), there are noteworthy differences between the two approaches. Firstly, the influence of IG waves on both the wave height and period at the toe is considered—using empirical formulae that are valid for a wide range of foreshore slopes (10 ≤ cot(m) ≤ 1000). Secondly, the effect of wind on wave transformation is neglected here due to close proximity to the shoreline—within 1 km. Lastly, as it is very likely that almost all vegetation will flatten or break under extreme forcing (Möller et al., 2014; Vuik et al., 2018a), wave attenuation by vegetation is not included in the probabilistic analysis. That said, the effect of vegetation (should it remain standing) is demonstrated for one location (Uithuizerwad field site, Figure 1) in Section 3.2; and treated as part of the discussion (Section 4.1). The individual components of the model framework are described in detail below. A visual representation of the dike-foreshore system, and the various framework components, is provided in Figure 3.

![Diagram of dike-foreshore system](https://example.com/diagram.png)

**Figure 3:** Schematic representation of the dike-foreshore system, showing the areas where key parameters and tools are applied. These parameters and tools are introduced throughout Section 2.2.

### 2.2.1 Boundary Conditions

#### Offshore Waves and Water Levels

To obtain the offshore water levels ($\bar{\eta}$), wave heights ($H_{m0}$) and periods ($T_{m-1,0,deep}$), Hydra-NL (Duits, 2019) was applied—a probabilistic application designed specifically for the assessment and design of flood defences in the Netherlands. It uses statistics of wind speed, wind direction, water level and their respective correlations to find the corresponding wave characteristics in a pre-calculated database, obtained using the phase-averaged numerical model, SWAN. The tool provides estimates along the entire Dutch Wadden coast, every 250 m, a few hundred meters offshore. To reduce the overall computational demand, the output locations were reduced to one every 1.5 km (Figure 1).

The wave heights estimated by Hydra-NL were reverse shoaled, using linear wave theory, to an offshore point approximately 1 km from the dike toe. This was done so as to remove the influence of the foreshore, similar to the approach of Van Osselen...
(2016), to obtain estimates of the offshore wave height \( (H_{m0,deep}) \) for use with Equation 12. This approach also allowed additional foreshore processes (e.g. vegetation) that were not considered in the original SWAN calculations in the Hydra-NL database to be modelled here (see Section 2.2.2). Compared to the original Hydra-NL estimates, the average change in significant wave height by reverse shoaling was +2%, with a maximum of +4% and minimum of -9%. As this approach neglects friction, refraction and any local generation or breaking that may occur, it is considered here to be an approximate estimate. For each location, five exceedance probabilities were considered in Hydra-NL: 1/100, 1/300, 1/1000, 1/3000 and 1/10000 per year (Table 1). Using the \( \bar{\eta} \), \( H_{m0,deep} \) and \( T_{m-1,0,deep} \) estimates for each probability of exceedance, Weibull distribution parameters—namely, scale and shape parameters—were derived to accurately describe the extremes. The range of the scale and shape parameter values is provided in Table A. 1 (Appendix A).

### Table 1: Characteristic values for offshore waves and water levels at the Uithuizerwad field site.

<table>
<thead>
<tr>
<th>Variables</th>
<th>Unit</th>
<th>1/year</th>
<th>1/100</th>
<th>1/300</th>
<th>1/1000</th>
<th>1/3000</th>
<th>1/10000</th>
</tr>
</thead>
<tbody>
<tr>
<td>( \bar{\eta} )</td>
<td>m</td>
<td>4.29</td>
<td>4.62</td>
<td>4.95</td>
<td>5.26</td>
<td>5.59</td>
<td></td>
</tr>
<tr>
<td>( H_{m0,deep} )</td>
<td></td>
<td>1.40</td>
<td>1.58</td>
<td>1.77</td>
<td>1.94</td>
<td>2.14</td>
<td></td>
</tr>
<tr>
<td>( T_{m-1,0,deep} )</td>
<td>s</td>
<td>5.02</td>
<td>5.55</td>
<td>6.04</td>
<td>6.52</td>
<td>6.99</td>
<td></td>
</tr>
</tbody>
</table>

**Dike-foreshore characteristics**

The bathymetry of the Dutch coast, from dry land up to the 20 m isobath in the North Sea, is continuously measured (at least once every seven years) by the Dutch government (Rijkswaterstaat). This dataset, referred to as “Vaklodingen” (Wiegmann et al., 2005), covers the Wadden Sea with a 20-m grid resolution (Figure 4). Cross-shore transects of approximately 1 km, at intervals corresponding to the Hydra-NL output locations, were extracted considering a NW to SE orientation—in-line with the dominant wind/wave direction during storms (NW) (Vuik et al., 2018b), similar to the approach taken by Willemsen et al. (2020). By aligning the transect with the dominant wind/wave direction, the influence of wave obliqueness on wave propagation and overtopping may be neglected.

For each transect, the mean elevations \( (z_b) \) and average foreshore slopes \( (\tan(m)) \) were obtained. To account for variations in bathymetry and measurement inaccuracies, \( z_b \) was treated as a normally distributed parameter with a standard deviation of 0.2 m. While the actual bathymetry is used for the numerical modelling of the SS waves (Section 2.2.2), the estimated \( \tan(m) \) is necessary for use with the empirical formulae for the influence of the IG waves (Section 2.2.2). Given the range of validity of the empirical formulae (Equations 12 and 13), a minimum foreshore slope of 1/1000 is considered here.
Figure 4: Subset of the Vaklodingen bathymetry dataset showing a) a NW-SE oriented transect at the field site location, Uithuizerwad (black line) and b) the cross-shore profile of this transect with the estimated average foreshore slope at the field site location.

Given the significant influence of the dike geometry on the calculated probability of failure (Section 3.2.2), the crest level ($z_c$) and average dike slope ($\tan(\alpha)$) are treated here as deterministic parameters with the same values applied to each location. This was done to remove the influence of variations in dike geometry on the calculated failure probabilities and allow the analysis to focus on what occurs over the foreshore. The crest levels were set to 6 m + NAP, corresponding to the required safety level (probability of failure less than 1/1000 per year), but also varied as part of the sensitivity analysis. Similarly, the dike slopes were set to 1/7 to represent the average slope of a typical Wadden Sea dike, which is often characterised by 1:4 upper and lower slopes separated by a mildly sloping berm. Note that an analysis of the sensitivity of the estimated probability of failure to variations in $z_c$, including its treatment as deterministic versus stochastic, is provided in Section 3.2.2. It should be noted that the actual crest levels of the Dutch Wadden Sea dikes typically exceed 8 m + NAP; however, a crest level this high would result in extremely small failure probabilities, which would distract from the findings herein.
2.2.2 Wave Transformation

Numerical Model for SS Waves: SWAN

SWAN is a third-generation, phase-averaged wave model used to estimate the generation (by wind), propagation and dissipation (by depth-induced breaking and bottom friction) of waves from offshore to the structure toe. This includes wave-wave interactions, in both deep and shallow water, and wave-induced setup; but neglects wave-induced currents and the generation or propagation of IG waves. SWAN computes the spectral evolution of wave action density ($A$) in space and time. For stationary 1D simulations, the governing equations follow:

\[
\frac{\partial c_x A}{\partial x} = \frac{S_{tot}}{\omega}, \tag{1}
\]

where $c_x$ is the propagation velocity of wave energy in the $x$-direction, $\omega$ is the frequency and $S_{tot}$ may include dissipation terms due to depth-limited wave breaking ($\overline{D}_w$), vegetation ($\overline{D}_v$) and bottom friction; and energy transfer terms.

To simulate depth-limited wave breaking, SWAN uses the following parametric dissipation model, by default (Battjes and Janssen, 1978):

\[
\overline{D}_w = \frac{4}{\rho g} f_{\text{mean}} Q_b H_{max}^2, \tag{2}
\]

and $Q_b$ is estimated as:

\[
\frac{1 - Q_b}{\ln Q_b} = -8 \frac{E_{tot}}{H_{max}^2}, \tag{3}
\]

where $f_{\text{mean}}$ is the mean wave frequency, $H_{max} = \gamma_{bj} h$ and $E_{tot}$ is the total wave-energy variance. Here, the breaker parameter ($\gamma_{bj}$) is based on the offshore wave steepness, $s_0 = H_{rms}/L_0$ (Battjes and Stive, 1985):

\[
\gamma_{bj} = 0.5 + 0.4 \tanh(33 \cdot s_0), \tag{4}
\]

where $H_{rms}$ is the root mean square wave height, with $H_{rms} = 8 \cdot E_{tot}$. Following Vuik et al. (2018b), $\gamma_{bj}$ is treated as a normally-distributed parameter with a standard deviation of 0.05.

As recommended by Baron-Hyppolite et al. (2018), the explicit vegetation representation in SWAN—which was implemented by Suzuki et al. (2012)—is applied. This method represents vegetation as rigid cylinders, following the approach of Dalrymple et al. (1984) modified for irregular waves by Mendez and Losada (2004). In this approach, the mean rate of energy dissipation per unit horizontal area due to wave damping by vegetation ($\epsilon_v$) is given by:
\[ \bar{D}_v = \frac{1}{2g} \rho C_D b_v N_v \left( \frac{g k}{2f_{mean}} \right)^3 \frac{\sinh^3 k h_v + 3 \sinh k h_v}{3k \cosh^3 k h} H_{rms}^3, \]  

where \( \rho \) is the density of water, \( g \) is the gravitational acceleration, \( k \) is the mean wave number, \( h \) is the water depth, \( C_D(0.4) \) is the drag coefficient, \( b_v \) (3 mm) is the stem diameter, \( N_v \) (1200 stems/m\(^2\)) is the vegetation density and \( h_v \) (0.3 m) is the vegetation height; where the values in parentheses are representative of saltmarshes in the Netherlands (Vuik et al., 2016).

Deep-water processes such as white-capping, wind and quadruplet wave-wave interactions were disabled; while triad wave-wave interactions, a shallow water process, was activated. All other model parameters were kept at their default values. For all simulations, a constant grid spacing of 2.5 m was applied. This corresponded to approximately 15-20 grid cells per deep-water wavelength.

**Empirical Formulae for the Influence of IG Waves**

The influence of the IG waves may be represented by an increase in the design parameters, namely: the total significant wave height \( H_{m0,\text{toe}} \) and spectral wave period \( T_{m-1,0,\text{toe}} \) at the toe based on incident waves, where:

\[ H_{m0} = \sqrt{H_{m0,SS}^2 + H_{m0,IG}^2}, \]  

\[ H_{m0,SS} = 4 \int_{0.05}^{1} C_{\eta \eta} \, df, \]  

\[ H_{m0,IG} = 4 \int_{0.005}^{0.05} C_{\eta \eta} \, df, \]

where \( C_{\eta \eta}(f) \) is the wave energy density and 0.05 Hz is the frequency separating SS and IG motions. This choice of split frequency is consistent with previous studies in the area (Engelstad et al., 2017; De Bakker et al., 2014) and coincides with the minimum in spectral density in the observed wave spectra at the field site (Section 2.3); and

\[ T_{m-1,0} = \frac{m_{-1}}{m_0}, \]  

where,
\[ m_n = \sqrt{\int_{0.005}^{1} C_{\eta} \cdot f^n \, df}. \] \tag{10} 

**Influence on Significant Wave Height at Toe**

Lashley et al. (2020a) investigated the influence of various dike-foreshore parameters on wave conditions at the dike toe using the XBeach non-hydrostatic numerical model (Kingsday release) in both 1D and 2D mode. In each simulation, the offshore wave height \( H_{m0,\text{deep}} \), offshore spectral wave period \( T_{m-1,0,\text{deep}} \), wave directional spreading \( \sigma_{\text{dir}} \), water depth at the toe \( h_{\text{toe}} \), foreshore slope \( m \), width of vegetated cover \( W_{\text{veg}} \) and structure slope \( \alpha \) were systematically varied, following a one-factor-at-a-time (OAT approach). From the resulting dataset of 672 different numerical estimates of \( T_{m-1,0,\text{toe}} \), \( H_{m0,\text{IG},\text{toe}} \) and \( H_{m0,\text{SS},\text{toe}} \), the following formula for \( H_{m0,\text{IG},\text{toe}} \) was derived:

\[ H_{m0,\text{IG},\text{toe}} = \bar{H}_{\text{IG}} \cdot H_{m0,\text{SS},\text{toe}}, \] \tag{11}

where the relative magnitude of the IG waves,

\[ \bar{H}_{\text{IG}} = 0.36 \cdot H_{m0,\text{deep}}^{0.5} \cdot \bar{\gamma}_{\sigma} \cdot \bar{\gamma}_{h} \cdot \bar{\gamma}_{f} \cdot \bar{\gamma}_{v} \cdot \bar{\gamma}_{d}, \] \tag{12}

where \( \bar{\gamma}_{\sigma}, \bar{\gamma}_{h}, \bar{\gamma}_{f}, \bar{\gamma}_{v}, \bar{\gamma}_{d} \) are influence factors for wave directional spreading, water depth at the toe, foreshore slope, vegetated cover and structure slope, respectively; see Lashley et al. (2020a) for further details on how each influence factor is determined. However, it should be noted that the influence of vegetation on IG waves was assessed for very shallow conditions \( h_{\text{toe}}/h_v = 3.3 \); therefore, the performance of Equation 12 for vegetation with larger water depth to stem height ratios is yet to be verified. For the probabilistic analysis, Equation 12 is multiplied by normally distributed factor \( f_{\text{IG}} \) with a mean (0.99) and standard deviation (0.18) based on the bias and scatter observed during its derivation, respectively (Lashley et al., 2020a).

The above approach accounts for IG-wave generation by: i) bound-wave shoaling over mildly-sloping bathymetry; and ii) the temporal variation in the location of breaking waves, known as the break-point forcing mechanism (Battjes, 2004). However, it does not account for IG waves that may be refractively trapped, known as edge waves or those reflected from a distant coast, known as leaky waves (Elgar et al., 1992; Bertin et al., 2018; Reniers et al., 2021). The relevance of these free IG waves to the Dutch Wadden Sea coast is still to be confirmed.

**Influence on Spectral Wave Period at Toe**

The existing method to estimate the increase in spectral wave period at the toe due to IG waves in shallow water, developed by Hofland et al. (2017), was based on laboratory tests with foreshore slopes, \( 35 < \cot(m) < 250 \). While the method proved accurate within this range, as shown by (Lashley et al., 2020b), it tended to underestimate \( T_{m-1,0,\text{toe}} \) for slopes gentler than...
1:250 (Nguyen et al., 2020). As foreshores in the Wadden Sea are typically 1:500 and gentler, a new formulation for $T_{m-1,0,\text{toe}}$ is derived here—using the above-mentioned numerical dataset (Lashley et al., 2020a).

Since $T_{m-1,0,\text{toe}}$ and $\tilde{H}_{IG}$ both describe the amount of energy in the IG band compared to the SS band, it stands to reason that a simple relation should exist between the two parameters. From the Lashley et al. (2020a) numerical dataset with $10 < \cot(m) < 1000$, the following relationship between $T_{m-1,0,\text{toe}}$, $T_{m-1,0,\text{deep}}$, $\tilde{H}_{IG}$, and $\cot \alpha_{\text{fore}}$ was found ($R^2 = 0.92$):

$$\frac{T_{m-1,0,\text{toe}}}{T_{m-1,0,\text{deep}}} = \begin{cases} 1.59 \cdot \tilde{H}_{IG}^{0.69} \cdot \cot(m)^{0.17} & \frac{h_{\text{toe}}}{H_{m0,\text{deep}}} \leq 1, \\ 1 & \frac{h_{\text{toe}}}{H_{m0,\text{deep}}} > 1. \end{cases}$$

(13)

Further details on the derivation of Equation 13 and its performance in comparison to the Hofland et al. (2017) model are provided in Appendix B. For the probabilistic analysis, Equation 13 is multiplied by normally distributed factor ($f_{Tm}$) with a mean (0.99) and standard deviation (0.17) based on the bias and scatter observed during its derivation, respectively.

### 2.2.3 Wave Overtopping

**Empirical Formulae for Actual Wave Overtopping**

In the present study, the overtopping formula proposed by Van Gent (1999) (Equation 14) is applied. This formula was chosen because it was developed specifically for shallow foreshores, considers the influence of both SS and IG waves and is considered valid for a wide range of breaker parameter ($\xi_{m-1,0}$) values.

$$\frac{q_a}{\sqrt{g \cdot H_{m0,\text{toe}}^3}} = 10^c \cdot \exp \left( - \frac{R_c}{H_{m0,\text{toe}} \cdot (0.33 + 0.022 \cdot \xi_{m-1,0})} \right).$$

(14)

where,

$$\xi_{m-1,0} = \frac{\tan(\alpha)}{\sqrt{H_{m0,\text{toe}} / L_{m-1,0}}}$$

(15)

$$L_{m-1,0} = \frac{g \cdot T_{m-1,0,\text{toe}}^2}{2\pi}$$

(16)
where $g$ is the gravitational constant of acceleration, $\alpha_{dike}$ is the dike slope, $\xi_{m-1,0}$ is the Iribarren number (also referred to as the breaker parameter) and $L_{m-1,0}$ is a fictitious wavelength based on the spectral wave period at the toe. It is important to note that $H_{m0,\text{toe}}$ and $T_{m-1,0,\text{toe}}$ in the above equations are based on the incident waves (i.e. without the influence of wave reflection at the structure). The empirical coefficient ($c$) is a normally-distributed parameter with a mean of 0.92 and a standard deviation of 0.24. Here, Equation 14 here is applied to all locations, regardless of $\xi_{m-1,0}$ value. However, it should be noted that this approach does not coincide with the current standard (EurOtop, 2018). EurOtop (2018) recommends that different formulae be applied depending on the $\xi_{m-1,0}$ value (Van der Meer and Bruce, 2014; Altomare et al., 2016). However, due to the gentle dike (1:7) and foreshore slopes (1:600, on average) considered here, applying the EurOtop (2018) approach proved challenging. This is discussed in detail in Section 4.2.

**Critical Wave Overtopping**

The erosion resistance of the grass-covered landward slope of the dike is described by a critical or tolerable overtopping discharge ($q_c$). Given the significant influence of this parameter on the probability of dike failure by wave overtopping (Section 3.2.2), it is treated here as a deterministic parameter with a value of 50 l/s/m for each location. In this way, the influence of other parameters, such as those linked to the IG waves, can be better assessed. An analysis of the sensitivity of the estimated probability of failure to changes in $q_c$, is provided in Section 3.2.2. This analysis also demonstrates how the probability of failure would change if $q_c$ were instead treated as a stochastic parameter.

### 2.2.4 Probabilistic Methods

**FORM**

The open-source implementation of FORM, part of OpenEarthTools (Van Koningsveld et al., 2010), is used to evaluate the limit state function (LSF) for any possible combination of input variables, which are each described by probability distributions. The following LSF is considered here (Oosterlo et al., 2018):

$$ Z = \log q_c - \log q_a, $$

(17)

where $Z$ is the limit state considering the critical ($q_c$) and actual ($q_a$) overtopping discharges, which represent the resistance and load, respectively; and the probability of failure by wave overtopping, $P_f = P(Z < 0)$ or $P(q_a > q_c)$.

FORM simplifies the mathematical problem by linearizing the LSF and transforming all probability distributions to equivalent normal distributions. $P_f$ is then expressed in terms of a reliability index ($\beta$), which represents the minimum distance from the most probable failure point on the limit state surface ($Z = 0$), referred to as the design point, to the origin of the transformed coordinate system (Hasofer and Lind, 1974).
\[ \beta = \frac{\mu_z}{\sigma_z}, \quad (18) \]

where are the mean and standard deviation of the limit-state function \((Z)\); and

\[ P_f = \Phi(-\beta), \quad (19) \]

where \(\Phi\) is the cumulative distribution function for a standard normal variable.

FORM starts in a user-defined position in the probability density functions of all variables (e.g. the mean value). It then uses an iterative procedure to update the design point until convergence is achieved (Vuik et al., 2018b). In each iteration, FORM tests how strong the LSF responds to a perturbation of each individual variable, \(X_i\). The response is expressed in terms of the partial derivative \(\partial Z / \partial X_i\) which are then used to calculate sensitivity factors \((\alpha_{sf,i})\):

\[ \alpha_{sf,i} = \frac{\partial Z / \partial X_i \cdot \sigma_i}{\sigma_z}, \quad (20) \]

where \(\alpha_{sf,i}\) represents the relative importance of the uncertainty in each stochastic parameter, such that \(\sqrt{\sum_{i=1}^{n} \alpha_{sf,i}^2} = 1\). Uncertainties in parameters with large \(\alpha_{sf}\)-values—that is, values closer to 1—are considered to be significant, such that a small change in the uncertainty of that parameter would result in a relatively large change in the reliability index \((\beta)\). However, the uncertainty in parameters with \(\alpha_{sf}\)-values close to zero have minor relative importance and those parameters may be treated as deterministic (Kjerengtroen and Comer, 1996).

**Dependencies**

The following (Gaussian) dependencies between variables are imposed (Table 2); all other variables are considered independent.

<table>
<thead>
<tr>
<th>Variables</th>
<th>Source</th>
<th>(\rho)</th>
</tr>
</thead>
<tbody>
<tr>
<td>(\bar{\eta})</td>
<td>(H_{m0,deep})</td>
<td>0.97</td>
</tr>
<tr>
<td>(\bar{\eta})</td>
<td>(T_{m-1,0,deep})</td>
<td>0.96</td>
</tr>
<tr>
<td>(H_{m0,deep})</td>
<td>(T_{m-1,0,deep})</td>
<td>0.99</td>
</tr>
</tbody>
</table>

(*Table 2: Pearson correlation coefficients (\(\rho\)) for Gaussian dependence between boundary conditions (\(\bar{\eta}\), \(H_{m0,deep}\), and \(T_{m-1,0,deep}\)).*)
Foreshore Scenarios

In order to investigate the effect of IG waves on the $P_f$, we consider the following two scenarios:

1) SS wave breaking: where the influence of the foreshore bathymetry on incident SS waves is considered but IG waves are neglected; and

2) SS wave breaking and IG waves: where the influence of the foreshore bathymetry on both SS and IG waves are considered.

In the Netherlands, Scenario 1 represents standard practice, as the influence of IG waves are typically not considered during safety assessments. By assessing the difference in $P_f$ between Scenario 1 and 2—hereafter, referred to as $P_{f1}$ and $P_{f2}$, respectively—the influence of the IG waves may be quantified. Note that in both scenarios vegetation is assumed to be flattened, broken or not present (mudflats) in the analysis of the wider Wadden Sea coast. However, the influence of standing vegetation on the $P_f$ is demonstrated for a single case at the Uithuizerwad location in Section 3.2.

2.3 Field Data for Model Validation

The performance of the combined numerical and empirical wave modelling approach is assessed by comparing estimates to storm data measured at Uithuizerwad, the Dutch Wadden Sea (Figure 1)—where the dike is fronted by vegetated foreshore with an average foreshore slope of 1:600. In this way, the ability of the approach to accurately represent the processes occurring over the foreshore is verified, namely: i) the decrease in SS waves due to depth-induced breaking over the foreshore and ii) the increase in wave height and period at the toe due to IG waves. This dataset is described below.

Two field campaigns were carried out in winter 2014/2015 and 2016/2017 at Uithuizerwad (Figure 1), capturing severe storms on January 11, 2015 and January 13, 2017, both with exceedance probabilities of approximately 1/5 per year (Zhu et al., 2020). Here, we consider two transects of wave gauges that captured the change in wave conditions from the marsh edge to dike toe (Figure 5). In January 2015, a transect of 5 pressure gauges (Ocean Sensor Systems, Inc., USA) was deployed nearshore, each sampling at 5 Hz over a period of 7 minutes, every 15 minutes (Figure 5a). In January 2017, the set-up of the experiment was slightly altered with four gauges deployed, each sampling continuously at 5 Hz (Figure 5b).
The pressure signal from each gauge was translated into time series of surface elevations, using linear wave theory to adjust for attenuation of the pressure signal with depth. After that, a Fourier transform was performed, to transform the data from the time domain to the frequency domain (Hann window, 50% overlap). To improve the frequency resolution of the resulting wave spectra, measurements from two successive bursts were combined into a single time record. For the 2015 dataset, this yielded spectra with 19 degrees of freedom and a frequency resolution of 0.011 Hz; while the analysis of the 2017 dataset yielded spectra with 31 degrees of freedom and a frequency resolution of 0.0089 Hz. The measured wave and water level conditions at the marsh edge for the 2015 and 2017 winter storms are summarized in Table 3.

**Table 3: Measured wave and water level conditions at the marsh edge during the 2015 and 2017 winter storms at the Uithuizerwad field site.**

<table>
<thead>
<tr>
<th>Variables</th>
<th>Unit</th>
<th>Year 2015</th>
<th>Year 2017</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\bar{\eta}$</td>
<td>m</td>
<td>3.12</td>
<td>3.25</td>
</tr>
<tr>
<td>$H_{m0, deep}$</td>
<td></td>
<td>0.71</td>
<td>0.84</td>
</tr>
<tr>
<td>$T_{m1, deep}$</td>
<td>s</td>
<td>5.02</td>
<td>5.31</td>
</tr>
</tbody>
</table>
3 Results

3.1 Validation of Wave Modelling

In this section, the comparison between the combined numerical and empirical modelling approach for wave transformation (Section 2.2.2) and the field measurements (Section 2.3) is presented. For both the 2015 and 2017 storms, SWAN is set up as a transect (1D) in line with wave sensors (Figure 5). In each simulation, the numerical model is forced at its boundary with the measured wave spectra at the most offshore wave sensor. SWAN is able to capture the dissipation of SS waves due to the combined effects of the shallow bathymetry and vegetation (Figure 6). In 2015, the modelled SS-wave attenuation from the most offshore gauge to the dike toe was 56%, half of which (28%) was due to depth-induced wave breaking over the shallow bathymetry alone. Similarly, modelled SS-wave attenuation in 2017 was 63% with 39% due to depth-induced wave breaking alone.
Figure 6: Comparison of measured and modelled significant wave heights in the SS bands ($H_{m0,SS}$) at the peaks of the a) January 2015 and b) January 2017 storms at Uithuizerwad (see Figure 5 for gauge locations). Error bars represent the uncertainty in the estimates based on the standard deviation of Equation 4. Panel ‘c’ shows the bed level and vegetated cover.

At the toe of the dike, Equations 12 and 13 are used to estimate the increase in the relative magnitude of the IG waves, $\tilde{H}_{IG}$ (Figure 7) and the associated increase in spectral wave period relative to its deep-water value, $T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}}$ (Figure 8), respectively. Compared to the measurements at the toe, Equation 12 produced an average error of -5%; that is, predictions were, on average, 5% lower than the measurements. In Figure 8, estimates of $T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}}$ made by SWAN and the Hofland et al. (2017) formula are also presented for comparison. For the two storms, SWAN produced an average error of -48% compared to Equation 13 with 11% error; thus indicating the relevance of IG waves. Similarly, the Hofland et al. (2017) formula produced an average error of -55%. As the Hofland et al. (2017) formula is based on tests with $35 \leq \cot(m) \leq 250$, these results further indicate that it should not be applied outside of this range and highlights the added value of Equation 13—which considers slopes as gentle as 1:1000.
Figure 7: Comparison of measured and modelled relative magnitude of the IG waves ($\tilde{H}_{IG}$) at the peaks of the a) January 2015 and b) January 2017 storms at Uithuizerwad (see Figure 5 for gauge locations). Error bars represent the uncertainty in the estimates based on the standard deviation of Equation 12. Panel ‘c’ shows the bed level and vegetated cover.
Figure 8: Comparison of measured and modelled relative spectral wave period ($T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}}$) at the peaks of the a) January 2015 and b) January 2017 storms at Uithuizerwad (see Figure 5 for gauge locations). Error bars represent the uncertainty in the estimates based on the standard deviations of Equations 12, 13 and the Hofland et al. (2017) formula (Appendix B). Panel ‘c’ shows the bed level and vegetated cover.

3.2 Probability of Failure: Uithuizerwad Case

3.2.1 Influence of IG waves on Probability of Failure

Using the modified probabilistic framework (Section 2.2), the annual probabilities of failure due to wave overtopping ($P_f$) at Uithuizerwad are presented in Figure 9a for the two scenarios considered and an additional scenario to assess the influence of standing vegetation. The calculated $P_f$ is presented alongside estimates of the wave height and period at the toe for a proxy storm with an exceedance probability of 1/3000 per year (Figure 9b), which is in line with the safety standard.

Scenario 2 (SS + IG) results in a $P_f$ 1.3 times larger than that of scenario 1 (SS) (Figure 9a). An increase which corresponds well with the increase in the spectral wave period at the toe ($T_{m-1,0,\text{toe}}$), and to a lesser extent, the increase in wave height at
the toe \((H_{m0,\text{toe}})\) due to IG waves (Figure 9b). If the effects of standing vegetation are considered (SS + IG + Veg), the \(P_f\) is reduced by one order of magnitude (Figure 9a). This is due to the wave attenuation effect of the vegetation, which reduces both \(H_{m0,\text{toe}}\) and \(T_{m-1,0,\text{toe}}\) compared to scenario 2 (SS + IG) alone (Figure 9b).

![Figure 9: Relationship between a) annual failure probabilities at Uithuizerwad for the two scenarios considered and b) physical parameters for a proxy storm with an exceedance probability of 1/3000 per year and a still water level of 5.26 m + NAP. Crest level of 6 m + NAP and dike slope of 1:7 considered. Note that the influence of vegetation may be overestimated (see Section 4.1).](image)

### 3.2.2 Influence of Parameter Values and Uncertainty on the Probability of Failure

The influence of the dike crest level \((z_c)\) on the calculated \(P_f\) is presented in Figure 10a. It can be seen that the influence of the IG waves increases with increasing \(z_c\) value. This is due to the fact that the large load \((q_a)\) needed for failure of a higher dike is reached earlier when IG waves are included. On the other hand, the difference between scenario 1 and scenario 2 remains rather constant with varying values for the critical wave overtopping discharge, \(q_{\text{crit}}\) while the magnitude of the calculated \(P_f\) decreases by a factor of \(O(10)\) when \(q_{\text{crit}}\) is increased by the same magnitude (Figure 10b).

With respect to the stochastic parameters, the sensitivity of the calculated \(P_f\) to uncertainties in each parameter is assessed using the FORM \(\alpha_{sf}\)-values (Section 2.2.4, Figure 11). Negative \(\alpha_{sf}\)-values represent variables that contribute to the load by increasing the actual overtopping discharge \((q_a)\). In both scenarios, the uncertainty in the offshore water level \((\bar{\eta})\) dominates the probability of failure with \(\alpha \leq -0.96\). This is expected since the dike is unlikely to fail without extreme water levels (i.e. a severe storm). In scenario 1 (Figure 11a), the variables that also contribute to the load are: the empirical wave overtopping
coefficient (c)—since larger c values increase $q_a$ (Section 2.2.3); the SWAN breaker parameters ($\gamma_{BJ}$), which controls the magnitude of breaking waves, such that higher $\gamma_{BJ}$ lead to larger wave heights at the structure toe and thus larger $q_a$; and the offshore wave forcing parameters ($H_{m0,deep}$ and $T_{m-1,0,deep}$).

As expected, when the influence of the IG waves is included in the analysis (scenario 2, Figure 11b) the uncertainty in factors for the relative IG-wave height, $f_{IG}$ ($\alpha = -0.03$) and spectral wave period, $f_{TM}$ ($\alpha = -0.02$) also contribute to the load, as larger $H_{m0, toe}$ and $T_{m-1,0, toe}$ values increase $q_a$. This suggests that the calculated $P_f$ is indeed sensitive to the accuracy of Equations
12 and 13. The uncertainty in bed level \((z_b)\) also contributes to the load due to its influence on the water depth at the toe, which directly influences the relative magnitude of the IG waves at the dike toe \((\tilde{H}_{IG})\).

![Figure 11: Sensitivity factors \((\alpha_{sf})\) ranked from left to right in order of importance, for a) scenario 1: SS and b) scenario 2: SS + IG. Negative values indicate load parameters.](image)

Overall, the results of the validation suggest that Equations 12 and 13 may be applied to the area with reasonable accuracy. Likewise, the results of the application of the model framework to the case at Uithuizerwad are in line with expectations. The calculated failure probability for scenario 1 (SS) are similar to the assumed safety standard (less than 1/1000 per year) and the differences observed between the scenarios show clear relationships with physical wave parameters at the dike toe, namely the significant wave height and spectral wave period that determine the magnitude of wave overtopping Figure 9b. With confidence in the model framework, it is applied to the wider Wadden Sea area (Figure 1) for a spatial analysis of the \(P_f\).
3.3 Probability of Failure for the wider Wadden Sea Area

As a next step, the dikes of the wider Wadden Sea area are considered from the city of Harlingen to those west of Eemshaven in the city of Groningen. Again, we apply the assumption of a constant dike height of 6 m (above NAP) and a slope 1:7 for all the dikes in the area. For scenario 1 (only SS waves), the probability of failure due to wave overtopping \( P_{f1} \) ranges from \( 1.1 \times 10^{-4} \) to \( 3.1 \times 10^{-2} \) per year with an average value of \( 4.6 \times 10^{-3} \) per year (Figure 12a). These variations in \( P_{f1} \) are due to: i) the level of exposure—where areas behind inlets are exposed to higher values of \( H_{m0}\text{,deep} \) and \( T_{m-1,0}\text{,deep} \) compared to those behind the barrier islands; ii) variations in the mean water level \( \bar{\eta} \), where values in the West can be approximately 0.5 m lower than those in the East for the same return period event; and iii) the amount of wave dissipation that occurs due to depth-induced wave breaking over the foreshore—where attenuation is greater at locations with higher foreshore elevations.

Figure 12: Spatial variation in the probability of dike failure by wave overtopping for a) scenario 1: SS \( (P_{f1}) \) and b) scenario 2: SS + IG relative to scenario 1 \( (P_{f2} / P_{f1}) \) across the wider Dutch Wadden Sea area for dikes with identical crest heights \( z_c = 6 \)m) and dike slopes \( \cot (\alpha) = 7 \).
To demonstrate this, we examine the variations in $P_{f1}$ against physical parameters for a proxy storm with an exceedance probability of 1/3000 per year. In Figure 13a, an offshore forcing parameter ($H_{m0,\text{deep}}^2T_{m-1,0,\text{deep}}$), which is proportional to the offshore energy flux, is used to represent the combined influence of $H_{m0,\text{deep}}$ and $T_{m-1,0,\text{deep}}$. In Figure 13b, the influence of variations in $\bar{\eta}$ and the bed level at the toe ($z_{b,\text{toe}}$) are represented by $h_{\text{toe}} = \bar{\eta} - z_{b,\text{toe}}$. The calculated $P_{f1}$ shows a strong positive relationship with $H_{m0,\text{deep}}^2T_{m-1,0,\text{deep}}$ ($R^2 = 0.65$), meaning that higher forcing results in higher failure probabilities. Though the correlation with $h_{\text{toe}}$ is lower ($R^2 = 0.43$), there is a trend of increasing $P_{f1}$ with increasing $h_{\text{toe}}$. This is because larger $h_{\text{toe}}$ (lower $z_{b,\text{toe}}$) values lead to higher wave heights at the toe due to less wave breaking. Likewise, higher water levels ($\bar{\eta}$) associated with larger $h_{\text{toe}}$ values, also lead to lower freeboards which results in higher overtopping volumes. Figure 13 also highlights that dikes fronted by mudflats typically have higher $P_{f1}$-values than those with saltmarshes, as saltmarshes accrete higher bed levels which in-turn promote more SS-wave attenuation by breaking.

Figure 13: Relationship between probability of failure for scenario 1 (SS) and: a) an offshore forcing parameter; and b) the water depth at the dike toe ($h_{\text{toe}}$), across the wider Dutch Wadden Sea area. Lines indicate best-fit through the data.
In order to identify the influence of the IG waves, the probability of failure by wave overtopping in scenario 2 relative to that of scenario 1 (\(P_f^2/P_f^1\)) is assessed. Figure 12b shows that \(P_f^2/P_f^1\) ranges from 1.1 to 1.6, with an average value of 1.2. This increase in \(P_f\) is due predominantly to the increase in \(T_{m-1.0,\text{toe}}/T_{m-1.0,\text{deep}}\) (at the design point) between scenario 2 and scenario 1—represented by \(\tilde{T}_2/\tilde{T}_1\) (Figure 14), where \(\tilde{T} = T_{m-1.0,\text{toe}}/T_{m-1.0,\text{deep}}\) and subscripts 1 and 2 represent scenarios 1 and 2, respectively. Figure 14a shows a strong positive relationship between \(P_f^2/P_f^1\) and \(\tilde{T}_2/\tilde{T}_1\) \((R^2 = 0.76)\), where a factor 2 increase in spectral wave period \((\tilde{T}_2/\tilde{T}_1)\) corresponds to a 1.5 times increase in the failure probability \((P_f^2/P_f^1)\). On the other hand, the increase in wave height at the toe (due to the IG waves) between the two scenarios was negligible (0.5 to 4.5%) compared to the increase in wave period.

Figure 14: a) Relationship between the change in probability of failure due to IG waves \((P_f^2/P_f^1)\) and the increase in relative spectral wave period at the toe \((\tilde{T}_2/\tilde{T}_1)\); and b) relationship between \(\tilde{T}_2/\tilde{T}_1\) and the relative water depth \((h_{\text{toe}}/H_{m0,\text{deep}})\), across the wider Wadden area. Note that \(\tilde{T}\) is a stand-in for \(T_{m-1.0,\text{toe}}/T_{m-1.0,\text{deep}}\). Subscripts 1 and 2 refer to scenarios 1 and 2, respectively.
As $T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}}$ depends largely on the offshore wave height, water depth at the toe and foreshore slope (Equation 13), the spatial variations in $P_{f_2}/P_{f_1}$ (Figure 12b) are due to variations in local bathymetric and forcing conditions. This is further demonstrated in Figure 14b by examining the relationship between the increase in spectral wave period ($\bar{T}_2/\bar{T}_1$) and the relative water depth under proxy storm conditions ($1/3000$ per year). The relative water depth parameter, which takes into account the variations in both local water depth ($h_{\text{toe}}$) and offshore wave height ($H_{m0,\text{deep}}$), shows a clear negative relationship with $\bar{T}_2/\bar{T}_1$ ($R^2 = 0.61$). Therefore, areas with low water depths at the toe and/or large offshore waves are expected to have a greater IG-wave influence on $T_{m-1,0,\text{toe}}$. It should be noted that the increase in spectral wave period due to IG waves is also sensitive to the estimated foreshore slope (Equation 13). However, $\bar{T}_2/\bar{T}_1$ showed little correlation with the foreshore slope here ($R^2 < 0.1$) since the foreshores of the Wadden Sea can all be considered very gentle ($1:600$, on average).

4 Discussion

4.1 Modelling Approach

The combined numerical and empirical approach to wave transformation proved accurate when compared to the 2015 and 2017 storm data at Uithuizerwad, also highlighting the growth of $H_{IG}$ (Figure 7) and the associated increase in $T_{m-1,0}/T_{m-1,0,\text{deep}}$ (Figure 8) as the water depth becomes shallower. Of particular note, is the difference in $T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}}$ calculated by the phase-averaged wave model SWAN compared to measurements. While the measurement data is likely contaminated by IG waves reflected from the dike, leading to longer wave periods, there is still a gross underestimation of $T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}}$ by SWAN due to its exclusion of IG-wave dynamics (Lashley et al., 2020b). Recent works also indicate that the mismatch between the $T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}}$ predictions made by SWAN and measurements may be partially explained by its misrepresentation of the frequency-dependence of wave dissipation by vegetation (Ascencio, 2020; Jacobsen and McFall, 2019)—where the presence of vegetation significantly influences the shape of the wave spectrum. However, this topic is still under investigation. Despite this underestimation of $T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}}$, SWAN was able to accurately model SS-wave transformation over the foreshore (Figure 6). Likewise, the growth of $H_{IG}$ (Figure 7) and $T_{m-1,0}/T_{m-1,0,\text{deep}}$ (Figure 8) at the dike toe are accurately captured using Equations 12 and 13, respectively.

The probabilistic method FORM was able to compute $P_r$ within 20 to 30 iterations with a computation time of under 10 minutes per dike section. Other methods, such as Crude Monte Carlo or Numerical Integration are known to be much more computationally demanding. However, other approaches such as Adaptive Directional Importance Sampling (Den Bieman et al., 2014) may also prove to be equally suitable for this application. This short computation time is also attributed to the use
of a phase-averaged wave model (SWAN), which is roughly 100 times faster than its phase-resolving counterparts (e.g. SWASH or XBeach Nonhydrostatic) (Lashley et al., 2020b).

As the dike characteristics (crest level, slope and critical overtopping discharge) typically dominate the probabilistic analysis, their treatment as deterministic variables here allowed the analysis to focus on the influence of foreshore parameters. Furthermore, by treating the influence of the IG waves as a separate module (Section 2.2.2), calculations with and without IG waves could be easily performed. Such a modular approach allows the framework to be easily modified or adapted to varying conditions. For example, the module to calculate the actual overtopping discharge could be extended with the formulae of Lashley et al. (In press) for environments where the conditions at the structure toe are extremely shallow, or in the case of vertical seawalls rather than sloping structures. Likewise, another numerical or empirical model more suited to the specific area of application could replace the model used here for SS-wave transformation (SWAN). This makes the overall approach easily adaptable and applicable to other coastlines where IG waves may play a critical role, such as the Belgian coast (Altomare et al., 2016), Japanese coast (Mase et al., 2013) and north and south coasts of Vietnam (Nguyen et al., 2020).

4.2 Applicability of Formulae for the Actual Overtopping Discharge

In the present study, the original overtopping formula of Van Gent (1999) (Equation 14) was applied to all locations. This formula was selected because it was developed specifically for shallow foreshores considering the influence of both SS and IG waves and is considered valid for a wide range of breaker parameter ($\xi_{m-1,0}$) values. However, applying Equation 14 here—to locations with $\xi_{m-1,0} < 5$ at the design point—does not coincide with the current standard (EurOtop, 2018). In EurOtop (2018), different formulae are applied depending on $\xi_{m-1,0}$ value (Van der Meer and Bruce, 2014; Altomare et al., 2016). An analysis of the different approaches revealed the following points:

- For $\xi_{m-1,0} < 1.8$, which is typical for $\cot(\alpha) = 7$, the spectral wave period showed a considerable influence on the overtopping discharge ($q_o$) calculated using the EurOtop (2018) approach. Figure 15a shows that a 1.5 times increase in wave period (due to IG waves) ($\tilde{T}_2/\tilde{T}_1$), resulted in an order of magnitude increase in $P_f$ using the EurOtop (2018) approach. Since the EurOtop (2018) formula for $\xi_{m-1,0} < 1.8$ (Van der Meer and Bruce, 2014) was not derived for shallow foreshore conditions (with IG waves), this significant increase in the $P_f$ is likely incorrect and requires further research.

- For $1.8 < \xi_{m-1,0} < 5$, which is typical for $\cot(\alpha) = 3$, the wave period no longer influences the EurOtop (2018) calculation, as a maximum $q_o$ is reached. This is evident in Figure 15b as no clear trend between $\tilde{T}_2/\tilde{T}_1$ and $P_f$ is visible for the EurOtop (2018) calculations. In these cases, the differences between EurOtop (2018) and the original Van Gent (1999) calculations are much smaller (Figure 15b) compared to Figure 15a.
• For $\xi_{m-1,0} > 7$ or wave steepness at the toe $< 0.01$, the modified version of the Van Gent (1999) formula (Altomare et al., 2016), described in EurOtop (2018) are valid for foreshore slopes steeper than or equal to 1:250. However, the modified approach, based on an equivalent slope concept (Altomare et al., 2016), is only applicable to foreshore slopes steeper than 1:250. As the foreshore slopes of the Wadden Sea are typically gentler than 1:500, the modified formulae could not be used here. Therefore, locations meeting these criteria were excluded for the EurOtop (2018) calculations.

• The results using the original Van Gent (1999) formula were of the same order of magnitude for both dike slopes considered (Figure 15), suggesting that the formula was not very sensitive to changes in $\xi_{m-1,0}$. However, it should be noted that this formula was derived using a limited dataset with $\cot(\alpha) = 2.5$ and 4 and $\cot(m) = 100$ and 250. Therefore, future studies should verify its performance for conditions with $\cot(\alpha) > 4$ and $\cot(m) > 250$.

![Figure 15](https://doi.org/10.5194/nhess-2021-211)

Figure 15: Relationship between the change in probability of failure due to IG waves ($P_{f2}/P_{f1}$) and the increase in relative spectral wave period at the toe ($\tilde{T}_2/\tilde{T}_1$) calculated using the original Van Gent (1999) and EurOtop (2018) approaches for the actual overtopping discharge for a) $\cot(\alpha) = 7$ and b) $\cot(\alpha) = 3$. Note that cases with wave steepness at the toe $< 0.01$ were excluded from the EurOtop (2018) calculations since the equivalent slope concept could not be applied (Section 4.2).
The above findings suggest that the EurOtop (2018) approach may be incorrect for shallow foreshore conditions with gentle dike slopes (e.g. 1:7), which often have $\xi_{m-1,0} < 1.8$. The source of this uncertainty lies in the sensitivity of the formulae to $T_{m-1,0,\text{toe}}$, a parameter whose magnitude increases proportionally with the magnitude of the IG waves (Figure B.2).

4.3 Influence of IG waves on Design Parameters

The influence of IG waves may be represented as an increase in the magnitude of both design parameters ($H_{m0,\text{toe}}$ and $T_{m-1,0,\text{toe}}$), compared to a situation where the IG waves are neglected. This was demonstrated by Lashley et al. (2020b), where the relative magnitude of the IG waves ($\tilde{H}_{IG} = H_{m0,\text{IG, toe}}/H_{m0,\text{SS, toe}}$) > 0.5 and the IG waves had a notable influence on both parameters. In the present study, $\tilde{H}_{IG}$ was much lower, ranging from 0.14 to 0.35 with a mean value of 0.19 (considering proxy storm conditions with 1/3000 per year exceedance probability). As a result, the impact of the IG waves on the total wave height at the toe was negligible (0.5 to 4.5%), since $H_{m0,\text{toe}} = \sqrt{H_{m0,\text{IG, toe}}^2 + H_{m0,\text{SS, toe}}^2}$. That said, there was still a notable increase in $T_{m-1,0,\text{toe}}$. This is attributed to: i) the sensitivity of $T_{m-1,0}$ to wave energy density at low frequencies, by definition (Equation 9); and ii) the influence of the foreshore slope on the shape of the wave spectrum at the toe—where gentler foreshore slopes lead to wider surf zones and more energy transfer to lower frequencies. This is demonstrated in Figure 16 using the results of two numerical simulations (XBeach Nonhydrostatic). The increase in $T_{m-1,0,\text{toe}}$ due to IG waves is larger for the 1:500 foreshore slope than the 1:50, despite having similar $\tilde{H}_{IG}$ values (Figure 16, Table 4). Table 4 also shows that while the influence of the IG waves on $T_{m-1,0,\text{toe}}$ is noteworthy, their influence on the total wave height at the toe ($H_{m0,\text{toe}}$) is negligible.

![Figure 16: Results of XBeach numerical simulations showing wave spectra at the dike toe for a 1:50 and 1:500 foreshore slope, under the same offshore forcing conditions ($H_{m0,\text{deep}} = 1 \text{ m, } T_{m-1,0,\text{deep}} = 4.54 \text{ s and } h_{\text{toe}} = 1 \text{ m}$). Dashed vertical line indicates the frequency separating IG and SS wave motions.](https://doi.org/10.5194/nhess-2021-211)
Table 4: Results of XBeach Nonhydrostatic simulations taken at the dike toe for two different foreshore slopes under the same offshore forcing conditions ($H_{m0,deep} = 1 \text{ m}, T_{m-1,0,deep} = 4.54 \text{ s} \text{ and } h_{toe} = 1 \text{ m}$).

<table>
<thead>
<tr>
<th>$\cot(m)$</th>
<th>$H_{m0,\text{toe}}$</th>
<th>$H_{m0,\text{SS,\text{toe}}}$</th>
<th>$H_{m0,\text{IG,\text{toe}}}$</th>
<th>$H_{IG}$</th>
<th>$T_{m-1,0,\text{toe}}$ (SS only)</th>
<th>$T_{m-1,0,\text{toe}}$ (SS + IG)</th>
</tr>
</thead>
<tbody>
<tr>
<td>50</td>
<td>0.98 m</td>
<td>0.97 m</td>
<td>0.15 m</td>
<td>0.16</td>
<td>4.61 s</td>
<td>5.98 s</td>
</tr>
<tr>
<td>500</td>
<td>0.80 m</td>
<td>0.79 m</td>
<td>0.14 m</td>
<td>0.17</td>
<td>4.80 s</td>
<td>7.22 s</td>
</tr>
</tbody>
</table>

The main takeaway here is that while IG waves may have a negligible influence on the design wave height at the structure, their influence on the design wave period can be considerable and should therefore not be neglected, particularly on gentle foreshore slopes.

4.4 Influence of Saltmarsh Vegetation

Another discussion point is the influence of saltmarsh vegetation and whether its effects should be considered for very high return period events. Figure 9 suggests that safety could be significantly improved by standing saltmarsh vegetation; however, these findings must be interpreted with caution. In their analysis based on dikes with foreshores in the Wadden Sea, (Vuik et al., 2018b) included a stem-breakage model and concluded that it was very likely that almost all vegetation would break at this location under extreme forcing—resulting in $P_f$ values similar to that of a non-vegetated foreshores. This flattening and breaking of saltmarsh vegetation under storm conditions was also reported by Möller et al. (2014) who conducted large-scale flume experiments with transplanted Wadden Sea vegetation. Moreover, though the vegetation component of Equation 12 was able to capture the influence of vegetation on IG waves for the two storms considered here (Figure 6 to Figure 8), its performance for more extreme events requires further validation. This is due to the low stem height to water depth ratio ($h_{\text{toe}}/h_v = 3.3$) considered in its derivation. As a result, Equation 12 may overestimate the influence of vegetation for high return-period events with $h_{\text{toe}}/h_v \approx 13$. Thus, the influence of saltmarsh vegetation on coastal safety under extreme forcing remains an important issue for future research.

4.5 Implications for Practice

Including the effects of IG waves (scenario 2) increases the $P_f$ by up to 1.6 times, compared to scenario 1 (Figure 9 and Figure 12b). This effect is considerably smaller than that reported by Oosterlo et al. (2018), where including the IG waves increased the $P_f$ by two orders of magnitude. This suggests that the findings of Oosterlo et al. (2018) were indeed case-specific and could also be due to the (potentially inappropriate) use of empirical overtopping formulae that were not formulated specifically for situations with IG waves (Van der Meer, 2002). But what does the presence of IG waves mean for practice? In general, the reliability of the existing defences may be overestimated since IG waves are largely neglected in their assessment. By
interpolating the results of Figure 10 logarithmically, the required crest level at Uithuizerwad for a fixed target probability of failure can be determined. For a target annual failure probability of 1/1000 per year (which corresponds to the safety standard), a crest level of 6.3 m (+ NAP) is needed for scenario 1 (SS). For scenario 2 (SS + IG), the required crest level is 6.5 m. Therefore, the influence of the IG waves may be alternatively seen as an increase in the required crest level of around 0.2 m with a cost in the order of magnitude of M€1/per km (Jönkman et al., 2013). If the influence of the IG waves on the $P_f$ were one order of magnitude larger, as suggested by the EurOtop (2018) formula (Figure 15a), then the increase in the required crest level would be around 0.8 m with an order of magnitude increase in cost (M€10/per km).

This increase in $P_f$ is attributed to the growth of $T_{m-1,0,\text{toe}}$ due to the IG waves and the well-known relationship between wave overtopping and wave period, where longer waves (larger $T_{m-1,0,\text{toe}}$ values) result in more overtopping (Section 2.2.3) and, by extension, higher $P_f$ values. These findings suggest that attention should also be given to changes in wave period, and not on wave height attenuation alone, when considering the influence of shallow foreshores on safety. However, it is important to stress that this effect is highly dependent on local conditions, as $T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}}$ (Equation 13) is dependent on the offshore wave height, the water depth at the toe and the estimated foreshore slope. Therefore, it should be assessed on a case-by-case basis rather than assumed constant over a large area. This spatial variability is demonstrated in Figure 12b. Additionally, the calculated $P_{f2}$ was found to be sensitive to the uncertainty in Equations 12 and 13 (Figure 11), which are based primarily on numerical simulations since field and physical model data are lacking. Future studies should carry out experiments to further validate and improve the empirical formulations presented here and, if possible, reduce the uncertainty (scatter) in their estimates.

Figure 17: Influence of raised bed levels due to saltmarshes on the spatial variation in the probability of dike failure by wave overtopping for scenario 2: SS + IG waves ($P_{f2}$) across the wider Dutch Wadden Sea area for dikes with identical crest heights ($z_c = 6$ m) and dike slopes ($\cot(\alpha) = 7$).
Even though vegetation itself was neglected in the probabilistic analysis of the wider Wadden Sea area, findings here advocate the importance of maintaining saltmarshes; as (higher) saltmarshes trap sediment, their raised platforms attenuate more SS waves than lower mudflats, which results in lower \( P_f \) values—even when IG waves are taken into account (Figure 17). These findings support the arguments of Zhu et al. (2020) for the net positive impact of shallow foreshores on coastal safety. However, the estimated increase in safety due to the foreshore may be reduced when IG waves are included in the analysis—in particular where wave overtopping is concerned. In planning and implementing foreshore systems, it is therefore important to consider the effects of IG waves on safety as well.

5 Conclusion

In this paper insight has been given into the influence of infragravity (IG) waves on the probability of failure by wave overtopping (\( P_f \)) for dikes fronted by shallow foreshores. This was achieved by incorporating an empirical expression for the IG-wave height at the toe (Equation 12 (Lashley et al., 2020a)) and a newly derived expression for the spectral wave period at the toe (Equation 13) into a probabilistic framework (Vuik et al., 2018b). The combined numerical and empirical wave modelling approach was first validated using data collected at Uithuizerwad, at a saltmarsh in the Wadden Sea (the Netherlands), during two storms with an exceedance probability of approximately 1/5 per year. It was able to reproduce the wave transformation processes associated with shallow foreshores under extreme conditions with reasonable accuracy, namely: i) the dissipation of incident SS waves by depth-induced breaking; ii) the increase in the relative magnitude of the IG waves (\( \tilde{H}_{IG} \)); and iii) the increase in the relative spectral wave periods (\( T_{m-1,0/toe} / T_{m-1,0/deep} \)).

For the Uithuizerwad case, it has been shown that \( P_f \) increases by 1.5 times if IG waves are included. This was attributed to the increase in wave period (\( T_{m-1,0/toe} \)) and, to a lesser extent, wave height (\( H_{m0/toe} \)) at the toe, which in turn leads to larger values of wave overtopping. The probabilistic framework was then applied to several locations along the Dutch Wadden Sea for a spatial analysis. Overall, including the influence of IG waves increased the \( P_f \) across all locations with the magnitude of the increase (\( P_{f2}/P_{f1} \)) varying from 1.1 to 1.6. Findings indicate that 76% of this variation is explained by the increase in \( T_{m-1,0/toe} \), as larger values lead to an increase in overtopping discharge. This increase is due to the IG waves and is highly dependent on the water depth and offshore wave height at each location. The sensitivity of the calculated \( P_f \) to changes in \( T_{m-1,0/toe} \) (Figure 13a) compared to changes in \( H_{m0/toe} \) also suggests that wave height attenuation alone is not a sufficient metric to evaluate the influence of shallow foreshores on safety. It is recommended that future studies assess changes in both wave height and period over the foreshore, which can be estimated using the empirical formulations proposed Lashley et al. (2020a) and this paper, respectively.
The main conclusion of the present work is that IG waves can have an impact on safety. The magnitude of the impact varied considerably depending on the empirical wave overtopping formula applied—where a maximum increase of 1.6 times the $P_f$ was found with the Van Gent (1999) formula compared to an 88 times increase when EurOtop (2018) was applied. Here, the Van Gent (1999) formula was considered more appropriate since both SS and IG waves were considered in its derivation. It is thus important that coastal engineers, dike managers and decision makers consider both the impact of IG waves and the appropriateness of the wave overtopping model used when assessing flood risk along coastlines with shallow foreshores. The methods proposed in this paper can aid in this by allowing practitioners to quickly identify areas where IG waves—and therefore tools which account for them—should be included in the analysis. Furthermore, given the modular characteristic of the approach, it could be easily fitted with different tools or adapted to other coastlines where IG waves may play a significant role.

The generalisability of these results is however subject to certain limitations. For instance, the analyses were conducted assuming a fixed dike height with a constant slope at each location. Therefore, it is recommended that the analyses be repeated using the actual dike geometries where the crest level is notably higher, as findings here suggest that the influence of the IG waves would likely be higher (Figure 10a). Likewise, in this study, we have validated and applied the framework to the Dutch Wadden Sea; but it is recommended that the framework be further applied at tested for other sites where IG waves play a role. Examples of such sites include the sandy foreshores along the Belgian coast (Altomare et al., 2016), the wide shelves of the Mekong Delta, Vietnam (Nguyen et al., 2020) and the steep foreshores found in Japan (Mase et al., 2013). It must also be noted that the current study did not consider edge or leaky (free) IG waves (Reniers et al., 2021). Therefore, additional field campaigns focused on measuring IG waves would provide much needed insight into IG-wave dynamics in the Dutch Wadden Sea and provide more data to validate the tools implemented here; particularly, for more extreme conditions. It is also very important that the applicability of the existing empirical formulae for wave overtopping to conditions with IG waves be critically assessed given the sensitivity of the formulae to the presence of IG waves. Finally, while vegetation had a notable influence on wave attenuation for storms with relatively high probability of exceedance (1/5 per year, Figure 6), it was assumed to be flattened or broken under more extreme conditions (Vuik et al., 2018a). Further research is required to assess the attenuation effects of saltmarsh vegetation under extreme water level and wave forcing.

**Data availability:** All data, models, or code that support the findings of this study are available from the corresponding author upon reasonable request.

**Author contribution:** CL prepared the manuscript with contributions from all co-authors. CL designed and carried out the numerical simulations. CL performed the formal analysis. CL, JB, SJ and JvM conceptualized the research. SJ, JvM and JB provided supervision. VV and CL developed the model code (MATLAB). VV designed and carried out the field experiments.

**Competing interests:** The authors declare that they have no conflict of interest.
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Appendix A: Overview of Stochastic Variables

Table A. 1 Extreme parameters for offshore wave and water level characteristics (Weibull distributions). Note that the scale (sc) and shape (sh) parameters, derived from Hydra-NL estimates, are dependent on location along the Wadden coast; the range of values is provided here.

<table>
<thead>
<tr>
<th>Variable</th>
<th>Symbol</th>
<th>Units</th>
<th>sc</th>
<th>sh</th>
</tr>
</thead>
<tbody>
<tr>
<td>Offshore significant wave height</td>
<td>(H_{m0})</td>
<td>m</td>
<td>0.31 – 1.11</td>
<td>1.21 – 3.01</td>
</tr>
<tr>
<td>Offshore spectral wave period</td>
<td>(T_{m-1,0})</td>
<td>s</td>
<td>1.81 – 3.91</td>
<td>1.81 – 3.51</td>
</tr>
<tr>
<td>Offshore mean water level</td>
<td>(\bar{\eta})</td>
<td>m+NAP</td>
<td>2.51 – 2.71</td>
<td>2.81 – 3.31</td>
</tr>
</tbody>
</table>

Table A. 2 Normally distributed foreshore parameters. Note that the mean value (\(\mu^*\)) is dependent on location along the Wadden coast.

<table>
<thead>
<tr>
<th>Variable</th>
<th>Symbol</th>
<th>Units</th>
<th>(\mu^*)</th>
<th>(\sigma)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Foreshore bed level</td>
<td>(z_b)</td>
<td>m+NAP</td>
<td>0.2</td>
<td></td>
</tr>
<tr>
<td>Factor for relative magnitude of IG waves at toe</td>
<td>(f_{IG})</td>
<td>-</td>
<td>0.99</td>
<td>0.18</td>
</tr>
<tr>
<td>Factor for relative magnitude of spectral wave period at toe</td>
<td>(f_{Tm})</td>
<td>-</td>
<td>0.99</td>
<td>0.17</td>
</tr>
<tr>
<td>Empirical overtopping coefficients</td>
<td>(c)</td>
<td>-</td>
<td>-0.92</td>
<td>0.24</td>
</tr>
</tbody>
</table>

Appendix B: Derivation of Formula for Spectral Wave Period at Toe

Hofland et al. (2017) showed that the ratio of spectral wave period at the structure toe to its deep-water equivalent \((T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}})\) may be empirically modelled as a function of relative water depth and foreshore slope (Equations B.1 to B.3). For long-crested waves (no directional spreading):

\[
\frac{T_{m-1,0,\text{toe}}}{T_{m-1,0,\text{deep}}} - 1 = 6 \cdot \exp(-4\hat{h}) + \exp(-\hat{h}),
\]

and for cases with short-crested waves:
\[ \frac{T_{m-1,0,\text{toe,in}}}{T_{m-1,0,\text{deep}}} - 1 = 6 \cdot \exp(-6\tilde{h}) + 0.25 \cdot \exp(-0.75\tilde{h}), \]  

(B.2)

where,

\[ \tilde{h} = \frac{h_{\text{toe}}}{H_{m0,\text{deep}}} \left( \frac{\cot\alpha_{\text{fore}}}{100} \right)^{0.2}. \]  

(B.3)

However, as Equations B.1 to B.3 were based on tests with \(35 \leq \cot\alpha_{\text{fore}} \leq 250\), they tend to over- and under-estimate \(T_{m-1,0,\text{toe}}\) for steep (\(\cot\alpha_{\text{fore}} < 35\)) and very gentle slopes (\(\cot\alpha_{\text{fore}} > 250\)), respectively—with \(R^2 = 0.30\) when applied to the numerical dataset (Figure B.1).

770 This inaccuracy, particularly for very gentle slopes (\(\cot\alpha_{\text{fore}} > 250\)), has also been reported by Nguyen et al. (2020) and suggests that a new formulation is required for application to the Dutch Wadden Sea—where foreshore slopes are typically 1:500 or gentler.

Since both \(T_{m-1,0,\text{toe}}\) and \(\tilde{H}_{Ig}\) both describe the amount of energy in the IG band compared to the SS band, it stands to reason that a simple relation should exist between the two parameters. From the Lashley et al. (2020a) numerical dataset, it can be seen that \(T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}}\) increases with increasing \(\tilde{H}_{Ig}\) \(R^2 = 0.76\), but with scatter related to the foreshore slope (Figure B.2a). Based on these trends, the following relation is proposed:

\[ \frac{T_{m-1,0,\text{toe}}}{T_{m-1,0,\text{deep}}} = \begin{cases} 1.59 \cdot \tilde{H}_{Ig}^{0.69} \cdot (\cot\alpha_{\text{fore}})^{0.17} & \frac{h_{\text{toe}}}{H_{m0,\text{deep}}} \leq 1, \\ 1 & \frac{h_{\text{toe}}}{H_{m0,\text{deep}}} > 1 \end{cases}. \]  

(B.4)

where the exponents were determined empirically, by minimizing scatter. Including the foreshore slope term significantly reduces the scatter in the data \((R^2 = 0.92, \text{Figure B.2b})\) and gives a better representation for mild slopes. This is due to the influence of the foreshore slope, not only on the relative magnitude of the IG waves—as already accounted for in Equation 12—but also on the spectral shape. As the area over which shoaling occurs increases with gentler foreshore slopes, energy transfer by nonlinear (difference) triad interactions occurs over a longer duration than on steeper slopes. This causes the spectral peak to migrate to lower frequencies and results in larger values of \(T_{m-1,0,\text{toe}}\) for gentler foreshore slopes (Battjes, 2004), despite having similar \(\tilde{H}_{Ig}\)-values.
It should also be noted that for deep-water cases, where \( h_{\text{toe}} / H_{m0,\text{deep}} > 1 \), \( T_{m-1,0,\text{toe}} / T_{m-1,0,\text{deep}} \approx 1 \) and is independent of the foreshore slope and \( \tilde{H}_{IG} \) parameters (Figure B.2b). This is consistent with the findings of Lashley et al. (In press) which suggest that the foreshore’s influence only becomes significant for cases with \( h_{\text{toe}} / H_{m0,\text{deep}} \leq 1 \).

Figure B.1 Numerically modelled relative spectral wave period as a function of relative water depth and foreshore slope, following Hofland et al. (2017). Black and red lines represent Equations B.1 and B.2, respectively.
Figure B. 2 Numerically modelled relative spectral wave period as a function of a) $\tilde{H}_{IG}$ alone; and b) $\tilde{H}_{IG}$ and an additional foreshore-slope term. Solid line represents Equation 13. Dashed vertical line indicates $h_{toe}/H_{m0,deep} = 1$ and dashed horizontal line represents the deep-water limit where $T_{m-1,0,\text{toe}}/T_{m-1,0,\text{deep}} \approx 1$. 