

1 Quantification of extremal dependence in spatial
2 natural hazard footprints: independence of windstorm
3 gust speeds and its impact on aggregate losses

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Abstract

Natural hazards, such as European windstorms, have wide spread effects, causing insured losses at multiple locations throughout a continent. Multivariate extreme-value statistical models for such environmental phenomena must therefore accommodate very high dimensional spatial data, as well as correctly represent dependence in the extremes to ensure accurate estimation of these losses. Ideally one would employ a flexible model, able to characterise all forms of extremal dependence. However, such models are restricted to a few dozen dimensions, hence an a priori diagnostic approach must be used to identify the dominant form of extremal dependence. Current approaches for doing so are, however, also based on relatively low dimensional data.

Here, we present an approach for systematically exploring the dominant extremal dependence class in a very high dimensional spatial hazard field. In addition, we contribute a further, natural hazards relevant diagnostic by exploring the impact of extremal dependence misspecification on conceptual aggregate hazard loss estimation. These approaches are illustrated by application to a dataset of high dimensional historical European windstorm footprints (spatial maps of 3-day maximum gust speeds at ~ 15000 locations).

We find there is little evidence of asymptotic extremal dependency in windstorm footprints. Furthermore, empirical extremal properties and conceptual losses are shown to be well reproduced using Gaussian copulas but not by extremally-dependent models such as Gumbel copulas. It is conjectured that the lack of asymptotic dependence is a generic property of turbulent flows, which may extend to other spatially continuous hazards. These results motivate the potential of using Gaussian process (geostatistical) models for efficient simulation of hazard fields.

Key Words: Natural hazards; Windstorm footprint; Extremal dependence; Reinsurance; Copulas

1 Introduction

Multivariate statistical models are increasingly used to explore the spatial characteristics of natural hazard footprints and quantify potential aggregate losses. For example, such

36 models for European windstorms are used by academics and re/insurers to create cata-
37 logues of possible events, explore loss potentials, and benchmark synthetic events from
38 atmospheric models (Bonazzi et al. 2012; Youngman and Stephenson 2016).

39 Natural hazards, such as European windstorms, have wide spread effects, often causing
40 insured losses at many locations throughout a continent. Therefore, statistical models
41 for such hazards must accommodate very high dimensional data in order to represent the
42 full hazard domain. Moreover, since natural hazards are rare events in the tail of the
43 distribution, these statistical models must also correctly represent the dependence in the
44 extremes to ensure valid inference and, hence a realistic representation of the hazard's
45 aggregate losses.

46 When modelling multivariate extremes, variables can be described as being either
47 asymptotically dependent, where large values of the variables tend to occur simulta-
48 neously, or asymptotically independent, where the largest values rarely occur together
49 (Coles et al., 1999). As noted by Wadsworth et al. (2017), examples of modelling joint
50 extremes often assume asymptotic dependence in order to accommodate asymptotically
51 justified extreme value max-stable models, potentially leading to over-estimation of the
52 joint occurrences of extremes, if incorrect. This assumption is common in the field of
53 natural hazard research. Coles and Walshaw (1994) used a max-stable model for the
54 dependence in maximum wind speeds in different directions; Blanchet et al. (2009) to
55 model snow fall in the Swiss Alps; Huser and Davison (2013) to model extreme rainfall
56 and Bonazzi et al. (2012) to model windstorm hazard fields at pairs of locations in Eu-
57 rope. Indeed, Bonazzi et al. (2012) simply base this modelling assumption on being “in
58 line with many examples found in the literature”. Therefore, it is important to ask: how
59 valid is this assumption of asymptotic dependence? And how much of an effect might a
60 misspecification of extremal dependence have on the resulting hazard loss representation
61 in the model?

62 Two approaches for exploring, and correctly representing, extremal dependence are
63 present in the literature. These involve using either a flexible model, able to represent
64 both forms of extremal dependence, or a set of diagnostic measures to identify extremal
65 dependence class prior to fitting a model with the diagnosed form of extremal dependence.

66 There is a growing literature in the area of flexible models for extremal dependence,
67 originating from the bivariate tail model of Ledford and Tawn (1996), varying in their

68 merits and limitations. Wadsworth and Tawn (2012) developed a spatial model, involving
69 inverted max-stable and max-stable models, able to incorporate both forms of extremal
70 dependence. This model, however, requires the estimation of a large number of parame-
71 ters and is only able to transition between dependence classes at a boundary point of the
72 parameter space. Following this, Wadsworth et al. (2017) explored more flexible tran-
73 sitions between extremal dependence classes and developed a model able to represent a
74 wider variety of dependence structures, although limited to the bivariate case. Huser et al.
75 (2017) went on to develop a flexible extension of the Wadsworth et al. (2017) model using
76 Gaussian scale mixtures, in which the two asymptotic dependence regimes are smoothly
77 bridged between, and estimated from the data. As noted by Huser and Wadsworth (2018),
78 however, this model either makes the transition between dependence class at a boundary
79 point of the parameter space (as in Wadsworth and Tawn 2012), or is inflexible in its
80 representation of the asymptotic independence structure. Huser and Wadsworth (2018)
81 presents the most recent advancement, in a flexible model able to capture both extremal
82 dependence classes in a parsimonious manner, provide a smooth transition between the
83 two cases and cover a wide range of possible dependence structures, all based on a small
84 number of parameters.

85 While these models provide a great advantage in terms of flexibility and are growing
86 in their applicability to higher dimensions, none are computationally feasible for very
87 high-dimensional datasets (Huser and Wadsworth, 2018), as required for natural hazards
88 modelling over a large spatial domain. Indeed, max-stable models for asymptotic de-
89 pendence are limited in application to a few dozen variables due to the computational
90 demand of existing fitting methods (de Fondeville and Davison, 2018). Hence, as noted
91 by Huser and Wadsworth (2018), with the exception of the specific high-dimensional
92 peaks-over-threshold model of de Fondeville and Davison (2018), truly high-dimensional
93 inference for spatial extreme-value models has yet to be achieved.

94 As a result, when aiming to model very high-dimensional data, the alternative, a priori
95 identification of extremal dependence class approach must be taken, and an appropriate
96 model then selected based on this identification. For example the model of de Fondeville
97 and Davison (2018) for asymptotic dependence or a geostatistical, multivariate Gaussian
98 model for asymptotic independence.

99 A number of papers have developed and/or employed diagnostic measures to identify

100 the form of extremal dependence between variables, prior to model fitting. Ledford
 101 and Tawn (1996) and Ledford and Tawn (1997) developed a bivariate tail model in
 102 which one of the parameters, named the coefficient of tail dependence, is used within a
 103 diagnostic approach to help identify the bivariate extremal dependence class. Coles et al.
 104 (1999) introduced two extremal dependence coefficients, $\chi(p)$ and $\bar{\chi}(p)$, characterising the
 105 conditional probability of a pair of locations exceeding the same high quantile threshold
 106 $1 - p$, for which the asymptotic limit (as $p \rightarrow 0$) provides a diagnostic of bivariate
 107 extremal dependence. Bortot et al. (2000) used pairwise scatter plots and empirical
 108 estimates of $\chi(p)$ and $\bar{\chi}(p)$ to diagnose the form of extremal dependence present in a
 109 3-dimensional dataset of sea surge, wave height and wave period in south-west England.
 110 They found evidence for asymptotic independence, and hence developed a multivariate
 111 Gaussian tail model for their data, derived from the joint tail of a multivariate Gaussian
 112 distribution with margins based on univariate extreme value distributions. Similarly,
 113 Eastoe et al. (2013) applied the coefficient of tail dependence, the χ and $\bar{\chi}$ measures, and
 114 the conditional extremes model of Heffernan and Tawn (2004) to estimate the form of
 115 extremal dependence in 3 hourly sea surface elevation maxima at 15 locations, identifying,
 116 in general, asymptotic dependence. Similarly, more recently, Kereszturi et al. (2015)
 117 employed the coefficient of tail dependence and χ and $\bar{\chi}$ measures within a comprehensive
 118 theoretical framework to assess extremal dependence of North Sea storm severity along
 119 four strips of 14 locations within the North Sea.

120 In all of the above examples these diagnostic approaches are applied to a relatively
 121 small number of locations. Here we present an approach for systematically exploring
 122 the dominant form of extremal dependence within a high dimensional natural hazards
 123 dataset. Specifically, we demonstrate this approach using a large ($\sim 6,000$ events) and
 124 very high-dimensional dataset ($\sim 15,000$ locations) of climate model generated European
 125 windstorm footprints.

126 We introduce the bivariate diagnostic measures of Ledford and Tawn (1996) and Coles
 127 (2001) in the context of our approach by initially using them to explore the bivariate
 128 extremal dependence in two pairs of locations within the European domain (London-
 129 Amsterdam and London-Madrid), and subsequently present an approach for systemat-
 130 ically applying the same diagnostics throughout the high dimensional domain. We use
 131 the simple extremal dependence measures of Ledford and Tawn (1996) and Coles (2001)

132 as they are quick to compute and can therefore be calculated for many thousands of pairs
133 of locations, important when exploring high dimensional data.

134 In addition, we contribute a further diagnostic, relevant for natural hazards modelling,
135 by presenting an approach for exploring the impact of extremal dependence misspecifica-
136 tion on conceptual aggregate hazard loss estimation. We use the Gaussian and Gumbel
137 copula models, representing asymptotic independence and dependence respectively, to
138 model pairs of locations, and quantify the discrepancy between modelled and observed
139 joint conceptual losses. This approach is introduced for one central location, paired with
140 all other locations in the high dimensional domain, and then extended to systematically
141 explore the full domain. In the case where a combination of asymptotic independence
142 and dependence is identified within the domain, this diagnostic is beneficial in under-
143 standing how using a model for one form of extremal dependence, necessary due to the
144 high dimensionality of the data, effects this important natural hazards model output,
145 hence providing further justification of the selected dependence model. The approaches
146 presented in this paper could be used to explore extremal dependence and develop an
147 appropriate multivariate statistical model for any alternative high-dimensional natural
148 hazard dataset.

149 The remaining paper is organised as follows. The windstorm hazard dataset used
150 throughout this paper, is described in Section 2. In Section 3 we introduce and ap-
151 ply the extremal dependence diagnostics of Ledford and Tawn (1996) and Coles et al.
152 (1999), firstly to two pairs of locations and secondly to systematically explore the high-
153 dimensional data domain. Section 4 describes our additional, natural hazards relevant,
154 conceptual aggregate loss extremal dependence diagnostic approach. Section 5 con-
155 tributes a physical explanation for the form of extremal dependence identified in the
156 windstorm hazard fields, and finally, Section 6 concludes.

157 **2 Data**

158 The windstorm footprint data set used in this paper is the same as that used in Dawkins
159 et al. (2016) and an extended version of the data set used in Roberts et al. (2014), con-
160 sisting of the 6103 high resolution model generated windstorm footprints, for windstorm
161 events that occurred within the European domain during the 35 extended winters (Oc-

162 tober - March) 1979/80 - 2013/14 (kindly provided by J. Standen and J. F. Lockwood at
163 the Met Office).

164 The windstorm footprint is defined as the maximum three second wind gust speed
165 (in ms^{-1}) at grid points in the region 15°W to 25°E in longitude and 35°N to 70°N in
166 latitude over a 72 hour period centred on the time at which the maximum 925hPa wind
167 speed occurred over land. The 925hPa wind speed is taken from ERA-interim reanalysis
168 (Dee et al., 2011). The three second wind gust speed has a robust relationship with storm
169 damage (Klawa and Ulbrich, 2003), and is commonly used in catastrophe models for risk
170 quantification (Roberts et al., 2014). A 72 hour windstorm duration is commonly used
171 in the insurance industry (Haylock, 2011), and is thought to capture the most damaging
172 phase of the windstorm (Roberts et al., 2014).

173 These 6103 historical windstorm events have been identified using the objective track-
174 ing approach of Hodges (1995) and the associated footprints are created by dynamically
175 downscaling ERA-Interim reanalysis to a horizontal resolution of 25km using the Met Of-
176 fice unified model (MetUM). As described by Roberts et al. (2014), the wind gust speeds
177 are calculated from wind speeds in the MetUM model, based on a simple gust parame-
178 terisation $U_{gust} = U_{10m} + C\sigma$, where U_{10m} is the wind speed at 10 metre altitude and σ is
179 the standard deviation of the horizontal wind, estimated from the friction velocity using
180 the similarity relation of Panofsky et al. (1977). The roughness constant C is determined
181 from the universal turbulence spectra and is larger over rough terrain.

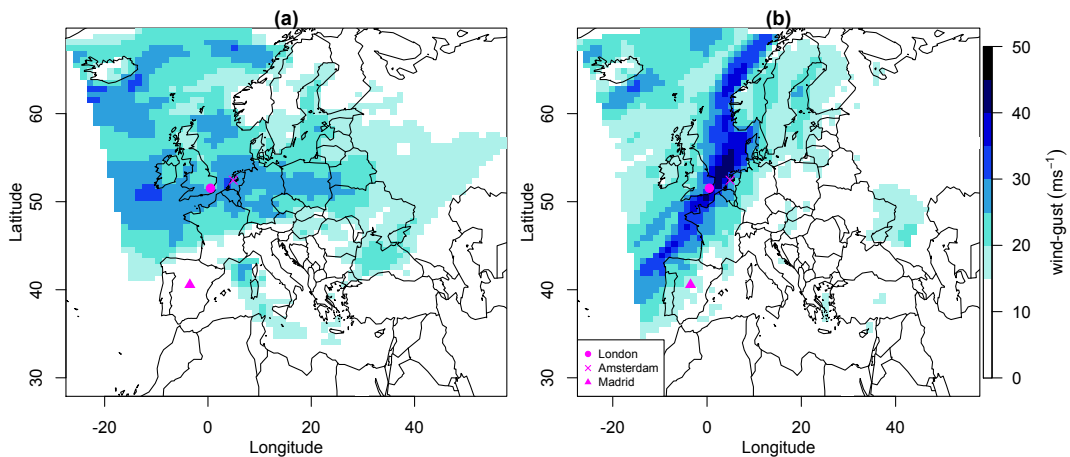


Figure 1: Hazard footprints for windstorms (a) Kyrill and (b) the Great Storm of October '87, with the location of the cities of London, Amsterdam and Madrid indicated.

182 The MetUM generated footprints for Kyrill (17th – 19th January 2007) and the Great
 183 Storm of October ‘87 (15th – 17th October 1987) are shown in Fig. 1. The variability in
 184 the intensity and location of extreme, damaging winds in these footprints demonstrate
 185 the potential importance of correctly modelling the spatial dependence between locations
 186 for realistically representing joint losses.

187 Using model generated windstorm footprints for representing historical storms has
 188 benefit in terms of spatial and temporal coverage, however these estimated maximum
 189 wind-gust speeds will inevitably differ from the those observed at nearby weather sta-
 190 tions. For example, as noted by Roberts et al. (2014), several alternative methods for
 191 parameterising wind gust speeds are available (see Sheridan (2011) for a review), which
 192 can lead to large differences in estimated gusts (10-20ms⁻¹). The validity of simplistic
 193 gust parameterisation stated above was evaluated by Roberts et al. (2014), who found
 194 an overestimation in the effect of surface roughness at stations greater than ~ 500 metre
 195 altitude, leading to an underestimation in MetUM modelled extreme winds in these lo-
 196 cations. In addition, Roberts et al. (2014) identified a slight underestimation in extreme
 197 wind gust speeds greater than $\sim 25\text{ms}^{-1}$, found to be due to a number of mechanisms
 198 including the underestimation of convective effects and strong pressure gradients, leading
 199 to the underdevelopment of fast moving storms (Roberts et al., 2014).

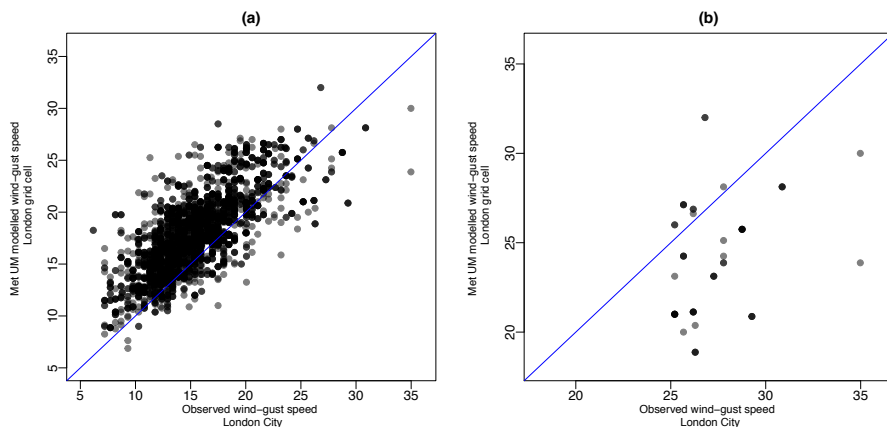


Figure 2: (a) The relationship between MetUM windstorm footprint wind gust speeds in the London grid cell and the corresponding observed wind-gust speeds at the London City weather station within the Global Summary Of the Day dataset, and (b) the same relationship for the 50 most extreme windstorm events at the London City weather station. In both plots the line $y = x$ has been added for reference (blue).

200 To explore the possible discrepancy in the MetUM wind gust speed data relevant for
201 this study, we extract daily maximum observed wind gust speed recorded at the London
202 City weather station (the station located within the London grid cell used throughout this
203 study) from the Global Summary Of the Day (GSOD) data repository, and, for each of the
204 6103 windstorm events in our dataset, find the maximum observed gust in the 3 day period
205 centred on the same date as in the MetUM model generated footprints. A comparison of
206 the observed and MetUM modelled footprint wind gusts in London is presented in Fig.
207 2 (a), indicating a general overestimation in modelled wind-gust speeds below 25m^{-1}
208 and a slight underestimation for wind-gust speeds above 25m^{-1} , reflecting the findings of
209 Roberts et al. (2014). Figure 2 (b) presents this same relationship for the 50 most extreme
210 events in the observed dataset, highlighting this underestimation of modelled extreme
211 wind-gust speeds. Indeed, the root mean squared difference between the observed and
212 modelled footprint wind-gust speeds for these 50 extreme events is 4.57ms^{-1} , giving an
213 indication of the model uncertainty in representing extreme windstorm footprint wind-
214 gust speeds.

215 The discrepancy in model generated wind-gust speeds compared to the observations
216 could lead to differences in results, namely the identification of the extremal dependence
217 class between locations. To explore this possibility we repeat the empirical analysis in
218 Section 3 (Fig. 4) for GSOD data at London City and Amsterdam Schiphol Airport,
219 shown in Figure 1 in the Supplementary Material. We find that for this pair of locations,
220 the weather station and MetUM data have very similar relationships in the extremes,
221 with the weather station data being slightly less dependent, therefore not changing the
222 conclusions of the analysis.

223 **3 Extremal dependency**

224 As a motivating example, the bivariate dependence in windstorm footprint wind gust
225 speeds for London paired with Amsterdam and Madrid are presented in Figures 3 (a)
226 and (c) respectively. These three locations are shown in Fig. 1, and these two pairings are
227 chosen because of their contrasting separation distances, and hence degrees of dependence
228 (as shown in Fig. 2 in the Supplementary Material). These scatter plots show a greater
229 degree of dependence between London and Amsterdam compared to London and Madrid.

230 Indeed, multiple windstorms have losses occurring in London and Amsterdam at the same
 231 time, when loss is associated with wind gust speeds exceeding the 99% quantile at a given
 232 location, characterised by the top right-hand corner of each plot in Fig. 3. However, does
 233 this level of dependence between London and Amsterdam necessarily suggest asymptotic
 dependence?

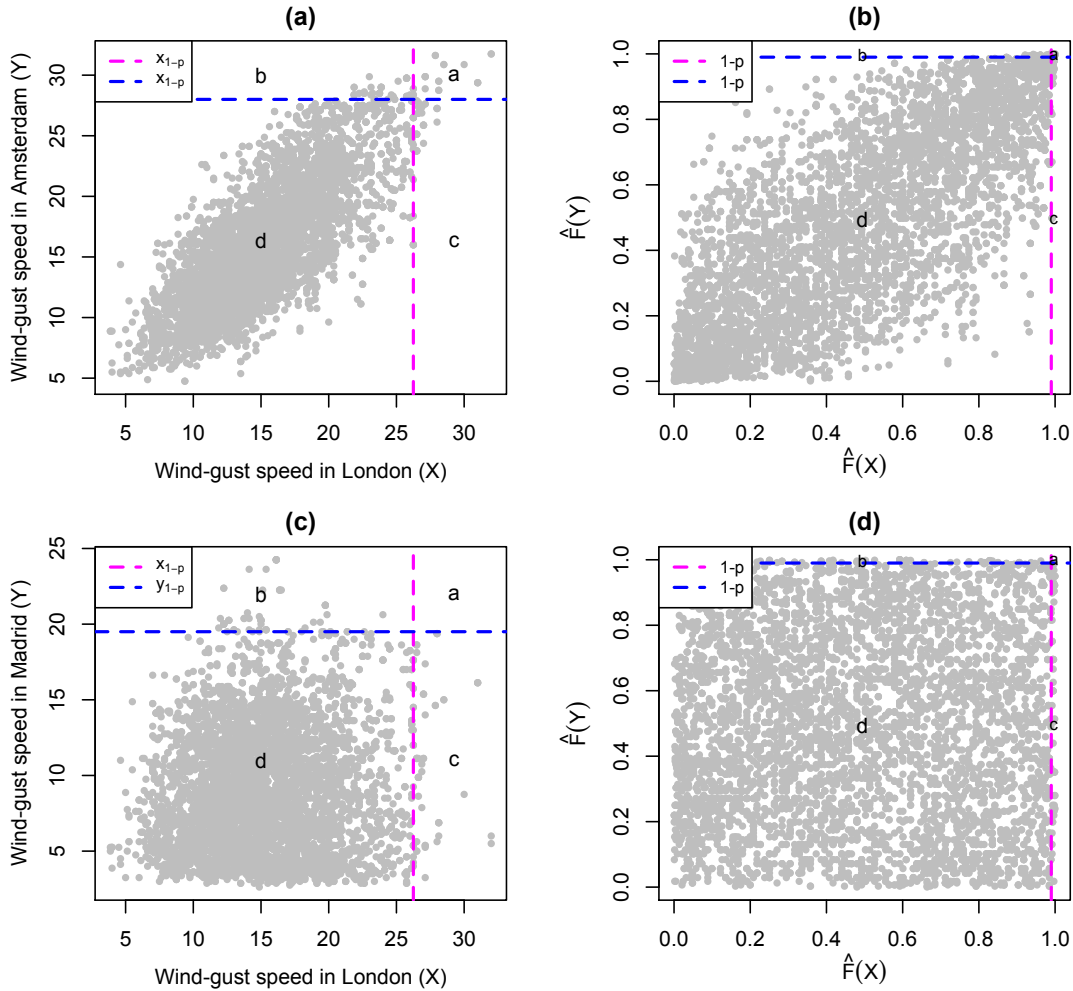


Figure 3: Scatter plot comparing historical windstorm footprint wind gust speeds (ms^{-1}) in London paired with (a) Amsterdam and (c) Madrid, and empirical copula plots for London paired with (b) Amsterdam and (d) Madrid. Dashed lines show the 99% quantile of wind gust speed at each location, and labels a-d represent the number of points in each section of each plot, related to being above or below these high quantile thresholds.

234

235 Let the $n \times 2$ variable (X, Y) represent the wind gust speeds associated with the
 236 $n = 6103$ windstorm events at any given pair of locations within the European domain,

237 e.g. London and Amsterdam. The bivariate relationship between X and Y can be
 238 represented by two components, the marginal distributions of each variable, and their
 239 joint dependence. The dependence component of the relationships shown in Fig. 3 (a)
 240 and (c) can therefore be isolated by, for each location, transforming wind gust speeds
 241 associated with each of the windstorm events, e.g. X_i for $i = 1, \dots, n$, to uniform margins
 242 using the estimator of the empirical distribution function ($\frac{1}{n} \sum_{j=1}^n \mathbb{1}_{X_j \leq X_i}$), shown in Fig.
 243 3 (b) and (d) respectively. This is known as the empirical copula.

244 3.1 Diagnostic measures

245 The degree of conditional dependence between locations, at a specified high quantile
 246 threshold, $1 - p$, can be explored, based on the empirical copula, using the Extremal De-
 247 pendence Coefficients, $\chi(p)$ and $\bar{\chi}(p)$, introduced by Coles et al. (1999), and the asymp-
 248 totic limit of these measures, as $p \rightarrow 0$, classifies the class of bivariate extremal depen-
 249 dence as either *asymptotically dependent* or *asymptotically independent*. These measures
 250 are defined as,

$$\chi(p) = \Pr(Y > y_{1-p} | X > x_{1-p}) = \frac{\Pr(Y > y_{1-p}, X > x_{1-p})}{p}, \quad (1)$$

251 where x_{1-p} and y_{1-p} are the $(1 - p)^{th}$ quantiles of X and Y respectively, $0 \leq \chi(p) < 1$
 252 for all $0 \leq (1 - p) \leq 1$, and,

$$\bar{\chi}(p) = \frac{2\log(\Pr(X > x_{1-p}))}{\log(\Pr(X > x_{1-p}, Y > y_{1-p}))} - 1 = \frac{2\log(p)}{\log(\chi(p)p)} - 1 = \frac{\log(p) - \log(\chi(p))}{\log(p) + \log(\chi(p))}, \quad (2)$$

253 where $-1 \leq \bar{\chi}(p) < 1$ for all $0 \leq (1 - p) \leq 1$. Hence, if $\lim_{p \rightarrow 0} \chi(p) = \chi(0) > 0$,
 254 $\lim_{p \rightarrow 0} \bar{\chi}(p) = \bar{\chi}(0) = 1$, and the pair (X, Y) are said to be asymptotically dependent
 255 with strength $\chi(0)$. If instead $\chi(0) = 0$, and hence, $\bar{\chi}(0) < 1$, the pair are said to be
 256 asymptotically independent, and the non-vanishing measure $\bar{\chi}(0)$ represents the strength
 257 of this non-asymptotic dependence.

258 As an initial exploration of bivariate extremal dependence class between variables,
 259 these conditional probability measures can be calculated empirically over a range of

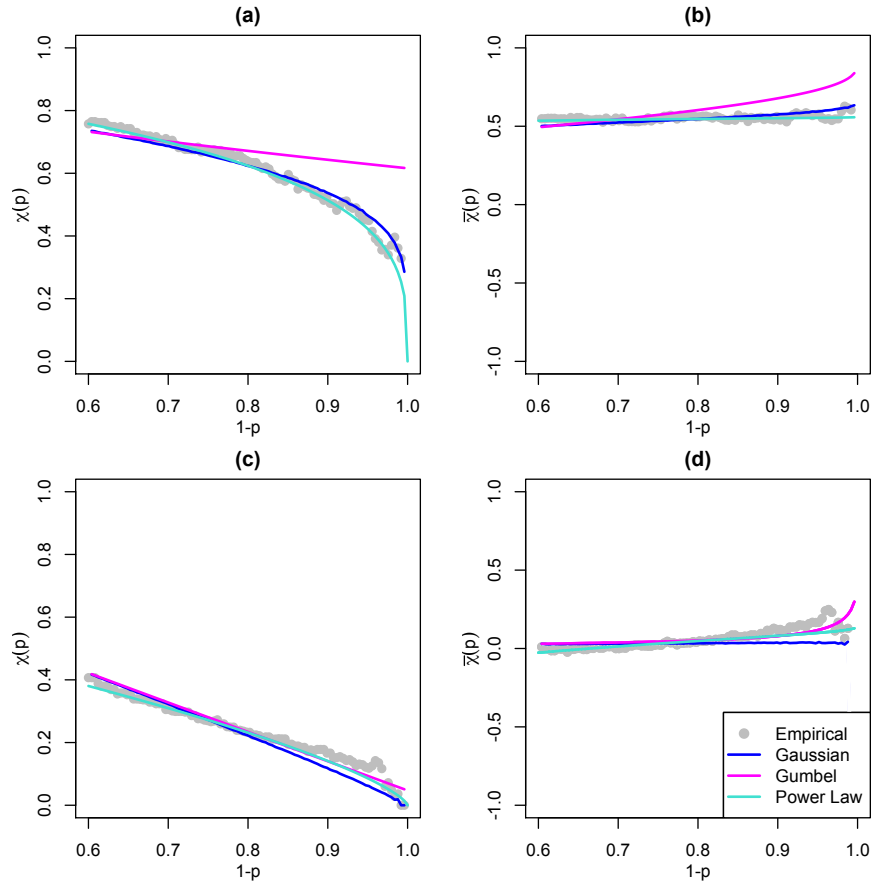


Figure 4: Extremal dependence measure $\chi(p)$, for $p \in [0, 0.4]$, for windstorm footprint wind gust speeds in London paired with (a) Amsterdam and (c) Madrid, and dependence measure $\bar{\chi}(p)$, for $p \in [0, 0.4]$, for windstorm footprint wind gust speeds in London paired with (b) Amsterdam and (d) Madrid, calculated empirically and based on the Gaussian, Gumbel and Power Law bivariate dependence functions, as defined in Table 1.

260 quantile thresholds, as shown in Fig. 4 for windstorm footprint wind gust speeds in
 261 London paired with Amsterdam and Madrid. These empirical estimates are calculated
 262 as functions of the counts (a,b,c,d) in Fig. 3, as defined in Table 1. Based on these em-
 263 pirical estimates, for both pairs of locations, $\chi(p) \rightarrow 0$ and $\bar{\chi}(p) < 1$ as $p \rightarrow 0$, suggesting
 264 asymptotic independence.

265 Here, however, and as in all datasets of environmental phenomena, the rarity of very
 266 extreme events makes it impossible to empirically quantify the asymptotic limits $\chi(0)$ and
 267 $\bar{\chi}(0)$, necessary for extremal dependence class identification. To overcome this, Ledford
 268 and Tawn (1996) developed a bivariate tail model, able to characterise both classes of
 269 extremal dependence, which when fit to a bivariate random variable can be used to predict
 270 the asymptotic limit of the conditional probability measures and hence give an estimate
 271 of the class of extremal dependence, based on the sub-asymptotic evidence in the data
 272 and the assumption that the model can be extrapolated to asymptotic levels.

273 As in Ledford and Tawn (1996), let Z_X and Z_Y denote X and Y transformed to unit
 274 Fréchet margins respectively, that is $\Pr(Z_X \leq z) = \Pr(Z_Y \leq z) = \exp(-1/z)$. Then the
 275 joint survivor function for Z_X and Z_Y , above some large quantile threshold z_{1-p} , takes
 276 the form,

$$\Pr(Z_X > z_{1-p}, Z_Y > z_{1-p}) \sim \mathcal{L}(z_{1-p})p^{1/\eta}, \quad (3)$$

277 where $p = \Pr(Z_X > z_{1-p}) = \Pr(Z_Y > z_{1-p})$, $\frac{1}{2} \leq \eta \leq 1$ is a constant and $\mathcal{L}(z_{1-p})$ is a
 278 slowly varying function as $p \rightarrow 0$. Based on this power law model, as shown by Coles
 279 et al. (1999),

$$\begin{aligned} \chi(p) &\sim \mathcal{L}(z_{1-p})p^{1/\eta-1}, \\ \bar{\chi}(p) &= \frac{2\log(p)}{\log(\mathcal{L}(z_{1-p})) + \frac{1}{\eta}\log(p)} - 1, \\ &\rightarrow 2\eta - 1 \quad \text{as } p \rightarrow 0. \end{aligned}$$

280 Hence, the parameter η , named the coefficient of tail dependence by Ledford and Tawn
 281 (1996), characterises the nature of the extremal dependence. When $\eta = 1$, $\chi(0) = \mathcal{L}(z_{1-p})$
 282 and $\bar{\chi}(0) = 1$, hence the pair (X, Y) are asymptotically dependent of degree $\mathcal{L}(z_{1-p})$.

283 Alternatively, if $\eta < 1$, $\chi(0) = 0$ and $\bar{\chi}(0) = 2\eta - 1$, and the pair are asymptotically
 284 independent with non-asymptotic dependence of degree $2\eta - 1$.

285 For a given pair, e.g. wind gust speeds in London and Amsterdam, the Ledford and
 286 Tawn (1996) model is fit to the joint survivor function along the diagonal, equivalent
 287 to the univariate distribution of $T = \min\{Z_X, Z_Y\}$, known as the structure variable,
 288 which has length n . Using the stable two parameter Poisson process representation of
 289 T , presented by Ferro (2007), who employed the Ledford and Tawn (1996) model for the
 290 verification of extreme weather forecasts, the exceedance of T above a high threshold w
 291 has the form,

$$\Pr(T > t) = \frac{1}{n} \exp \left[- \left(\frac{t - \alpha}{\eta} \right) \right] \quad \text{for all } t \geq w, \quad (4)$$

292 with location parameter α and scale parameter $0 < \eta \leq 1$, equivalent to η in Eqn. (3),
 293 estimated by maximum likelihood (Ferro, 2007).

294 We fit this model to the pairs London-Amsterdam and London-Madrid, requiring the
 295 specification of the high threshold, w , above which the Poisson process model is fit.
 296 As discussed by Ferro (2007), this threshold selection is a trade-off between being low
 297 enough that enough data is attained to ensure model precision, but high enough that the
 298 extreme-value theory that motivates the model provides accurate estimates, suggesting we
 299 should select the lowest level at which the extreme-value approximations are acceptable
 300 (Ferro, 2007). In a similar way to choosing the appropriate threshold when fitting a
 301 Generalised Pareto Distribution (see Coles 2001), empirical diagnostic plots can be used
 302 to inform this selection. For example parameter stability plots, in which the estimated
 303 model parameters and mean excess should be constant above the chosen high threshold;
 304 and quality of fit plots, in which for this model, the transformed excesses, $(Z - w)/\eta$,
 305 should be exponentially distributed if an appropriately high threshold has been chosen
 306 (see Ferro (2007) for more details).

307 Here, the 85% quantile of the structural variable T is selected, based on these diag-
 308 nostic plots (examples of these plots for London-Amsterdam are presented in Fig. 3 in
 309 the Supplementary Material). This choice is similar to the 0.88% and 0.9% thresholds
 310 selected in the applications of Ferro (2007) and Ledford and Tawn (1996) respectively.
 311 Based on this choice of w , $\eta = 0.78 < 1$ for London-Amsterdam and $\eta = 0.58 < 1$ for
 312 London-Madrid, indicating asymptotic independence for both pairs of locations. This

313 is further demonstrated in Figure 4 which shows how the Ledford and Tawn (1996)
314 model, referred to as the Power Law model, calculated as in Table 1, represents the the
315 conditional dependence measures $\chi(p)$ and $\bar{\chi}(p)$ as $p \rightarrow 0$, for London-Amsterdam and
316 London-Madrid.

317 In addition, as a comparison (included in Fig. 4), alternative parametric bivariate
318 dependence models known as the Gaussian and Gumbel copulas, can be used to model
319 the pair (X, Y) to give further indication of the extremal dependence class present.

320 The Gumbel bivariate copula model characterises asymptotic dependence with the
321 degree of dependence quantified by parameter r . For each pair of locations, this param-
322 eter is estimated via maximum likelihood using the `copula` R package. The Gaussian
323 bivariate model characterises asymptotic independence with dependence parameter ρ ,
324 here represented by the Spearman's rank correlation coefficient. Both models are fit to
325 the full bivariate data pair, as presented in Fig. 3. For the Gumbel model the data
326 is transformed to uniform margins using the empirical distribution function. The same
327 transformation is made for the Gaussian model, followed by a transformation to Gaussian
328 margins using the standard normal distribution function. The parametric forms of $\chi(p)$
329 and $\bar{\chi}(p)$ for these two opposing models are shown in Table 1. In Fig. 4, the Gumbel
330 model is calculated as in Table 1, however, since the closed form definition for the Gaus-
331 sian model in Table 1 only holds for the limit $p \rightarrow 0$, for this model $\chi(p)$ and $\bar{\chi}(p)$ are
332 estimated as the median of 1000 parametric bootstrap simulations from the associated
333 bivariate normal distribution.

334 For both pairs of locations in Fig. 4, all three parametric bivariate dependence models
335 indicate asymptotic independence, since for the Power Law model $\chi(0) = 0$ and $\bar{\chi}(0) < 1$,
336 the Gaussian model matches closely with the empirical estimates and the Power Law
337 model, and the Gumbel model overestimates the conditional probability of joint extremes.

338 As a final diagnostic, analogous to that used by Ledford and Tawn (1996, 1997),
339 the coefficient of tail dependence can be estimated for a range of high thresholds, w ,
340 to explore the sensitivity of the parameter estimate to this choice. As in Ledford and
341 Tawn (1996, 1997), here this diagnostic observes the proportion of time $\eta = 1$ is within
342 the profile likelihood confidence interval for η , when estimated using w in the interval
343 of $0.5 - 1$ quantile of T . The pair (X, Y) are said to be asymptotically dependent if
344 $\eta = 1$ is contained within these confidence intervals for a majority of the range of w ,

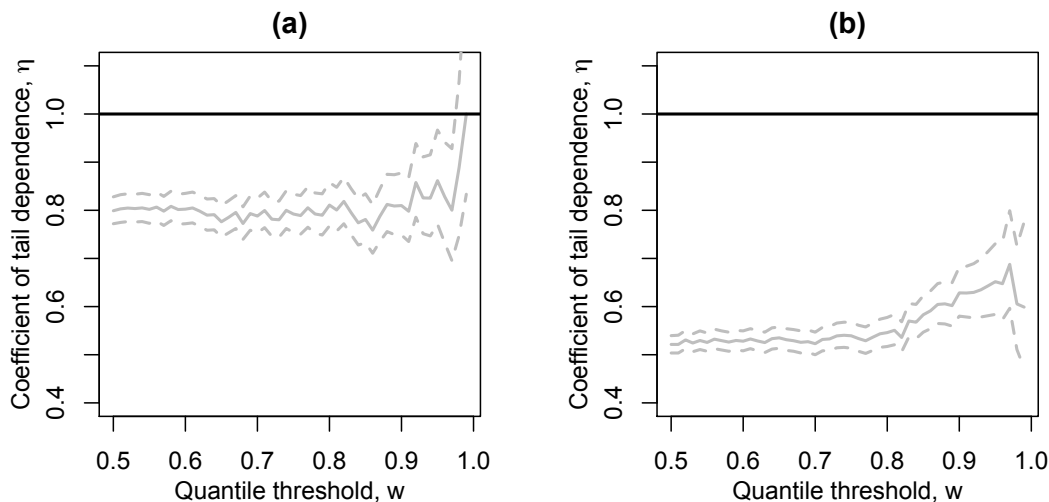


Figure 5: Diagnostic plots of maximum likelihood estimates (solid) and 95% profile likelihood confidence intervals (dashed) of η , in Eqn. (4), for threshold w in the range of the 0.5 – 1 quantile of T , for London paired with (a) Amsterdam and (b) Madrid.

345 and asymptotically independent otherwise. This exploration is presented for London
 346 paired with Amsterdam and Madrid in Fig. 5, providing further evidence of asymptotic
 347 independence for both pairs, based on this criterion.

348 3.2 Extending to high dimensions

349 We now present an approach for extending the quick-to-calculate coefficient of tail de-
 350 pendence diagnostic approach presented above to systematically explore the dominant
 351 extremal dependence class across locations in a high dimensional hazard field, demon-
 352 strated by application to our windstorm footprint data set.

353 We first take a stratified (based on the distribution of locations over longitude and lat-
 354 itude) sample of 100 locations within the European domain. One such sample is shown in
 355 Fig. 6 (a). Since the extremal dependence is likely to decrease with increasing separation
 356 distance (Wadsworth and Tawn, 2012) and we hope to understand if asymptotic indepen-
 357 dence is dominant and hence present at small separation distances, for each of these 100
 358 locations, we estimate the coefficient of tail dependence, η (and the associated 95% pro-
 359 file likelihood confidence interval) when paired with the 100 nearest locations within the
 360 full domain. Figure 6 (b) demonstrates how the 100 nearest locations are geographically

361 distributed for one such sampled location in our windstorm footprint dataset. For each
362 pairing, the coefficient of tail dependence is calculated using w as the 0.9 quantile thresh-
363 old of the structure variable, found to ensure stable estimates of η using diagnostic plots
364 as in Fig. 6 (c). The estimated η parameters and confidence intervals for these 100×100
365 pairs of locations are plotted against separation distance to explore how, throughout the
366 domain, η varies at small separation distances and changes with increasing separation
367 distance, shown in Fig. 6 (d). The parameter estimate related to the pair of locations in
368 pink and blue in Fig. 6 (b), is shown in pink. This method is repeated many times with
369 10 such repetitions shown in Fig. 4 of the Supplementary Material, showing very similar
370 results.

371 Figure 6 (d) shows that for small separation distances (<180 km) a proportion of pairs
372 of locations have coefficients of tail dependence parameter, η , estimates close to 1, with
373 $\eta = 1$ within the confidence interval, indicating asymptotic dependence. Within the range
374 (0-50 km) 69% of pairs of locations exhibit this behaviour, however this proportion reduces
375 rapidly as separation distance increases, to 30% for locations separated by (50-100 km),
376 13% for locations separated by (100-150 km) and 3% for locations separated by (150-200
377 km). Hence, while there is evidence of asymptotic dependence for some locations in close
378 proximity, even at very small separation distances (50 km) a larger proportion of locations
379 exhibit asymptotic independence. Indeed, here and in Fig. 4 of the Supplementary
380 Material, beyond a separation distance of approximately 200km the coefficient of tail
381 dependence parameter estimates drop well below 1, indicating asymptotic independence.
382 Therefore, since separation distances within the domain extend to up to 3500km, we
383 conclude that asymptotic independence is the dominant extremal dependence structure
384 across the spatial domain.

385 It is important to consider the validity of representing even this small proportion of
386 asymptotically dependent pairs of locations incorrectly as asymptotically independent.
387 To explore this, Bortot et al. (2000) carried out a simulation study in which they fit the
388 Gaussian, Ledford and Tawn (1996) and Gumbel models to bivariate data simulated from
389 three parent populations with different classes of extremal dependence. They conclude
390 that, for asymptotically independent parent populations the Gaussian copula is able to
391 provide accurate inferences for tail probability estimates, out performing the Gumbel
392 copula model, and even for asymptotically dependent parent populations, the estimation

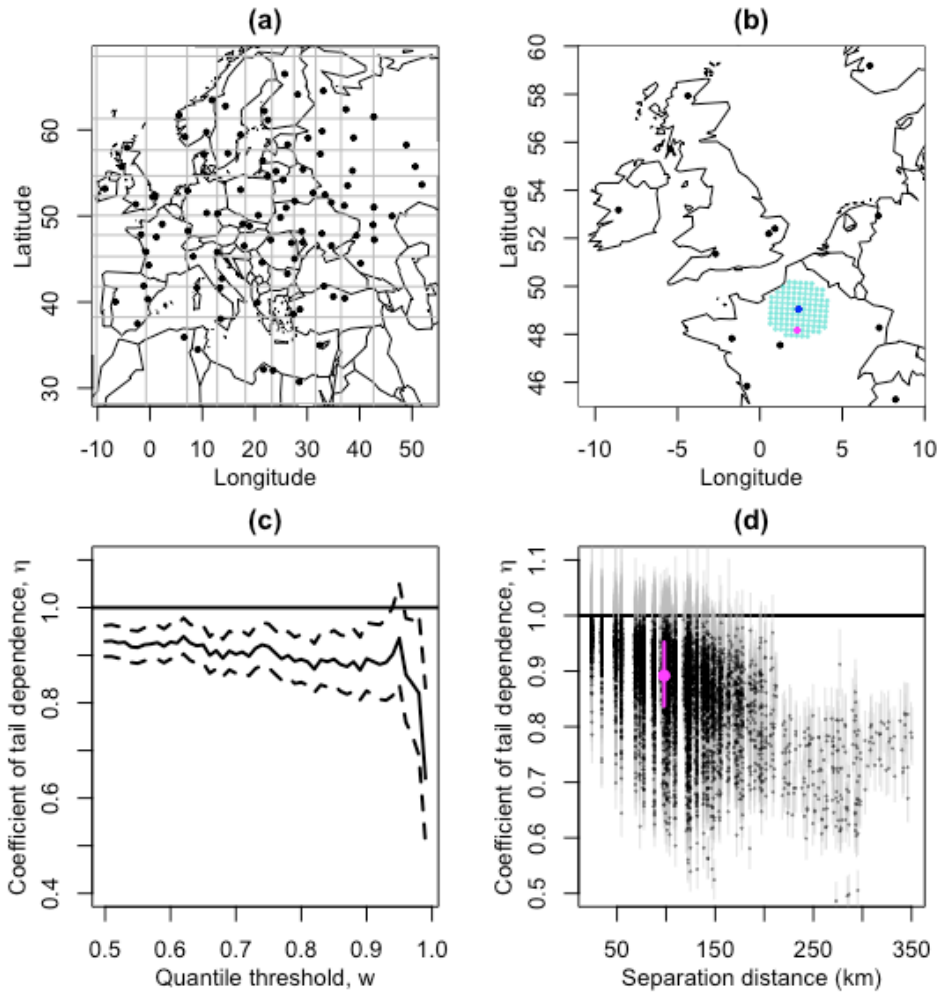


Figure 6: (a) A stratified (based on the distribution of locations over longitude and latitude) sample of locations within the European domain, with stratified grid shown in grey; (b) a demonstration of the 100 nearest locations [turquoise] to one of these sampled locations [blue], with one such point selected at random [pink]; (c) the coefficient of tail dependence diagnostic plot (as in Fig. 5) for wind gusts at the blue location paired with the pink location; (d) the coefficient of tail dependence (estimated using w as the 90% quantile threshold of the structure variable) and 95% profile likelihood confidence intervals, for each of the 100 sampled locations paired with their 100 nearest locations in the full domain, plotted against separation distance in kilometres, with the estimate based on the pair of locations in (b) and (c) added in pink.

393 error of the Gaussian copula model was deemed to be acceptably small. This suggests
 394 that, when data dimensionality prohibits the use of flexible extremal dependence mod-
 395 els, such as Huser and Wadsworth (2018), and asymptotic independence is found to be
 396 the dominant extremal dependence structure across the spatial domain, using an asymp-
 397 totically independent model, such as the Gaussian tail model, is preferable over using
 398 a model for asymptotic dependence throughout the domain. In Section 4 we present a
 399 further, natural hazards relevant, diagnostic approach for further validating this, based
 400 on estimates of the aggregate natural hazard losses.

401 **4 A conceptual loss diagnostic approach**

402 We now contribute an additional, natural hazards relevant diagnostic approach for ex-
 403 ploring the dominant extremal dependence class, providing further justification of the
 404 selected dependence model. We define a conceptual hazard loss function and explore the
 405 impact of misspecifying the extremal dependence class on aggregate hazard loss estima-
 406 tion, using the Gaussian and Gumbel copula models previously introduced. We present
 407 this approach initially based on one central location (London), and then demonstrate
 408 how this can be extended to systematically explore a high dimensional hazard field.

409 Similar to other natural hazard loss models, in the absence of confidential insurance
 410 industry exposure and vulnerability information, it has become common in the literature
 411 to define conceptual windstorm loss as a function of the footprint wind gust speeds (see
 412 Dawkins et al. (2016) for a review). While these conceptual windstorm loss functions
 413 vary in the detail of their composition, it is possible to express most in a general form,
 414 for the pair (X, Y) , as:

$$L(X, Y) = g[V(X)e(X)H\{X - U(X)\} + V(Y)e(Y)H\{Y - U(Y)\}] \quad (5)$$

415 where V is a function the wind gust speeds characterising the magnitude of the hazard, e
 416 represents exposure (e.g. population density), U quantifies a high threshold of the wind
 417 gust speed above which losses occur, H is a Heaviside function such that $H\{m\} = 1$ if
 418 $m > 0$ and $H\{m\} = 0$ otherwise, and g is an additional function applied in some cases
 419 to reduce skewness. For example, in the widely used and rigorously validated conceptual
 420 loss function of Klawa and Ulbrich (2003), $V(X) = (X - x_{0.98})^3$, $U(X) = x_{0.98}$ (where

421 $x_{0.98}$ is the 98% quantile of X) and $e(X)$ is represented by the population density at
422 the location (with equivalent expression for Y), while Cusack (2013) used a loss function
423 in which $V(X) = (X - x_{0.99})^3$, $U(X) = x_{0.99}$, the 99% quantile of X , and $g[\cdot] = \sqrt[3]{\cdot}$.
424 See Table 2.1 in Dawkins (2016) for a summary of previously published conceptual loss
425 functions in terms of the components of Eqn. 5.

426 More recently, Roberts et al. (2014) presented an exploration of the success of a num-
427 ber of these conceptual windstorm loss functions in representing insured loss throughout
428 the European domain, based on the same data set as in this study, with the aim of
429 developing a method for selecting extreme storms for the eXtreme WindStorms (XWS)
430 catalogue. While there is much published work on the existence of a relationship be-
431 tween loss severity and the magnitude of the wind, in particular the cubed excess wind
432 as used in the loss functions of Klawa and Ulbrich (2003) and Cusack (2013), Roberts
433 et al. (2014) found that a conceptual loss function representing just the area in which
434 the windstorm footprint exceeds a high loss threshold (i.e. $V(X) = 1$ and $e(X) = 1$ in
435 Eqn. 5) to be more successful at characterising a subset of extreme windstorms known
436 to have caused large insured losses. It should be noted however, that this exploration
437 did not include population density within the Klawa and Ulbrich (2003) loss function,
438 and was therefore not a direct comparison of this measure. In addition, an alternative
439 subjectively selected subset of extreme storms may have given an alternative result, and
440 the success of this simplistic ‘areal frequency of loss’ function in representing losses in
441 this climate model generated data set of windstorm footprints may be due to its relative
442 insensitivity to errors in other components of the loss estimates, such as estimated gusts,
443 and may not perform as well as other loss functions if applied to wind gust observations.

444 However, following the results of Roberts et al. (2014) in the context of this data set,
445 and in line with Dawkins et al. (2016), within this study we propose a similar threshold
446 exceedance conceptual loss function. Roberts et al. (2014) used an exceedance threshold
447 of 25ms^{-1} while Dawkins et al. (2016) used a threshold of 20ms^{-1} , as is commonly used
448 by German insurance companies (Klawa and Ulbrich, 2003). Here, similar to Klawa and
449 Ulbrich (2003) and Cusack (2013), we propose a locally varying wind gust speed quantile
450 threshold, accounting for local adaptation to varying wind intensity. We find that the
451 99% quantile of windstorm footprint wind gust speed is in excess of the commonly used
452 20ms^{-1} loss threshold for most land locations in Europe, with a higher loss threshold

453 used in regions where stronger winds occur (as shown in Figure 5 in the Supplementary
 454 Material).

455 Since, for a given storm event, $V(X)e(X)$ and $V(Y)e(Y)$ in Eqn. (5) are constants,
 456 this equation can be simplified to:

$$L(X, Y) \propto C_X H\{X - U(X)\} + C_Y H\{Y - U(Y)\} \quad (6)$$

457 where $C_X = V(X)e(X)$ and $C_Y = V(Y)e(Y)$. In our case $C_X = C_Y = 1$, and $U(X) =$
 458 $x_{0.99}$, $U(Y) = y_{0.99}$, the 99% quantiles of X and Y respectively. Therefore, while in this
 459 study we use just one conceptual loss function in which the magnitude of the loss is
 460 always equal to 1, it is simple to adapt the following analysis to accommodate alternative
 461 loss functions in which the size of the loss is included as a function of the excesses of the
 462 natural hazard, by incorporating a model for the marginal distribution of hazard at each
 463 location. This would be an interesting area of future exploration within this windstorm
 464 footprint application, beyond the scope of this analysis.

465 The probability mass function of the bivariate conceptual loss function can easily be
 466 obtained in terms of the Extremal Dependence Coefficient, $\chi(p)$, by considering the joint
 467 probability of (X, Y) in each of the quadrants shown in Fig. 3:

$$\begin{aligned} \Pr(L(X, Y) = C_X + C_Y) &= \chi(p)p, \\ \Pr(L(X, Y) = C_X) &= \Pr(L(X, Y) = C_Y) = 2(1 - \chi(p))p, \\ \Pr(L(X, Y) = 0) &= 1 + p(\chi(p) - 2), \end{aligned}$$

468 This indicates that the success of a given model in representing the bivariate con-
 469 ceptual loss for the pair (X, Y) closely relates to its characterisation of $\chi(p)$, where here
 470 $p = 0.01$, and hence the extremal dependence between X and Y .

471 To compare how well the Gaussian and Gumbel models represent our empirical bivari-
 472 ate conceptual loss function we can therefore compare estimates for $\chi(p)$ and $\bar{\chi}(p)$ for our
 473 specified loss threshold $p = 0.01$, calculated based on each model, with those calculated
 474 empirically (as in Table 1). We present the resulting difference in these estimates for
 475 London paired with all other land locations in the European domain in Fig. 7.

476 Figure 7 demonstrates how, for London paired with all other locations, the Gaussian
 477 model is able to represent empirical $\chi(0.01)$ well throughout the domain. Conversely

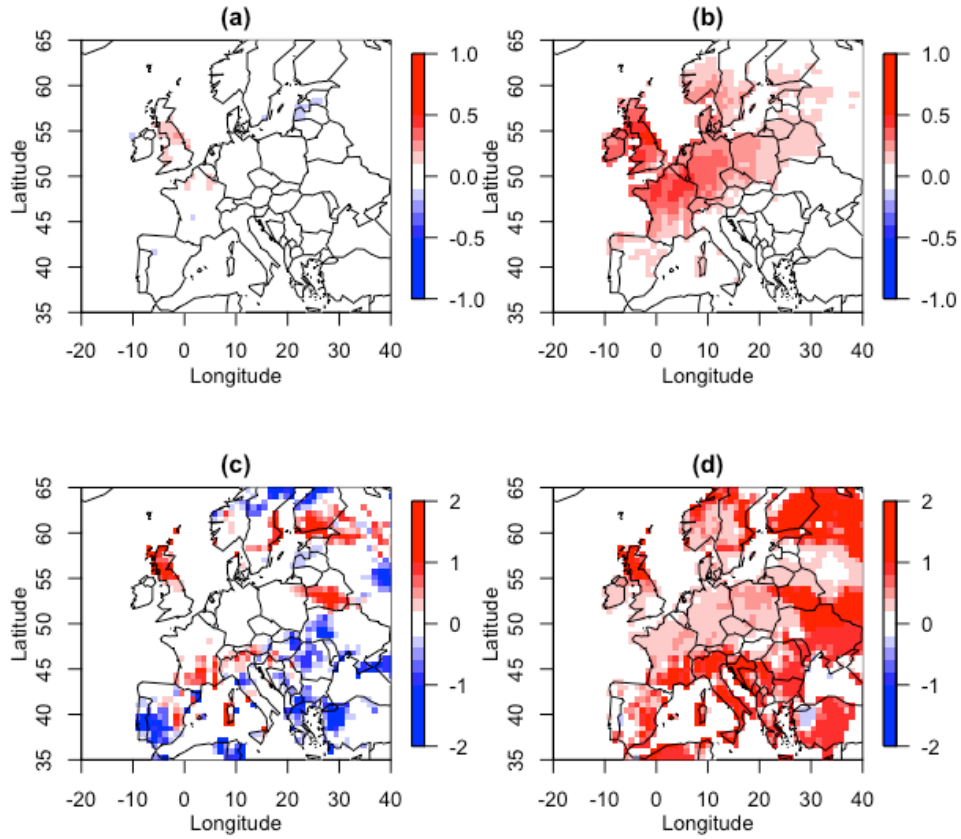


Figure 7: The difference between empirical and modelled $\chi(0.01)$ for (a) the Gaussian model and (b) the Gumbel model, and the difference between empirical and modelled $\bar{\chi}(0.01)$ for (c) the Gaussian model and (d) the Gumbel model, for London paired with all other locations over land.

478 the Gumbel model greatly over estimates $\chi(0.01)$ for all pairs of locations with non-zero
479 empirical $\chi(0.01)$, bar the neighbouring grid cell. However, this neighbouring grid cell is
480 also well represented by the Gaussian model. The Gaussian model reproduces $\bar{\chi}(0.01)$
481 well for locations within a small to medium separation distance from London, with this
482 distance being greater in the West-East direction, reflecting the common path of storms
483 over Europe (Hoskins and Hodges, 2002). The Gaussian model over and under estimates
484 $\bar{\chi}(0.01)$ for far away locations. This discrepancy is most likely due to the very small
485 sample of joint extremes at these pairs of locations making estimates of $\bar{\chi}(0.01)$ highly
486 uncertain. The Gumbel model greatly overestimates $\bar{\chi}(0.01)$, for all locations, except
487 again for those locations in very close proximity to London. This discrepancy in the
488 Gumbel model is likely due to a misspecification of asymptotic dependence between most
489 locations, resulting in an overestimation of the conditional dependencies in the extremes.

490 As well as being relevant for representing the probability mass function of the bi-
491 variate conceptual loss function, $\chi(p)$ can also be shown to characterise the conditional
492 expectation of joint loss:

$$\begin{aligned}\mathbb{E}(L(X, Y)) &= (C_X + C_Y)\chi(p)p + C_X(1 - \chi(p))p + C_Y(1 - \chi(p))p = (C_X + C_Y)p, \\ \Rightarrow \mathbb{E}(L(X, Y)|L(X) = C_X) &= (C_X + C_Y)\chi(p)p + C_X(1 - \chi(p))p = p(C_Y\chi(p) + C_X).\end{aligned}\tag{7}$$

493 The conditional first moment of the loss distribution in Eqn. (7) can therefore be
494 used to compare how well the opposing dependence models represent the size of the joint
495 losses, rather than just their conditional probability of occurrence, since the expression
496 includes C_X and C_Y . Here, $C_X = C_Y = 1$, hence the conditional expectation of joint
497 loss is equivalent to the conditional expectation of loss jointly occurring at both locations
498 given a loss has occurred at one location. It should be noted that the (non-conditional)
499 expected loss, $\mathbb{E}(L(X, Y))$, does not depend on $\chi(p)$. This is because the expectation of a
500 sum is the sum of the expectations, hence expected total loss over two or more locations
501 is simply the sum of the expected losses at each location, and so is unaffected by the
502 amount of dependency between sites.

503 Figure 8 presents a comparison of the distribution of the conditional expected joint
504 loss for London paired with each land location in our European domain, given a loss has
505 occurred in London, when calculated empirically and using the two opposing dependence

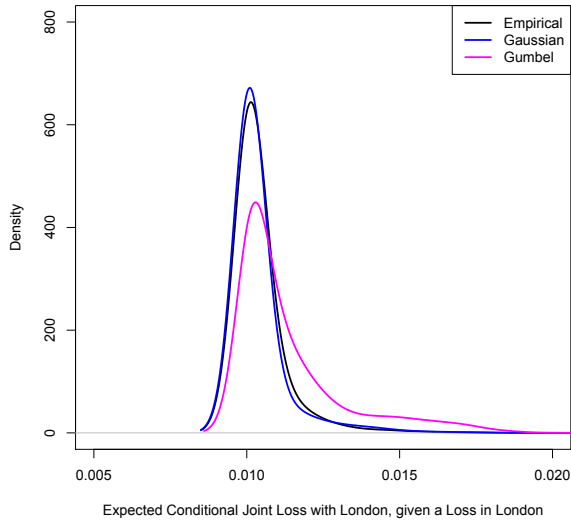


Figure 8: For all land locations in the European domain, the conditional expected joint loss with London, given a loss has occurred in London (Eqn. 7), calculated empirically and using the Gaussian and Gumbel copula models.

506 models.

507 Figure 8, further illustrates the importance of correctly specifying extremal depen-
 508 dence class when representing loss. When a conceptual loss occurs in London, the Gumbel
 509 dependence model over estimates the expected conditional joint loss with other European
 510 land locations, while conversely, the Gaussian model provides a very good estimate of the
 511 empirical expected conditional joint loss distribution.

512 4.1 Extending to high dimensions

513 We extend the analysis in Fig. 7 to systematically explore the high-dimensional domain
 514 by fitting both the Gaussian and Gumbel models to a stratified sample of 100 locations
 515 paired with each of the other 99 locations, and, for each pair, plot the difference between
 516 empirical and modelled $\chi(0.01)$ against their separation distance, shown in Fig. 9.

517 This domain-wide comparison indicates that, while the Gaussian model slightly over
 518 and under estimates empirical $\chi(0.01)$ at small separation distances, this model con-
 519 sistentlly outperforms the Gumbel model which overestimates $\chi(0.01)$ for all separation
 520 distance, even very small. This indicates, as in Fig. 6, that a majority of nearby lo-
 521 cations do not exhibit asymptotic dependence as they are not well represented by the

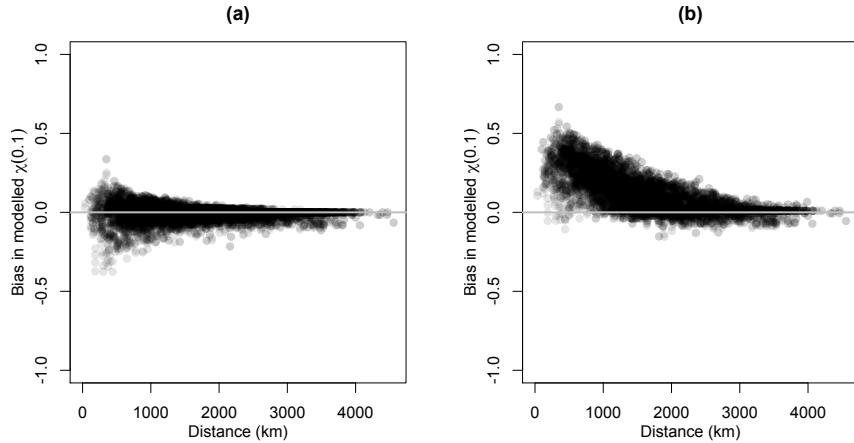


Figure 9: The difference between empirical and modelled $\chi(0.01)$ for a stratified sample of 100 locations paired with each of the other 99 locations, plotted against separation distance for (a) the Gaussian model and (b) the Gumbel model.

522 Gumbel model, further supporting the diagnosed dominance of extremal independence
 523 throughout the European domain.

524 Finally, we extend the analysis in Fig. 8 to systematically explore the high-dimensional
 525 domain by replacing London as the location of origin, with each location within a strat-
 526 ified sample of 100 locations. For each of these 100 locations, Fig. 10 presents the
 527 the difference between modelled and empirical relative frequencies of binned ranges of
 528 conditional expected joint loss, separately for the Gaussian and Gumbel models, i.e. rep-
 529 resenting the difference between the modelled and empirical density plots in Fig. 8, but
 530 for 100 locations rather than one. Fig. 10 (b) identifies that the discrepancy between the
 531 empirical and Gumbel estimates of conditional expected joint loss shown in Fig. 8 are
 532 consistent throughout the domain, with lower values being under-represented and higher
 533 values over-represented by the Gumbel model. In a similar way, Fig. 10 (a) shows that
 534 the Gaussian model performs equally well for these 100 locations, with much smaller
 535 discrepancy compared to the Gumbel model, as found in Fig. 8.

536 This novel conceptual aggregate loss diagnostic approach supports the use of the
 537 Gaussian model when asymptotic independence is found to be the dominant extremal
 538 dependence characteristic within a high dimensional natural hazards dataset. In this
 539 windstorm footprint application, we found that while the Gumbel model is able to repre-
 540 sent some pairs of locations at very small separation distances, where asymptotic depen-

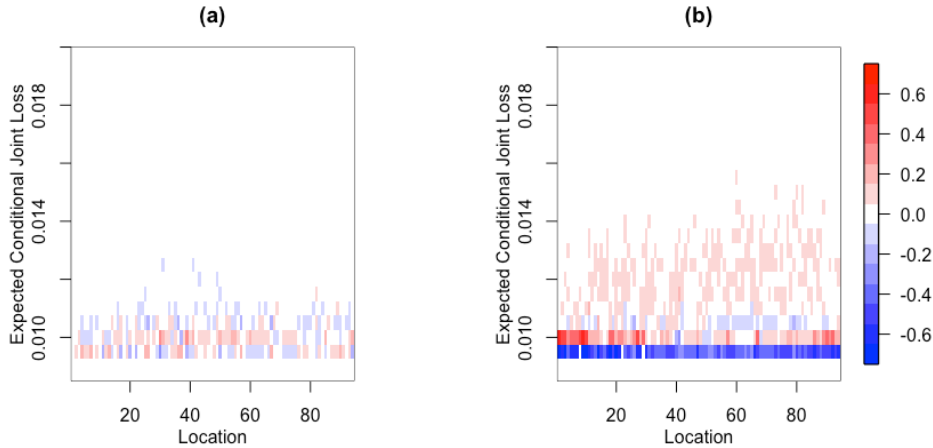


Figure 10: For a stratified sample of 100 locations within the windstorm footprint domain, the difference between modelled and empirical relative frequencies of binned ranges of expected conditional joint loss, for (a) the Gaussian model, (b) the Gumbel model.

541 dence is suggested by the coefficient of tail dependence, this model greatly misrepresents
 542 the joint tail behaviour and hence the conditional probability of joint loss for a majority
 543 of pairs and separation distances. Conversely, the Gaussian model is able to represent
 544 the joint tail behaviour relevant for loss estimation for locations within close proximity
 545 to each other, as well as further apart.

546 As previously mentioned, alternative windstorm loss thresholds have been imple-
 547 mented in other studies, for example the 98% quantile in Klawns and Ulbrich (2003),
 548 and the fixed thresholds of 20ms^{-1} in Bonazzi et al. (2012) and Dawkins et al. (2016) and
 549 25ms^{-1} in Lamb and Frydendahl (1991) and Roberts et al. (2014). An exploration of the
 550 effect of the choice of loss threshold and, indeed loss function, on how the opposing de-
 551 pendence models represent joint losses would be an extremely interesting area of further
 552 investigation, however beyond the scope of this study. Dawkins (2016) goes some way in
 553 addressing this by presenting a comparison for the 98% quantile and 25ms^{-1} fixed loss
 554 thresholds in the same form as Fig. 7. Dawkins (2016) found that the overall suitability
 555 of the opposing models remained the same for both threshold, although the discrepancy
 556 of the Gumbel model was slightly smaller for the lower, 98% quantile, loss threshold.
 557 This was thought to be because modelled exceedances further from the upper limit of the
 558 joint distribution were less sensitive to a mis-specification of the extremal dependence
 559 characteristic in the Gumbel model.

5 Why are wind gust speeds asymptotically independent?

It is of interest to ask whether there might be fundamental fluid dynamical reasons for why extreme wind gust speeds should be asymptotically independent at different spatial locations. One approach to answering this question is to consider extremal dependence in turbulent flows. The atmospheric flow in storm track regions is highly chaotic and irregular and is therefore turbulent rather than smoothly varying laminar flow (see Held 1999; and references therein). Furthermore, over short enough spatial distances, the horizontal flow in a storm may be considered to be stationary in space and directionally invariant, in other words, homogeneous isotropic turbulence.

It is useful to first consider the more tractable problem of dependency in simultaneous wind speeds rather than maximum wind speeds over a given time period. The dependency between maximum gust speeds over 3 days will not generally be less than the dependency between simultaneous wind gust speeds because maximum wind gusts for a storm do not occur at the same time at different locations. However, for locations that are close to one another, maximum gust speeds for fast moving extreme storms will occur within a short time window (e.g. within around 3 hours or less for extreme storms over the UK) and so simultaneous results become more relevant.

As originally proposed by Von Kármán (1937), turbulent wind fields can be efficiently and realistically simulated using stochastic processes (Mann, 1998). This approach is widely used for many applications such as testing loads on new aircraft designs. The basic assumption in homogeneous turbulence is that the Cartesian velocity components are independent Gaussian processes, each with a prescribed spatial covariance function. In the special case of isotropic turbulence, the spatial covariance functions are identical for each velocity component. Hence, for 2-dimensional windstorm gusts, the wind gust speed at spatial location, s , is given by $X(s) = \sqrt{u^2 + v^2}$, where $u = u(s)$ and $v = v(s)$ are independent Gaussian processes having identical covariance functions.

The distribution of each velocity component is expected, by the Central Limit Theorem, to be close to normally distributed since the net displacement of a particle in turbulence is the result of many irregular smaller displacements. The distribution of each component has zero skewness due to the symmetry of the fluid equations (negative

591 deviations are as likely as positive ones) but can have slightly more kurtosis (i.e. fatter
592 tails) than the normal distribution due to intermittency in the flow. Measurements of
593 velocity components in the atmospheric surface layer reveal that the distributions are
594 near to Gaussian (e.g. Chu et al. (1996)).

595 So what can be deduced about the extremal dependence class of wind speeds from
596 such turbulence models? Firstly, as shown in Example 5.32 of McNeil et al. (2005),
597 since the individual velocity components are bivariate normal, the individual velocity
598 components are asymptotically independent at different locations e.g. $u_1 = u(s_1)$ and
599 $u_2 = u(s_2)$ are asymptotically independent when s_1 differs from s_2 , and likewise for
600 $v(s)$. Furthermore, it can be shown that the square of each velocity component is also
601 asymptotically independent (see Appendix).

602 The squared wind speeds at pairs of locations are sums of two such independent com-
603 ponents, $(X^2, Y^2) = (u_1^2 + v_1^2, u_2^2 + v_2^2)$, and so it would be counter intuitive if somehow
604 these sums were not also asymptotically independent. Unfortunately a proof of asymp-
605 totic independence between (X^2, Y^2) (and hence (X, Y)) remains elusive. Nevertheless,
606 the conjecture can be explored using numerical simulation.

607 By simulating velocities from bivariate normal distributions, we have found no evi-
608 dence of extremal dependence in wind speeds even when each velocity component is highly
609 correlated. Figure 11 shows an example obtained by simulating a million wind speeds at
610 two locations where the u and v velocity components are independent standard normal
611 variates each with correlation of 0.9 between locations (i.e. the correlation between u_1
612 and u_2 is 0.9). The squared wind speeds at each location are chi-squared distributed
613 with 2 degrees of freedom but are not independent: there is positive association clearly
614 visible in Fig. 11(a). To assess extremal dependence, Fig. 11(b) shows how the joint
615 exceedance probability, $\Pr(X^2 > t^2, Y^2 > t^2)$, and the marginal exceedance probability,
616 $\Pr(X^2 > t^2) = \Pr(Y^2 > t^2)$, behave as threshold t^2 is varied. As the threshold is increased
617 the joint probability drops to zero faster than the marginal exceedance probability (the
618 curve in Fig. 11(b) is steeper than the dashed line), which suggests that the ratio, the
619 conditional probability of exceedance, equivalent to χ in Eqn. (1), will tend to zero in
620 the asymptotic limit.

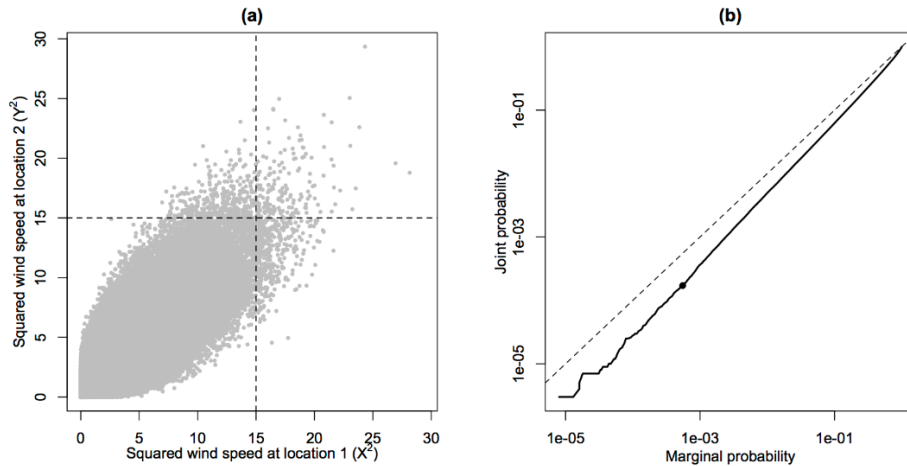


Figure 11: Simulation of wind speeds at two sites having highly correlated velocities (see main text for details): (a) scatter plot of squared wind speeds at the two sites (1000 points randomly sampled out of the million); (b) joint versus marginal exceedance probabilities (on logarithmic axes). The dot shows an example obtained by counting the fraction of points in the upper right and the right hand quadrants of (a). The curve has a steeper slope than the dashed line (equal probabilities denoting complete dependence) suggesting asymptotic independence.

6 Conclusion

This study has presented an approach for using the extremal dependence diagnostics of Coles et al. (1999) and Ledford and Tawn (1996) along the the Gaussian and Gumbel copula models to systematically explore the dominant extremal dependence class in a high dimensional natural hazards field. Within this analysis we contribute an additional, natural hazards relevant, aggregate conceptual loss extremal dependence diagnostic approach, again applied to explore extremal dependence in high dimensional spatial data. We find that when a combination of asymptotic independence and dependence is identified within the domain, this aggregate loss diagnostic is beneficial in understanding how using a model for one form of extremal dependence, necessary due to the high dimensionality of the data, effects the representation of this important natural hazards model output, hence providing further justification of the selected dependence model.

These methods reveal strong evidence of the dominance of asymptotic independence in windstorm footprint hazard fields, contrary to what has been assumed in previous studies such as Bonazzi et al. (2012), and that the mis-specification of this extremal dependency

636 (e.g. by using a Gumbel copula) leads to severe over-estimation of the probability of
637 joint losses. A reason for this lack of asymptotic dependency has been proposed based on
638 arguments from turbulence theory. These results provide justification that spatial rep-
639 resentation and simulation of windstorm hazard fields can be represented by a Gaussian
640 geostatistical model, such as that developed in Chapter 5 of Dawkins (2016).

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645 **Appendix**

Table 1: Empirical and Parametric forms for extremal dependence measures $\chi(p)$ and $\bar{\chi}(p)$.

	$\chi(p)$	$\bar{\chi}(p)$
Empirical	$\frac{a}{a+c}$	$\frac{2 \log(a+c)/n}{\log(a/n)} - 1$
Power Law	$\frac{1}{n} \exp\left(\frac{\alpha}{\eta}\right) p^{\frac{1}{\eta}-1}$	$\frac{2 \log(p)}{\log\left(\frac{1}{n} \exp\left(\frac{\alpha}{\eta}\right)\right) + \frac{1}{\eta} \log(p)} - 1$
Gumbel	$\sim 2 - \frac{(2 \log(1-p)^r)^{\frac{1}{r}}}{\log(1-p)} = 2 - 2^{\frac{1}{r}}$ (Coles et al., 1999)	$\frac{2 \log(p)}{\log(2p(1-p)^2)} - 1$
Gaussian	$\bar{F}(1-p, 1-p)/p$, where $\bar{F}(1-p, 1-p) = Pr(X > x_{1-p}, Y > y_{1-p}) \sim (1+\rho)^{\frac{3}{2}}(1-\rho)^{\frac{1}{2}}(4\pi)^{-\frac{\rho}{1+\rho}}(-\log(p))^{\frac{\rho}{1+\rho}}p^{\frac{2}{1+\rho}}$ as $p \rightarrow 0$ (Coles et al., 1999)	$\frac{2 \log(p)}{\log(\bar{F}(1-p, 1-p))} - 1$

646 **Proof of independence in stochastic models of turbulent flows**

647 Assume the velocity components (u_1, v_1) and (u_2, v_2) at two separate locations in an
 648 isotropic turbulent flow can be represented as bivariate normally distributed vectors
 649 (u_1, u_2) and (v_1, v_2) that are independent and identically distributed with zero expecta-
 650 tions.

651 The individual velocity components, (u_1, u_2) and (v_1, v_2) , are both asymptotically
 652 independent because of each being bivariate normally distributed.

653 The squares of the individual velocity components, e.g. (u_1^2, u_2^2) , are also asymptoti-
 654 cally independent. This is proven by rewriting the joint probability of exceedance:

$$\begin{aligned}
 & \Pr(u_1^2 > t^2, u_2^2 > t^2) \\
 &= \Pr(u_1 > t, u_2 > t) + \Pr(u_1 > t, u_2 \leq t) + \Pr(u_1 \leq t, u_2 > t) + \Pr(u_1 \leq t, u_2 \leq t) \\
 &= \chi_{++}\Pr(u_1 > t) + \chi_{-+}\Pr(u_1 > t) + \chi_{-+}\Pr(u_1 \leq t) + \chi_{--}\Pr(u_1 \leq t) \\
 &= \chi_{++}\Pr(u_1^2 > t^2) + \chi_{+-}\Pr(u_1^2 > t^2),
 \end{aligned}$$

655 which is obtained by noting that $\Pr(u_1^2 > t^2) = \Pr(u_1 > t) + \Pr(u_1 \leq t)$, and conditional
656 probabilities $\chi_{++} = \chi_{--}$ and $\chi_{+-} = \chi_{-+}$ by symmetry of the bivariate normal distri-
657 bution about $(0, 0)$. Since the components are bivariate normal, χ_{++} and $\chi_{+-} \rightarrow 0$ as
658 $t \rightarrow \infty$, and so $\Pr(u_1^2 > t^2, u_2^2 > t^2)/\Pr(u_1^2 > t^2) \rightarrow 0$. Hence, the square of the velocity
659 component is also asymptotically independent.

660 Perhaps rather counter-intuitively, the sum of two independent identically distributed
661 asymptotically independent variables is not necessarily asymptotically independent. It,
662 therefore, remains to be proven whether or not $(u_1^2 + v_1^2, u_2^2 + v_2^2)$ is asymptotically inde-
663 pendent.

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