



**Quantifying the effect  
of slr and flood  
defence on coastal  
flood damage**

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**Quantifying the effect of sea level rise and  
flood defence – a point process  
perspective on coastal flood damage**

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## Abstract

In contrast to recent advances in projecting sea levels, estimations about the economic impact of sea level rise are vague. Nonetheless, they are of great importance for policy making with regard to adaptation and greenhouse-gas mitigation. Since the damage is mainly caused by extreme events, we propose a stochastic framework to estimate the monetary losses from coastal floods in a confined region. For this purpose, we follow a Peak-over-Threshold approach employing a Poisson point process and the Generalised Pareto Distribution. By considering the effect of sea level rise as well as potential adaptation scenarios on the involved parameters, we are able to study the development of the annual damage. An application to the city of Copenhagen shows that a doubling of losses can be expected from a mean sea level increase of only 11 cm. In general, we find that for varying parameters the expected losses can be well approximated by one of three analytical expressions depending on the extreme value parameters. These findings reveal the complex interplay of the involved parameters and allow conclusions of fundamental relevance. For instance, we show that the damage always increases faster than the sea level rise itself. This in turn can be of great importance for the assessment of sea level rise impacts on the global scale. Our results are accompanied by an assessment of uncertainty, which reflects the stochastic nature of extreme events. While the uncertainty of flood damage increases with rising sea levels, we find that the error of our estimations in relation to the expected damage decreases.

## 1 Introduction

Considering current CO<sub>2</sub> emission pathways, severe climate change impacts need to be anticipated (IPCC, 2007; Nicholls and Cazenave, 2010). As one of the most perceivable effects of global warming, sea level rise will amplify the magnitude as well as the frequency of coastal floods (Rahmstorf and Coumou, 2011; Seneviratne et al., 2012) and is likely to have significant economic impacts (Hinkel et al., 2014; Nicholls and Tol,

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2006). Even in the case of temperature stabilisation, sea levels will continue to rise for many decades (Meehl et al., 2012). Accordingly, greenhouse-gas mitigation alone will not be sufficient, and additional preventive measures need to be considered to cope with the consequences (Petherick, 2012). The most common method to assess the efficiency of such measures is Cost-Benefit Analysis (Tol, 2002), where the benefits in terms of averted damage are compared to the investment costs. For this purpose, a concise assessment of potential economic consequences is indispensable.

Adverse effects from sea level rise are particularly expected from storm surges, presupposing the coincidence of extreme tidal and storm conditions (Woodworth et al., 2011). Accordingly, not the mean sea level itself but rather its effect on the tail of the sea level distribution needs to be studied. Since the actual distribution of sea levels is in general unknown, *extreme value theory* is commonly employed in order to characterise extreme events by a unifying tail distribution (Hawkes et al., 2008).

Estimating the *annual flood damage* at a specific site (that is the sum of all damages caused within one year), information on the occurrence of flood events, their magnitude as well as the corresponding damage is required. Due to the stochastic nature of extreme events, the annual damage cannot be predicted for a specific year and is characterised by its average value over a longer time period. In reality, the actual damage fluctuates around this *expected annual damage* with a certain variability. For instance, there are years without any damage and others where a very unlikely flood event (e.g. a 100- or 1000-year event) occurs. We measure this variability by means of the standard deviation and thus quantify the uncertainty of our damage estimations. Since environmental as well as climatic changes alter the statistics of extreme events, we assume non-stationarity and investigate the development of damage for specific parameter scenarios.

Considering sea level rise, we find analytic relations describing the damage for asymptotic parameter values (i.e. for very large changes) and show that they represent good approximations for the behaviour of damage under current conditions. Furthermore, studying the mitigation effects due to coastal protection measures in an anal-



ner et al., 2006), the excess water levels follow approximately a *Generalised Pareto Distribution* (GPD). In addition to  $u$ , the GPD is determined by a shape parameter  $\xi$  and a scale parameter  $\sigma$ . Its cumulative distribution is of the following form:

$$H_{u,\xi,\sigma}(x) = \begin{cases} 1 - \exp(-\frac{x-u}{\sigma}) & \text{if } \xi = 0 \\ 1 - (1 + \xi \frac{x-u}{\sigma})^{-1/\xi} & \text{if } \xi \neq 0 \end{cases}, \quad (1)$$

with  $x > u$ . In the case  $\xi < 0$ , the water level is bounded from above by a maximum possible water level  $x_{\max} := u - \sigma/\xi$  and we set  $H_{u,\xi,\sigma}(x) = 1$  for  $x \geq x_{\max}$  accordingly.

In our context,  $u$  is the critical water level above which damage occurs or corresponds to the given protection height at the site. However, bearing in mind that  $H_{u,\xi,\sigma}$  describes the limiting distribution of exceedances for an asymptotically increasing threshold,  $u$  needs to be large enough to obtain a good approximation of the true distribution. In particular, if no protection is given and  $u$  is freely chosen, a compromise between the adequacy of the statistical model (the larger  $u$ , the better the approximation) and the omission of smaller events ( $x < u$ ) needs to be found.

## 2.2 Point process

Section 2.1 provided the distribution of water levels given that the threshold  $u$  is exceeded. However, an estimation of annual damage requires additional information on *how often* the sea level exceeds  $u$ . Therefore, we define a *flood event* as such an exceedance and use a point process to model the incidence of these events (Coles, 2001; Embrechts et al., 1997). By employing a *Poisson process* (Reiss and Thomas, 2007), the number of flood events  $N$  within a specific year is Poisson-distributed with a certain mean value  $\Lambda$ , i.e.

$$N \sim \text{Poi}(\Lambda). \quad (2)$$

Consequently,  $\Lambda$  is the average number of flood events and  $P(N = k) = \frac{\Lambda^k}{k!} e^{-\Lambda}$  the probability of  $k$  events within one year. The Poisson property Eq. (2) is a strong assumption

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tion, which is ensured for independent water levels. However, it is commonly assumed in practice (e.g. Mudersbach and Jensen, 2010; Hawkes et al., 2008). The parameter  $\Lambda$  can be estimated by counting the number of observed events divided by the corresponding time period. Furthermore,  $\Lambda$  can be directly related to the GPD parameters: denoting  $\mu$  as the 1-year event, i.e. the water level that is exceeded on average once per year, it holds (Embrechts et al., 1997):

$$\Lambda = \begin{cases} \exp\left(-\frac{u-\mu}{\sigma}\right) & \text{if } \xi = 0 \\ \left(1 + \xi \frac{u-\mu}{\sigma-\xi(u-\mu)}\right)^{-1/\xi} & \text{if } \xi \neq 0 \end{cases}. \quad (3)$$

### 2.3 Parameter effects

We want to study the impact of sea level rise as well as potential protection measures on the flood damage. As illustrated in Fig. 2, two general effects can be observed within our framework. On the one hand, the frequency of flood events, i.e. the number of annual floods  $N_j$ , is expected to change. On the other hand, the intensities of occurring flood events can change, which would be represented in a change of the probability distribution of exceedances.

#### 2.3.1 Sea level rise

We assume that a rise in mean sea levels results in a shift of today's sea level distribution towards higher water levels without deformation of the distribution (Kauker and Langenberg, 2000; Mudersbach et al., 2013). This scenario is illustrated in Fig. 2b where the time series is shifted by the mean sea level rise leading to increased numbers of flood events  $N_j$  (which in turn change the parameter  $\Lambda = \sum_{i=1}^n N_i/n$ ) and a modified probability distribution of exceedances. Accordingly, we adjust our model in such a way that every event of certain annuality in a particular year is increased by the corresponding sea level rise (McInnes et al., 2013). This is achieved by a simple modification of the parameters. Firstly, the frequency of exceedances will increase. Using the 1-year

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where  $E_{D_i}$  and  $\text{Var}_{D_i}$  describe the expected damage of a single flood event and its variance, respectively. Please note that all event magnitudes  $D_i$  within one year are assumed to be independent and identically distributed (Coles, 2001).

## 2.6 Computational calculations

5 Given the extreme value parameters  $u, \xi, \sigma, \Lambda$  and a damage function  $F$ , we calculate the expected annual damage  $E_D$  and the standard deviation  $SD_D$  by virtue of Eqs. (7) and (8). For this purpose, the required information on the single events  $D_i$  is obtained via  $E_{D_i} = \int_u^\infty F(x) h_{u, \xi, \sigma}(x) dx$  and  $SD_{D_i}^2 = \int_u^\infty (E_{D_i} - F(x))^2 h_{u, \xi, \sigma}(x) dx$ , where  $h_{u, \xi, \sigma}(x) = \frac{d}{dx} H_{u, \xi, \sigma}(x)$  is the probability density function of exceedances. From the computational perspective, the mentioned integrals need to be discretised and the upper limit replaced by a finite value  $x_{\max}$ . In the case  $\xi < 0$ , the limit  $x_{\max}$  represents the maximal possible water level as described in Sect. 2.1, otherwise it is set to such a high value that the resulting error becomes negligible. Partitioning the range of integration  $[u, x_{\max}]$  by equidistant steps  $\Delta x$  with midpoints  $x_1, \dots, x_n$ , the following approximations are used:

$$E_{D_i} \approx \Delta x \sum_{j=1}^n F(x_j) h_{u, \xi, \sigma}(x_j) \quad \text{and} \quad (9)$$

$$SD_{D_i}^2 \approx \Delta x \sum_{j=1}^n (E_{D_i} - F(x_j))^2 h_{u, \xi, \sigma}(x_j). \quad (10)$$

## 3 Sea level rise impacts

20 We start in a general setting where the GPD parameters  $\xi, \sigma, u$  as well as the damage function exponent  $\gamma$  are given and investigate the behaviour of damage for rising sea levels. Again, we parameterise the mean sea level by the 1-year event  $\mu$ . The following

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results are derived analytically (see Appendix B for details) and hold in an asymptotic sense. More precisely, the expected damage  $E_D$  divided by the provided expression converges to a non-zero constant number for  $\mu \rightarrow \infty$  (or, if it is bounded, by a value  $\mu_{\max}$ , for  $\mu \rightarrow \mu_{\max}$ ). Hence, the following relations represent limit behaviours. Their practical use as approximations of the actual behaviour is examined in Sect. 4.

We find an increase of the annual damage by means of two separate effects (as described in Sect. 2.3.1): (i) higher frequency of events or (ii) higher severity of the events. Depending on the shape parameter  $\xi$ , three possible behaviours need to be distinguished:

(i) In the case  $\xi = 0$  (indicating an exponential tail in the sea level distribution), sea level rise leads to an exponentially increasing number of flood events while the alterations of single floods are negligible, overall implying an exponential dependence of the expected annual damage on the sea level:

$$E_D(\mu) \sim e^{\mu/\sigma}. \quad (11)$$

(ii) In contrast, we find a less steep relation if the water levels are bounded tailed (i.e.  $\xi < 0$ ):

$$E_D(\mu) \sim \mu^{\gamma-1/\xi}. \quad (12)$$

Here, the two effects are superposed: the average damage of an event increases with exponent  $\gamma$  and the number of events with exponent  $-1/\xi$ .

(iii) For the heavy-tailed case ( $\xi > 0$ ), the damage can be characterised by a power law:

$$E_D(\mu) \sim (\mu_{\max} - \mu)^{-1/\xi}, \quad (13)$$

which holds for  $\mu$  close to the maximum possible value  $\mu_{\max}$ . Approaching this value, the number of flood occurrences becomes very large and ends in a permanent flooding of the area under study for  $\mu = \mu_{\max}$ . As for  $\xi = 0$ , this behaviour

is solely caused by more frequent inundations and not by a changing severity of flood events.

It can be seen that each of the three possible relations involves a different set of parameters. Surprisingly, the damage function exponent  $\gamma$  is not involved (in the case  $\xi \geq 0$ ) or plays only a minor role (for  $\xi < 0$ , the term  $-1/\xi$  in Eq. 12 is predominant for typical parameter values), which implies that the functional behaviour of  $E_D(\mu)$  is mostly independent of the orography and the location of values in the case study area. In general, considering that  $|\xi|$  typically takes small values, the expected damage increases super-linearly in all cases.

The expected annual damage only represent average values, and the actually occurring losses fluctuate considerably. Therefore, we also examine the uncertainty of our estimations by means of the standard deviation of the damage,  $SD_D$ , and find expressions similar to the average value but with an additional factor of 0.5 in the exponents (see Appendix B for their derivations):

$$SD_D(\mu) \sim \begin{cases} e^{0.5\mu/\sigma} & \text{if } \xi = 0 \text{ (for large } \mu) \\ \mu^{\gamma-0.5/\xi} & \text{if } \xi < 0 \text{ (for large } \mu) \\ (\mu_{\max} - \mu)^{-0.5/\xi} & \text{if } \xi > 0 \text{ (for } \mu \text{ close to } \mu_{\max}) \end{cases} \quad (14)$$

This uncertainty measure represents just a lower bound since it includes only the aleatory uncertainty from the fact that one does not know *when* the extremes occur and does not take into account additional epistemic uncertainties due to a lack of knowledge, e.g. stemming from the stage–damage relation (Merz et al., 2004). However, the relations imply that although the variability increases in all cases, the relative error of the estimate,  $SD_D/E_D$ , *decreases* with rising sea levels. Surprisingly, this implies that, in a sense, flood damage becomes more foreseeable.

Besides sea level rise, which is regarded as the main driver for higher and more frequent extremes (Menéndez and Woodworth, 2010), meteorological changes can play an important role. Evolving wind patterns, for instance, can lead to a modified distribu-

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tion of water levels (Haigh et al., 2010), which in turn alters the damage distribution. Although this effect is not understood (Mudersbach and Jensen, 2010), the influence of a hypothetically changing scale parameter  $\sigma$  on the damage is studied in Appendix B.

## 4 Application

We would like to illustrate how the respective variables behave in real examples and compare our analytic derivations from the previous section with numerical calculations (as described in Sect. 2.6) for two Danish case studies – namely the city of Copenhagen and the municipality of Kalundborg. The two locations were chosen due to the availability of damage functions as well as sea level records. Details on the case studies can be found in Hallegatte et al. (2011) and Boettle et al. (2011), respectively.

For the estimation of extreme value parameters in the two case studies, extreme sea level records from closely located gauges were preprocessed by subtracting a linear trend of 0.45 cm (Copenhagen) and 0.16 cm (Kalundborg) per year (derived from mean sea level data, available at <http://www.psmsl.org>). Next, a threshold  $u$ , above which the behaviour of water levels is modelled, needs to be chosen in each case where a trade-off between bias and variance is necessary (Coles, 2001). Considering the *mean excess plots* (Fig. 3), we find that the mean excesses show an approximate linear behaviour for thresholds above 100 cm (Copenhagen) and 80 cm (Kalundborg). Choosing the thresholds correspondingly, the GPD parameters were estimated on the basis of the remaining 69 (Copenhagen) and 106 (Kalundborg) extreme sea levels. Using a maximum likelihood estimation (Embrechts et al., 1997), the parameters  $\xi = -0.14$  and  $\sigma = 15.79$  cm for Copenhagen and  $\xi = 0.08$ ,  $\sigma = 13.65$  cm for Kalundborg were obtained. Since the damage function of Kalundborg shows only a negligible damage for sea levels below 140 cm, the threshold was raised to  $u = 135$  cm, which entails a modified scale parameter of  $\sigma = 17.78$  cm according to Eq. (5). As the shape parameter  $\xi$  is negative for Copenhagen, the possible sea levels are bounded and a maximum possible sea level of  $x_{\max} = 215.28$  cm is deduced. Please note that these

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parameter deviate from the previously performed estimation of GEV parameters based on annual maximum sea level records (Boettle et al., 2013). The occurrence rates  $\Lambda$  for the description of the Poisson process were estimated by the average number of observed exceedances per year. These are  $\Lambda = 0.585$  (Copenhagen) and  $\Lambda = 0.083$  (Kalundborg). Finally, the 1-year events  $\mu = 91.21$  cm (Copenhagen) and  $\mu = 95.35$  cm (Kalundborg) were calculated by using Eq. (3).

The available damage functions support our presumption from Sect. 2.4 and follow roughly power laws with exponents  $\gamma = 1.6$  (Copenhagen) and  $\gamma = 4.1$  (Kalundborg). Further details on the damage functions can be found in Boettle et al. (2013). The exponent  $\gamma$  is used for two purposes: (i) the parameterisation of the damage function needed for the functional description of sea level rise effects (Sect. 3) and (ii) for an extrapolation of the damage function beyond the provided range which in some cases is required for numerical calculations (namely, if  $x_{\max} > 10$  m, Sect. 2.6).

Once having this information, the annual damage can be calculated numerically for varying mean sea levels (parameterised by the 1-year event  $\mu$ ). Figure 4 shows this annual mean damage and its standard deviation as a function of  $\mu$  and compares it with the asymptotic results. Since  $\xi > 0$  holds in the Kalundborg case, the parameter  $\mu$  is bounded from above by a value  $\mu_{\max}$  and we parameterise the damage by the difference  $\mu_{\max} - \mu$ . For  $\mu$  approaching  $\mu_{\max} = 332.04$  cm, the number of floods events becomes very large and ends in a permanent flooding of the considered region. It can be seen that rising mean sea levels lead to an increase in expected damage and standard deviation, which is well described by our asymptotic results already for moderate values of  $\mu$ . Accordingly, the asymptotic behaviours provide good estimates under the current conditions for both case studies. This shows that adequate projections of future flood damage can be obtained based on very few parameters. It needs to be highlighted again that in both cases the damage function exponent  $\gamma$  plays only a minor role. While for Kalundborg the approximation is even independent from it, the asymptotic projection for Copenhagen,  $\mu^{\gamma-1/\xi} \approx \mu^{1.6+7.1}$ , is clearly dominated by the shape parameter  $\xi = -0.14$ .

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In practice, one is often interested in a temporal development of damage. Therefore, mean sea level projections for the city of Copenhagen from the *Dynamical Interactive Vulnerability Assessment (DIVA)* tool (Hinkel and Klein, 2003; Vafeidis et al., 2008) have been used to study the damage as a function of time. Again, we suppose that changes in mean sea levels result in a shift of extreme events and add the estimated mean sea level rise to the corresponding events. Figure 5a displays the sea level projections for the SRES scenarios B1 (medium climate sensitivity) and A1B (high climate sensitivity) with a total rise of 11 cm and 26 cm by 2050 respectively. Panel B shows the resulting annual damage, exhibiting a steeper slope than the sea levels with an increase by a factor of roughly 2 (B1) and 4.6 (A1B) by 2050, respectively.

### 5 The effect of protection measures

At this point, it is important to bear in mind that the severity of a flood disaster is not only determined by environmental factors but also to a significant extent by human decisions (Pielke Jr and Downton, 2000). In particular, the implementation of a flood defence measure can counteract the increasing flood risk (Fig. 1c). However, identifying the appropriate height of the protection measure is crucial for choosing a cost-efficient solution, i.e. an investment that pays off within a considered time period. Therefore, we investigate the effect of varying protection heights  $\omega$  on the residual damage, assuming that inundations from flood levels below  $\omega$  are completely avoided (as suggested by Hallegatte et al., 2013). Since the distribution of sea levels is bounded in the case  $\xi < 0$ , the damage vanishes for a protection measure higher than the maximum possible water level  $x_{\max}$ . This is not the case for  $\xi \geq 0$ . In summary, we find the asymptotic relations

$$E_D(\omega) \sim \begin{cases} \omega^\gamma e^{-\omega/\sigma} & \text{if } \xi = 0 \text{ (for large } \omega) \\ (x_{\max} - \omega)^{-1/\xi} & \text{if } \xi < 0 \text{ (for } \omega \text{ close to } x_{\max}) \\ \omega^{\gamma-1/\xi} & \text{if } \xi > 0 \text{ (for large } \omega) \end{cases} \quad (15)$$

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As for rising sea levels, the behaviour of the expected damage depends fundamentally on the shape parameter  $\xi$ . While we find a decay that is dominated by an exponential component in the case  $\xi = 0$ , a power law relation independent of the scale parameter  $\sigma$  is found if  $\xi > 0$ . For  $\xi < 0$ , the expected damage follows a power law with the proximity of the protection height  $\omega$  to the maximum water level  $x_{\max}$ . Remarkably, the expressions differ not only in their functional forms but also in the parameters involved. For instance, the exponent  $\gamma$  of the damage function does not influence the behaviour in the case  $\xi < 0$  (as in Copenhagen). This highlights the decisive character of the shape parameter  $\xi$ , whose sign is not always unambiguous (Martins and Stedinger, 2000). Although a steep decrease in the damage is found in all cases, full flood safety can only be achieved if  $\xi < 0$ , and a residual risk needs to be dealt with otherwise even if potential protection failures, such as dyke breaches, are disregarded. Considering the standard deviations, similar expressions are found (again with an additional factor of 0.5 in the exponents):

$$SD_D(\omega) \sim \begin{cases} \omega^\gamma e^{-0.5\omega/\sigma} & \text{if } \xi = 0 \text{ (for large } \omega) \\ (x_{\max} - \omega)^{-0.5/\xi} & \text{if } \xi < 0 \text{ (for } \omega \text{ close to } x_{\max}) \\ \omega^{\gamma-0.5/\xi} & \text{if } \xi > 0 \text{ (for large } \omega) \end{cases} \quad (16)$$

Hence, in all cases, the relative variation of the damage,  $SD_D/E_D$ , grows with increasing protection levels. Consequently, damage in regions with high flood protection standards is subject to a wider range of relative uncertainty, indicating a higher contribution of low-probability high-impact events to the total damage (Merz et al., 2009). This shows that although coastal protection can reduce the average damage significantly, it cannot always avert the threat of very extreme floods.

Regarding the case studies Copenhagen and Kalundborg, Fig. 6 shows that the results from the numerical analyses of different protection levels can be very well approximated by our analytical relations from Eqs. (15) and (16). Accordingly, our results





intrinsic errors, it overcomes the major shortcoming of a block maxima approach and hence can be considered as superior.

Studying the effect of sea level rise, we find that in any case the expected damage increases super-linearly with the mean sea level, when considering typical values of the shape parameter. This means, the losses always increase at a higher rate than the sea levels – a universal result that needs to be explored when the climate change impacts of sea level rise are discussed economically.

Our work also shows that the upcoming losses from sea level rise are mostly determined by the type of sea level extremes (i.e. the sign of the parameter  $\xi$ ), which crucially dictates the power of  $E_D(\mu)$  (assuming constant coastal assets). This finding brings us to the following insights: (i) Since the steepness of the damage function (exponent  $\gamma$ ) is mostly irrelevant, potential policies aiming at changing the slope of the damage function via relocation of valuable assets can reduce the expected losses, but a priori have only a marginal mitigation effect on the development of future flood damage. I.e. such policies change the proportionality constant but hardly alter the proportionality. (ii) A reliable characterisation of sea level extremes is essential for a systematic assessment of climate change impacts due to sea level rise in the form of coastal floods. Thus, we plead for a high-resolution sea level network (Woodworth, 2010) if the losses from sea level rise are to be assessed on the regional or global scale.

In general, our results show how the complexity of climate change, adaptation and flood damage can be disentangled by surprisingly simple and general expressions which are applicable to arbitrary regions and case studies. These relations are the basis for understanding the effect of sea level rise on coastal flood damage and are of great importance for the development of broad scale assessment models in the context of climate change (e.g. Leimbach et al., 2010; Nordhaus and Yang, 1996).

The main text is complemented in two ways. Firstly, an additional analysis of flood damage in Copenhagen and Kalundborg as a function of the time is provided in Sect. A. Finally, all expressions describing the asymptotic behaviours of the damage and its standard deviation are mathematically proven in Sect. B.

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## Appendix A: Further results

In addition to varying the parameters  $\mu$  and  $\omega$  as discussed in the main text, the alteration of the scale parameter  $\sigma$  represents also a potential impact from climate change. Such an effect could be explained by changing wind patterns leading to a lower or higher variability of water levels. Although alterations of  $\sigma$  can be observed (Mudersbach and Jensen, 2010), the mechanism behind is still unexplained. Nevertheless, for the sake of completeness, we investigate the hypothetical effect of a varying scale parameter  $\sigma$  on the annual damage – analogous to the other parameter shifts. The corresponding asymptotic relations, Eqs. (B14) and (B15), are derived in the following section. Figure 7 illustrates the comparison of numerical calculations with the asymptotic results. It can be seen that in both case studies the increase of damage is less steep than the asymptotic behaviours for  $\sigma$  close to the present value  $\sigma_0$  and that a convergence is found for considerably larger values of  $\sigma$ .

## Appendix B: Analytical derivation of parameter effects

In this section we derive the asymptotic relations from the main text. The section comprises three parts, each part considering the effects of changing parameters  $\mu$ ,  $\sigma$  as well as  $\omega$  on one variable. First, in Sect. B1, we derive the effects of changing parameter values on the expected number of annual flood events. I.e. how many threshold exceedances of the sea level can be expected within one year? Given such an exceedance, the magnitude of the sea level is still random with a certain probability distribution, from which a probability distribution of the corresponding damage can be derived. How this distribution alters with changing parameters is described in Sect. B2. Finally, combining the number of flood events and the damage of a single flood event provides the annual damage. The derivation of the relations for the expectation value and the standard deviation of the annual damage is presented in Sect. B3.

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The provided expressions describe the damage for asymptotically large parameter values or, in case they are bounded, for parameters approaching their limit. This means that the numerically calculated values divided by the analytic result obtained converge to a non-zero constant number for increasing parameter values. In the whole section, the *Generalised Pareto* probability density function with regard to the threshold  $u$ , the shape parameter  $\xi$ , and the scale parameter  $\sigma$  is denoted by  $h_{u,\xi,\sigma}$ .

### B1 Effects on the occurrence rate

The investigation of the occurrence rate  $\Lambda$  is based on Eq. (3). It has to be noted that shifting the parameter  $\mu$  entails a modification of the scale parameter  $\sigma$  by virtue of Eq. (4). In particular, the denominator in the case  $\xi \neq 0$  is constant for varying  $\mu$ . One can see that for  $\xi > 0$  the parameter  $\mu$  is bounded from above by a value  $\mu_{\max} = \mu + \sigma/\xi$ , at which  $\Lambda$  becomes infinite. Accordingly, we study the asymptotic behaviour for  $\mu$  approaching  $\mu_{\max}$  in that case. Straightforward calculations provide the asymptotic relations for  $\Lambda$  as a function of  $\mu$ :

$$\Lambda(\mu) \sim \begin{cases} e^{\mu/\sigma} & \text{if } \xi = 0 \text{ (for large } \mu) \\ \mu^{-1/\xi} & \text{if } \xi < 0 \text{ (for large } \mu) \\ (\mu_{\max} - \mu)^{-1/\xi} & \text{if } \xi > 0 \text{ (for } \mu \text{ close to } \mu_{\max}) \end{cases}. \quad (\text{B1})$$

Considering a variable scale parameter  $\sigma$ , we obtain the relation

$$\Lambda(\sigma) \sim 1 \text{ (for large } \sigma), \quad (\text{B2})$$

which holds for arbitrary values of  $\xi$  and asymptotically large  $\sigma$ . In the last part of our analysis, we alter the protection height  $\omega$ , represented by setting the threshold to  $u = \omega$ , which also leads to a modification of the scale parameter  $\sigma$ . Using Eq. (5), it follows

asymptotically:

$$\Lambda(\omega) \sim \begin{cases} e^{-\omega/\sigma} & \text{if } \xi = 0 \text{ (for large } \omega) \\ (x_{\max} - \omega)^{-1/\xi} & \text{if } \xi < 0 \text{ (for } \omega \text{ close to } x_{\max}) \\ \omega^{-1/\xi} & \text{if } \xi > 0 \text{ (for large } \omega) \end{cases}, \quad (\text{B3})$$

where  $\omega$  is assumed to be below the maximum possible sea level  $x_{\max}$  in the case  $\xi < 0$ . If a protection level above this value is chosen, no inundation can occur and  $\Lambda$  is 0.

With these expressions, the frequency of events is fully described for large parameter values and, in combination with the following section, the behaviour of the annual damage is derived in Sect. B3. Note that Eq. (B1) holds only for small values of  $\Lambda$  and that the results are therefore only valid for the corresponding parameter values (Embrechts et al., 1997). Otherwise, if  $\Lambda$  becomes too large, the GPD is not an adequate estimation of the water levels.

## B2 Effects on the event damage distribution

Not only the number of flood events is affected by evolving parameters. As described in Sect. 2.2, in the case of a flood, its magnitude follows a GPD, which in turn is modified by changing parameters. In the following, all integrals are integrated over the whole support of the corresponding density function. For reasons of simplicity, we omit all integral limits in the text. Furthermore, we assume the shape parameter  $\xi$  to be small enough such that all integrals exist (Katz et al., 2002) – this is ensured for  $\xi < 1/\gamma$  (expectation value) and  $\xi < 0.5/\gamma$  (standard deviation). A divergence would imply an infinite variance or average value of the annual damage.

**Theorem 1** ( $\mu$  relations). *Let the water levels above a threshold  $u$  follow a GPD with parameters  $\xi$  and  $\sigma$  and let us suppose a power damage function  $F(x) = x^\gamma$  ( $\gamma \in \mathbb{R}^+$ ).*

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For the damage  $D_i$  of a single event we obtain asymptotically

$$E_{D_i}(\mu) \sim \begin{cases} 1 & \text{if } \xi = 0 \text{ (for large } \mu) \\ \mu^\gamma & \text{if } \xi < 0 \text{ (for large } \mu) \\ 1 & \text{if } \xi > 0 \text{ (for } \mu \text{ close to } \mu_{\max}) \end{cases} \quad \text{and} \quad (B4)$$

$$SD_{D_i}(\mu) \sim \begin{cases} 1 & \text{if } \xi = 0 \text{ (for large } \mu) \\ \mu^\gamma & \text{if } \xi < 0 \text{ (for large } \mu) \\ \mu_{\max} - \mu & \text{if } \xi > 0 \text{ (for } \mu \text{ close to } \mu_{\max}) \end{cases} \quad (B5)$$

with  $\mu_{\max} := \mu + \sigma/\xi$ .

- 5 *Proof.* The relations for  $\xi = 0$  follow immediately from equations provided in Sect. 2.6. In the case  $\xi < 0$ , a varying  $\mu$  leads to a linear increase in  $\sigma$  according to Eq. (4). Therefore, the relations are equivalent to  $E_{D_i}(\sigma) \sim \sigma^\gamma$  and  $SD_{D_i}(\sigma) \sim \sigma^\gamma$ . The definition of the expectation value now provides

$$E_{D_i}(\sigma)/\sigma^\gamma = \int F(z + u/\sigma) h_{0,\xi,1}(z) dz \xrightarrow{\sigma \rightarrow \infty} \int z^\gamma h_{0,\xi,1}(z) dz = \text{const.} \neq 0.$$

- 10 Furthermore, using the notation  $m_k := \int z^k h_{0,\xi,1}(z) dz$ , we obtain

$$\left( SD_{D_i}(\sigma)/\sigma^\gamma \right)^2 = \int F(z + u/\sigma)^2 h_{0,\xi,1}(z) dz - \left( \int F(z + u/\sigma) h_{0,\xi,1}(z) dz \right)^2 \\ \xrightarrow{\sigma \rightarrow \infty} m_{2\gamma} - m_\gamma^2 = \text{const.} \neq 0,$$

which shows the asymptotic relations for  $\xi < 0$ . In both cases we used uniform convergence to swap the integral and the limit.

- 15 As mentioned above,  $\mu$  is bounded by  $\mu_{\max} = \mu + \sigma/\xi$  in the case  $\xi > 0$  and we study the asymptotic behaviour for  $\mu$  approaching  $\mu_{\max}$ . Considering  $\mu \rightarrow \mu_{\max}$ , Eq. (4)

implies  $\sigma \rightarrow 0$  and by using the uniform convergence of the integrand, it follows

$$E_{D_i}(\mu) = \int F(\sigma z + u) h_{0,\xi,1}(z) dz \xrightarrow{\mu \rightarrow \mu_{\max}} F(u),$$

which shows  $E_{D_i} \sim 1$  for  $\mu$  approaching  $\mu_{\max}$ . In order to investigate the standard deviation for  $\xi > 0$ , we make use of the Taylor expansion around  $z = 0$ :

$$(\sigma z + u)^\gamma = u^\gamma + \gamma \sigma u^{\gamma-1} z + \gamma(\gamma-1) \sigma^2 u^{\gamma-2} z^2 + \mathcal{O}(\sigma^3).$$

We obtain:

$$\begin{aligned} \text{Var}_{D_i}(\sigma) &= \int (\sigma z + u)^{2\gamma} h_{0,\xi,1}(z) dz - \left( \int (\sigma z + u)^\gamma h_{0,\xi,1}(z) dz \right)^2 \\ &= \int \left( u^{2\gamma} + 2\gamma u^{2\gamma-1} \sigma z + 2\gamma(2\gamma-1) u^{2\gamma-2} \sigma^2 z^2 + \mathcal{O}(\sigma^3) \right) h_{0,\xi,1}(z) dz \\ &\quad - \left( \int \left( u^\gamma + \gamma u^{\gamma-1} \sigma z + \gamma(\gamma-1) u^{\gamma-2} \sigma^2 z^2 + \mathcal{O}(\sigma^3) \right) h_{0,\xi,1}(z) dz \right)^2 \\ &= u^{2\gamma} + 2\gamma u^{2\gamma-1} \sigma m_1 + 2\gamma(2\gamma-1) u^{2\gamma-2} \sigma^2 m_2 + \mathcal{O}(\sigma^3) \\ &\quad - \left( u^{2\gamma} + 2\gamma u^{2\gamma-1} \sigma m_1 + \gamma^2 u^{2\gamma-2} \sigma^2 m_1^2 + 2\gamma(\gamma-1) u^{2\gamma-2} \sigma^2 m_2 + \mathcal{O}(\sigma^3) \right) \\ &= \underbrace{\text{const.}}_{\neq 0} \cdot \sigma^2 + \mathcal{O}(\sigma^3) \end{aligned}$$

with the  $n$ th moments  $m_n := \int z^n h_{0,\xi,1}(z) dz$ . Considering that  $\sigma$  converges linearly to 0 for  $\mu \rightarrow \mu_{\max}$ , it follows that  $SD_{D_i}(\sigma)/\sigma \rightarrow \text{const.} \neq 0$  for  $\sigma \rightarrow 0$  and in turn  $SD_{D_i}(\mu) \sim \mu_{\max} - \mu$  for  $\mu \rightarrow \mu_{\max}$ .  $\square$

**Theorem 2** ( $\sigma$  relations). *Let the water levels above a threshold  $u$  follow a GPD with parameters  $\xi$  and  $\sigma$  and let us suppose a power damage function  $F(x) = x^\gamma$  ( $\gamma \in \mathbb{R}^+$ ). For the damage  $D_i$  of a single flood event we obtain*

$$E_{D_i}(\sigma) \sim \sigma^\gamma \text{ and } SD_{D_i}(\sigma) \sim \sigma^\gamma \tag{B6}$$

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for asymptotically large  $\sigma$ .

*Proof.* The proof corresponds to the case of increasing  $\mu$  for  $\xi < 0$  in Theorem 1.  $\square$

Finally, we derive expressions for the dependence on the protection height. As can be found in Coles (2001), a change of the threshold from  $u$  to  $u'$  affects the scale parameter  $\sigma$  by virtue of Eq. (5), which leaves the annualities of sea levels above the threshold unchanged.

**Theorem 3** ( $\omega$  relations). *Let the water levels above a threshold  $u = \omega$  follow a GPD with parameters  $\xi$  and  $\sigma$ , and let us suppose a power damage function  $F(x) = x^\gamma$  ( $\gamma \in \mathbb{R}^+$ ). For the damage  $D_i$  of a single event we obtain asymptotically*

$$E_{D_i}(\omega) \sim \begin{cases} \omega^\gamma & \text{if } \xi = 0 \text{ (for large } \omega) \\ 1 & \text{if } \xi < 0 \text{ (for } \omega \text{ close to } x_{\max}) \\ \omega^\gamma & \text{if } \xi > 0 \text{ (for large } \omega) \end{cases} \quad \text{and} \quad (\text{B7})$$

$$SD_{D_i}(\omega) \sim \begin{cases} \omega^{\gamma-1} & \text{if } \xi = 0 \text{ (for large } \omega) \\ x_{\max} - \omega & \text{if } \xi < 0 \text{ (for } \omega \text{ close to } x_{\max}) \\ \omega^\gamma & \text{if } \xi > 0 \text{ (for large } \omega) \end{cases}, \quad (\text{B8})$$

where  $x_{\max} := u - \sigma/\xi$  denotes the maximum possible sea level in the case  $\xi < 0$ .

*Proof.* Let  $u$  denote the current value of the threshold and  $\omega$  the variable protection height corresponding to the new threshold  $u'$ . We use Eq. (4) to calculate the scale parameter  $\sigma'$  which describes the excesses above the threshold  $\omega = u'$  and obtain

$$\begin{aligned} E_{D_i}(\omega) &= \int F(x) h_{\omega, \xi, \sigma}(x) dx \\ &= \int (\sigma'z + \omega)^\gamma h_{0, \xi, 1}(z) dz \\ &= \int (\sigma z + \xi z(\omega - u) + \omega)^\gamma h_{0, \xi, 1}(z) dz. \end{aligned} \quad (\text{B9})$$

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In the case  $\xi \geq 0$ , the uniform convergence of the integrand provides

$$\begin{aligned} E_{D_i}(\omega)/\omega^\gamma &= \int (\xi z + 1 + z(\sigma - \xi u)/\omega)^\gamma h_{0,\xi,1}(z) dz \\ &\xrightarrow{\omega \rightarrow \infty} \int (\xi z + 1)^\gamma h_{0,\xi,1}(z) dz \\ &= \text{const.} \neq 0, \end{aligned}$$

- 5 which shows the asymptotic relation  $E_{D_i}(\omega) \sim \omega^\gamma$  for  $\xi \geq 0$ . For the corresponding standard deviation in the case  $\xi = 0$  it follows that

$$\text{Var}_{D_i}(\omega) = \int (\sigma z + \omega)^{2\gamma} h_{0,0,1}(z) dz - \left( \int (\sigma z + \omega)^\gamma h_{0,0,1}(z) dz \right)^2.$$

Using the Taylor expansion around  $z = 0$ ,

$$(\sigma z + \omega)^\gamma = \omega^\gamma + \gamma \omega^{\gamma-1} \sigma z + \gamma(\gamma-1)/2 \omega^{\gamma-2} \sigma^2 z^2 + \mathcal{O}(\omega^{\gamma-3}),$$

- 10 we obtain

$$\begin{aligned} \text{Var}_{D_i}(\omega) &= \int \left( \omega^{2\gamma} + 2\gamma \omega^{2\gamma-1} \sigma z + \gamma(2\gamma-1) \omega^{2\gamma-2} \sigma^2 z^2 + \mathcal{O}(\omega^{2\gamma-3}) \right) h_{0,0,1}(z) dz \\ &\quad - \left( \int \left( \omega^\gamma + \gamma \omega^{\gamma-1} \sigma z + \gamma(\gamma-1)/2 \omega^{\gamma-2} \sigma^2 z^2 + \mathcal{O}(\omega^{\gamma-3}) \right) h_{0,0,1}(z) dz \right)^2 \end{aligned}$$

and straightforward calculations provide

$$\text{Var}_{D_i}(\omega) = (m_2 - m_1^2) \gamma^2 \sigma^2 \omega^{2\gamma-2} + \int \mathcal{O}(\omega^{2\gamma-3}) h_{0,0,1}(z) dz,$$

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again using  $m_k := \int z^k h_{0,0,1}(z) dz$ . Now,

$$\begin{aligned} \lim_{\omega \rightarrow \infty} \text{Var}_{D_i}(\omega) / \omega^{2\gamma-2} &= \lim_{\omega \rightarrow \infty} (m_2 - m_1^2) \gamma^2 \sigma^2 + \lim_{\omega \rightarrow \infty} \int \mathcal{O}(\omega^{-1}) h_{0,0,1}(z) dz \\ &= (m_2 - m_1^2) \gamma^2 \sigma^2 \\ &= \text{const.} \neq 0 \end{aligned}$$

5 proves the expression for  $\text{SD}_{D_i}$  in the case  $\xi = 0$ .

For  $\xi > 0$  holds

$$\begin{aligned} \text{Var}_{D_i}(\omega) / \omega^{2\gamma} &\int \frac{1}{\omega^{2\gamma}} (z(\sigma + \xi\omega - \xi u) + \omega)^{2\gamma} h_{0,\xi,1}(z) dz \\ &- \left( \int \frac{1}{\omega^\gamma} (z(\sigma + \xi\omega - \xi u) + \omega)^\gamma h_{0,\xi,1}(z) dz \right)^2 \\ &\xrightarrow{\omega \rightarrow \infty} \int (\xi z + 1)^{2\gamma} h_{0,\xi,1}(z) dz - \left( \int (\xi z + 1)^\gamma h_{0,\xi,1}(z) dz \right)^2 \\ &= \text{const.} \neq 0, \end{aligned}$$

10 which proves the asymptotic relation  $\text{SD}_{D_i}(\omega) \sim \omega^\gamma$ .

For the case  $\xi < 0$ , we consider Eq. (B9):

$$\begin{aligned} E_{D_i}(\omega) &= \int (\sigma z + \xi z(\omega - u) + \omega)^\gamma h_{0,\xi,1}(z) dz \\ &= \int (\xi z(\omega - x_{\max}) + \omega)^\gamma h_{0,\xi,1}(z) dz \\ &\xrightarrow{\omega \rightarrow x_{\max}} \int x_{\max}^\gamma h_{0,\xi,1}(z) dz = x_{\max}^\gamma, \end{aligned}$$

where we use the uniform convergence of the integrand to swap the integral and the limit. This proves the relation for  $E_{D_i}$ . In order to investigate the standard deviation, we

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define  $\Delta\omega := x_{\max} - \omega$  and examine the limit  $\Delta\omega \rightarrow 0$ . A Taylor expansion of  $(\omega - \xi\Delta\omega z)^\gamma$  around  $z = 0$  provides

$$(\omega - \xi\Delta\omega z)^\gamma = \omega^\gamma + \gamma\omega^{\gamma-1}(-\xi\Delta\omega)z + \gamma(\gamma-1)/2\omega^{\gamma-2}(-\xi\Delta\omega)^2 z^2 + \mathcal{O}(\Delta\omega^3) \quad (\text{B10})$$

and for the variance holds

$$\begin{aligned} \text{Var}_{D_i}(\omega) &= \int (\omega - \xi\Delta\omega z)^{2\gamma} h_{0,\xi,1}(z) dz - \left( \int (\omega - \xi\Delta\omega z)^\gamma h_{0,\xi,1}(z) dz \right)^2 \\ &\stackrel{\text{Eq. (B10)}}{=} \int_0^\infty (\omega^{2\gamma} + 2\gamma\omega^{2\gamma-1}(-\xi\Delta\omega)z + 2\gamma(2\gamma-1)\omega^{2\gamma-2}(-\xi\Delta\omega)^2 z^2 \\ &\quad + \mathcal{O}(\Delta\omega^3)) h_0(z; \xi, 1) dz \\ &\quad - \left( \int_0^\infty (\omega^\gamma + \gamma\omega^{\gamma-1}(-\xi\Delta\omega)z + \gamma(\gamma-1)\omega^{\gamma-2}(-\xi\Delta\omega)^2 z^2 \right. \\ &\quad \left. + \mathcal{O}(\Delta\omega^3)) h_0(z; \xi, 1) dz \right)^2 \\ &= \omega^{2\gamma} - 2\gamma\omega^{2\gamma-1}\xi\Delta\omega m_1 + 2\gamma(2\gamma-1)\omega^{2\gamma-2}\xi^2\Delta\omega^2 m_2 + \mathcal{O}(\Delta\omega^3) \\ &\quad - \left( \omega^\gamma - \gamma\omega^{\gamma-1}\xi\Delta\omega m_1 + \gamma(\gamma-1)\omega^{\gamma-2}\xi^2\Delta\omega^2 m_2 + \mathcal{O}(\Delta\omega^3) \right)^2 \\ &= \omega^{2\gamma} - 2\gamma\omega^{2\gamma-1}\xi\Delta\omega m_1 + 2\gamma(2\gamma-1)\omega^{2\gamma-2}\xi^2\Delta\omega^2 m_2 + \mathcal{O}(\Delta\omega^3) \\ &\quad - \omega^{2\gamma} + 2\gamma\omega^{2\gamma-1}\xi\Delta\omega m_1 - (2\gamma(\gamma-1)\omega^{2\gamma-2}\xi^2\Delta\omega^2 m_2 \\ &\quad - \gamma^2\omega^{2\gamma-2}\xi^2\Delta\omega^2 m_1^2) + \mathcal{O}(\Delta\omega^3) \\ &= \gamma^2\omega^{2\gamma-2}\xi^2\Delta\omega^2 m_1^2 + \mathcal{O}(\Delta\omega^3) \end{aligned}$$

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and therefore

$$SD_{D_i}(\omega)/\Delta\omega = \left( \gamma^2 \omega^{2\gamma-2} \xi^2 m_1^2 + \mathcal{O}(\Delta\omega) \right)^{1/2} \xrightarrow{\Delta\omega \rightarrow 0} \text{const.} \neq 0,$$

which shows the statement of the theorem. □

### B3 Effects on the annual damage

- 5 As stated in the main text, the total annual damage  $D$  is calculated as the sum of all single event damages  $D_i$ , i.e.  $D = D_1 + \dots + D_N$ , where  $N \sim \text{Poi}(\Lambda)$  is the number of flood events in one year. This implies  $E_N = \text{Var}_N = \Lambda$  and using Wald's identities (Beichelt, 2006), it follows

$$E_D = \Lambda E_{D_i} \text{ as well as } \text{Var}_D = \Lambda(\text{Var}_{D_i} + E_{D_i}^2) \quad (\text{B11})$$

- 10 and together with the results from Sects. B1 and B2 we obtain

$$E_D(\mu) \sim \begin{cases} e^{\mu/\sigma} & \text{if } \xi = 0 \text{ (for large } \mu) \\ \mu^{\gamma-1/\xi} & \text{if } \xi < 0 \text{ (for large } \mu) \\ (\mu_{\max} - \mu)^{-1/\xi} & \text{if } \xi > 0 \text{ (for } \mu \text{ close to } \mu_{\max}) \end{cases} \text{ as well as} \quad (\text{B12})$$

$$SD_D(\mu) \sim \begin{cases} e^{0.5\mu/\sigma} & \text{if } \xi = 0 \text{ (for large } \mu) \\ \mu^{\gamma-0.5/\xi} & \text{if } \xi < 0 \text{ (for large } \mu) \\ (\mu_{\max} - \mu)^{-0.5/\xi} & \text{if } \xi > 0 \text{ (for } \mu \text{ close to } \mu_{\max}) \end{cases} \quad (\text{B13})$$

as asymptotic relations for varying mean sea levels. An altering scale parameter  $\sigma$  leads to

$$15 \quad E_D(\sigma) \sim \sigma^\gamma \quad \text{and} \quad (\text{B14})$$

$$SD_D(\sigma) \sim \sigma^\gamma \quad (\text{B15})$$

for asymptotically large values of  $\sigma$  and for changing protection levels  $\omega$  holds asymptotically

$$E_D(\omega) \sim \begin{cases} \omega^Y e^{-\omega/\sigma} & \text{if } \xi = 0 \text{ (for large } \omega) \\ (X_{\max} - \omega)^{-1/\xi} & \text{if } \xi < 0 \text{ (for } \omega \text{ close to } x_{\max}) \\ \omega^{Y-1/\xi} & \text{if } \xi > 0 \text{ (for large } \omega) \end{cases} \quad (\text{B16})$$

$$SD_D(\omega) \sim \begin{cases} \omega^Y e^{-0.5\omega/\sigma} & \text{if } \xi = 0 \text{ (for large } \omega) \\ (X_{\max} - \omega)^{-0.5/\xi} & \text{if } \xi < 0 \text{ (for } \omega \text{ close to } x_{\max}) \\ \omega^{Y-0.5/\xi} & \text{if } \xi > 0 \text{ (for large } \omega) \end{cases} \quad (\text{B17})$$

5 All results from the previous sections are summarised in Table 1.

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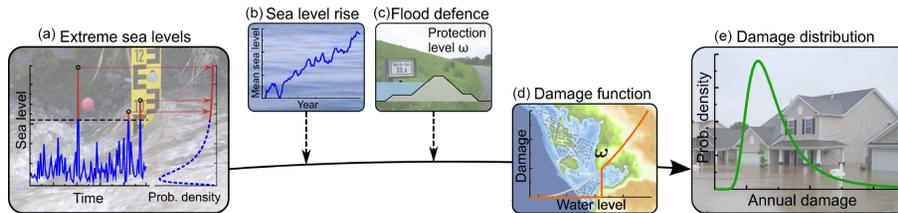
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**Figure 1.** From extreme sea levels to damage. **(a)** The analysis of extreme sea levels provides parameter estimations for the *generalised Pareto distribution*. **(b)** The distribution of sea levels is influenced by mean sea level rise. **(c)** Flood defence measures, such as dykes, set the threshold below which any damage is prevented. **(d)** The distribution of extreme sea levels is combined with the corresponding damage via a damage function, providing the total damage in the region under study at a certain maximum flood level. **(e)** From the resulting distribution of total annual damage the expected annual damage and its standard deviation can be derived. (Photographs: “Ilmepgel Ilmenau” by Michael Sander (2006), “Sea” by Dedda71 (2008), “Kilometermarkierung Deich” by Georg HH (2006), and “Nashville Flood” by Eric Hamiter (2010) from Wikimedia Commons – CC:BY-SA.)

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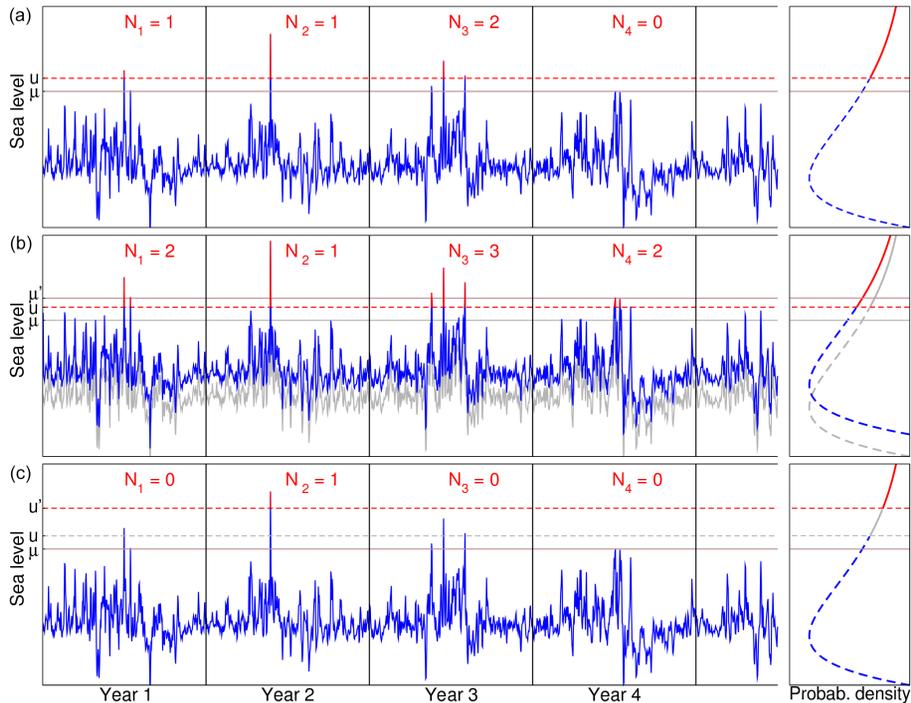
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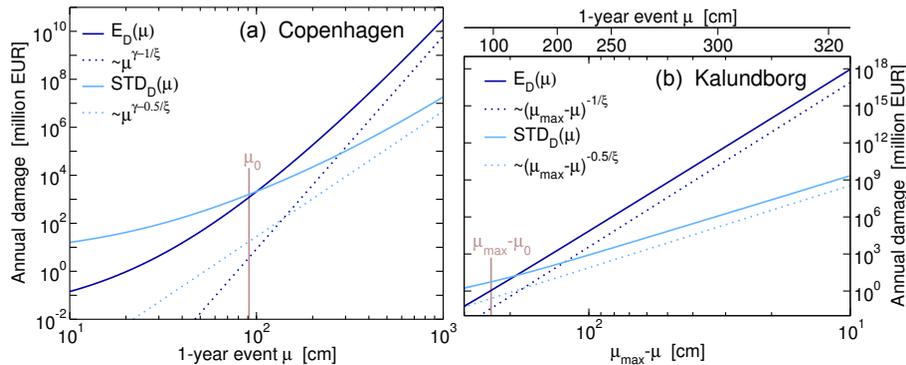


**Figure 2.** Illustrative time series of sea levels and their probability density function for several scenarios. **(a)** Current conditions with a threshold  $u$  and the 1-year event  $\mu$ . **(b)** Increased mean sea level with a corresponding shift of the time series and thus the 1-year event from  $\mu$  to  $\mu'$ . **(c)** Supposing a protection height of  $u'$  implying an adjustment of the threshold from  $u$  to  $u'$ . The values  $N_1, N_2, \dots$  represent in all cases the number of exceedances within the corresponding year. The average value of the  $N_i$  provides an estimator for occurrence rate  $\lambda$ .



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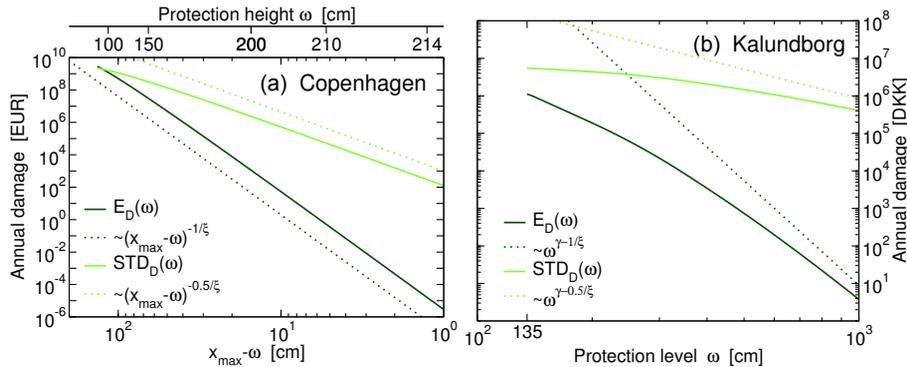


**Figure 4.** Expected annual damage (dark blue) and standard deviations (light blue) in **(a)** Copenhagen and **(b)** Kalundborg as a function of the mean sea level (parameterised by the 1-year flood  $\mu$ ). The dotted lines show the asymptotic relations Eqs. (12), (13) and (14) with  $\gamma = 1.6$  and  $\xi = -0.14$  (Copenhagen) and  $\gamma = 4.1$  and  $\xi = 0.08$  (Kalundborg). The values for the current 1-year floods  $\mu_0 = 91.21$  cm (Copenhagen) and  $\mu_0 = 95.35$  cm (Kalundborg) are indicated by brown vertical lines. The abscissa in the right panel is inverted and shows the difference between the 1-year flood  $\mu$  and  $\mu_{\max} = 332.04$  cm (at the top, the corresponding 1-year floods are displayed).



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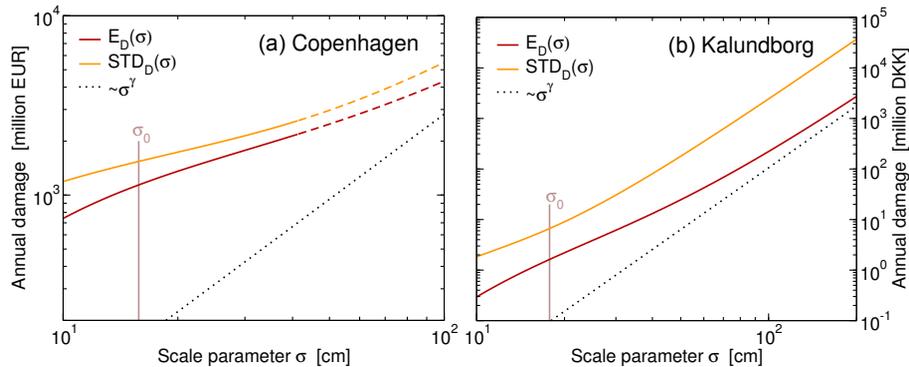


**Figure 6.** Expected annual damage (dark green) and standard deviations (light green) in **(a)** Copenhagen and **(b)** Kalundborg as a function of the protection level  $\omega$ . The abscissa in the left panel is inverted and shows the difference between the protection level  $\omega$  and the maximum possible water level  $x_{\max} = 215.28$  cm (at the top, the corresponding protection heights are displayed). Since no considerable damage occurs in Kalundborg for sea levels below 135 cm, only protection levels above  $\omega = 135$  cm are considered. The dotted lines follow the power laws from Eqs. (15) and (16) with the estimated damage function exponents  $\gamma = 1.6$  (Copenhagen) and  $\gamma = 4.1$  (Kalundborg).

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**Figure 7.** Expected annual damage (red) and standard deviations (orange) in **(a)** Copenhagen and **(b)** Kalundborg as a function of the scale parameter  $\sigma$ . The solid lines were numerically calculated with the available damage functions; the dashed continuations use an extrapolation of the damage function as a power law with exponent  $\gamma = 1.6$  (Copenhagen) and  $\gamma = 4.1$  (Kalundborg). The dotted line shows the asymptotic results from Eqs. (B14) and (B15) and the current values of the scale parameter  $\sigma_0 = 15.79$  cm (Copenhagen) and  $\sigma_0 = 17.78$  cm (Kalundborg) are displayed as brown vertical lines.

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