



**Reduction of
maximum tsunami
run-up**

N. Goseberg

Reduction of maximum tsunami run-up due to the interaction with beachfront development – application of single sinusoidal waves

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Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Abstract

Experiments are presented that focus on the interaction of single sinusoidal long waves with beachfront development on the shore. A pump-driven methodology is applied to generate the tested waves in the wave flume. The approaching waves firstly propagate over a horizontal bottom, then climbing up a 1 in 40 beach slope. The experiments reported here are confined to the surf similarity parameter of the waves ranging from $\xi = 7.69 - 10.49$. The maximum run-up of the tested waves under undisturbed conditions agrees well with analytical results of Madsen and Schäffer (2010). Beachfront development is modelled with cubic concrete blocks (macro-roughness (MR) elements). The obstruction ratio, the number of element rows parallel to the shoreline as well as the way of arranging the MR elements influences the overall reduction of maximum run-up compared to the undisturbed run-up conditions. Staggered and aligned as well as rotated and non-rotated arrangements are tested. As a result, nomograms are finally compiled to depict the maximum run-up reduction over the surf similarity parameter.

1 Introduction

Facing sea level rise the implications of coastal hazards such as cyclones, storm surges, floods or tsunami to endangered stretches of coastlines will increase in future. By combining geophysical and socio-economic data sets Nicholls and Small (2002) found that the population density of the near-shore coastal zone is three times higher than the average global population density. More interestingly, the highest number of inhabitants around 4×10^8 to 6×10^8 was found within the first 10 km inland of the shoreline based on census data of 1990 (Small et al., 2000; Nicholls and Small, 2002). Accelerated urbanisation and increasing construction activity is usually fostered in close proximity to the shoreline when population is substantially growing. Concurrently, many of these coastlines worldwide are situated in close proximity to subduction zones from where earthquake induced tsunami originate which is dramatically underlined by recent

NHESSD

1, 1119–1171, 2013

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



life-taking events such as the Indian ocean tsunami in 2004 or the 2011 Tohoku-Oki tsunami in Japan. In this regard the question arises how beachfront development in urbanized areas affect the run-up of tsunami. Some researchers recently made advancements in deducing water levels and flow velocities from eye-witness reports and video observations during the run-up and draw-down phase of tsunami indicating that flow velocities are in the range of 2 m s^{-1} to 11 m s^{-1} (Fritz et al., 2006, 2012; Shimozono et al., 2012). These water level and velocity measurements are most important for building design purposes since it is inevitable to validate analytical and numerical models which are capable of computing maximum fluid forces in the tsunami run-up zone (Yeh, 2006). Besides the purpose to establish and enhance current design guidelines aiming at constructing tsunami-proof infrastructure in a tsunami run-up zone, it is equally important to understand how maximum wave run-up is affected since evacuation and hazard mitigation planning often relies on the information about the spatio-temporal extent of the tsunami-induced inundation. In the tsunami run-up zone the flow is often observed to consist of an air, water and sediment phase which increases the difficulty to find analytical or simplified numerical solutions to the problem. Hence in literature a variety of laboratory studies documents the effort to understand how bores interact with surface-piercing vertical structures (Ramsden and Raichlen, 1990; Ramsden, 1996), the sediment transport initiated by solitary or cnoidal waves (Moronkeji, 2007), the forces on a single square column modelling free standing infrastructure located close to the shoreline (Nistor et al., 2009), the process of dam-break flow through an idealised city (Soares-Frazão and Zech, 2008) or the interaction between solitary-type wave (Cox et al., 2009) or error-function wave (Rueben et al., 2011) and the town Seaside, Oregon, US. Although some improvements were reported on the numerical modelling of wave-structure interaction (i.e. Tomita et al., 2007), it nowadays still remains cumbersome to cover the entire domain of a tsunami approaching the shoreline while interacting with beachfront development in three dimensions numerically. Hence in this study scale experiments were chosen methodically where the approaching tsunami was modelled by means of single sinusoidal waves in a water depth of

**Reduction of
maximum tsunami
run-up**N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



30 m (prototype condition). Although beachfront development and urban pattern are usually randomly distributed in prototype, it is decided to approximate those building pattern by regular structures, since from literature no hints on the parametrization of building density were determined. In total, four different pattern of MR elements modelling beachfront development were combined; (a) aligned, non-rotated, (b) aligned, rotated, (c) staggered, non-rotated and (d) staggered, rotated. The region of interest in this study is restricted to the shallow water region from where the wave starts climbing up a plane beach and the tsunami run-up zone where the long wave interacts with beachfront development. Even though it is unclear under which conditions building collapse is stimulated in prototype conditions, it is herein assumed that MR elements remain stable all the time. Currently, the flow processes incorporating non-stationary conditions during wave run-up with blocking MR elements on the shore have received the least attention from hydraulic engineering perspective and hence remain an active area of research. Hence, in the framework of the above model concept the following research questions should be addressed:

- How is the run-up of long waves caused by submarine earthquakes basically modified under the influence of beachfront development?
- How large is the reduction of maximum wave run-up compared to an undisturbed case and which processes or parameters dominate this reduction?

The problem of long waves which are climbing up a plane, sloping beach after propagating over a horizontal bottom is depicted in Fig. 1. Firstly, before the wave front reaches the shoreline, shoaling and refraction constitutes the main wave modification. At a specific water depth or for specific nonlinearity wave breaking may occur. Yet, Madsen and Fuhrman (2008) pointed out that wave breaking is mostly dedicated to shorter waves riding on top of lower frequent tsunami waves. Then, during their run-up and subsequent draw-down, the approaching tsunami may interact with beachfront development represented in this study by MR elements positioned at different distances to each other and in different arrangements on the beach above the still water line. In

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



$F(\Pi_1, \Pi_2, \dots, \Pi_{n-r}) = 0$ by recasting into dimensionless Π terms. The dimensionless variables Π are composed of power products of the former dimensional variables x . Undisturbed wave run-up was investigated by various authors (i.e. Hunt, 1959). Dimensional analysis for this problem was addressed in detail by Müller (1995) who derived dimensional analysis for the problem of impulsive wave propagation and run-up in alpine lakes induced by rock fall. Limiting the problem to two-dimensionality, neglecting model, scale as well as frictional effects, relative run-up might be related to the following dimensional variables:

$$f \left(\underbrace{g, \rho, \nu, \zeta}_{\text{universal properties}}, \underbrace{a, T, \beta, c}_{\text{plain wave run-up}}, \underbrace{b_{\text{mr}}, b_{\text{st}}, n, \phi, \Gamma}_{\text{MR related}} \right) = 0 \quad (1)$$

where g is gravitational acceleration, ρ is density of the fluid, ν is kinematic viscosity, ζ is surface tension. $H/2 = a$ describes wave amplitude which for sinusoidal waves equals height $H/2$ of a wave with period T and with propagation speed c running up on a beach with slope $\tan \beta$. A third group of variables accounts for wave structure interaction due to MR elements on the shore. These are b_{mr} describing the edge length of an MR element, the distance between the MR elements b_{st} , the number of MR element rows n , ϕ defining the angle of rotation of MR elements with respect to the incident wave direction and Γ indicates various sets of MR configurations. By means of the theorem of Buckingham it is hence possible to recast Eq. (1) in the following form containing

nondimensional variables, which will be defined and discussed subsequent to Eq. (2):

$$F \left(\underbrace{\frac{c}{\sqrt{gh_0}}, \frac{ch_0}{\nu}, \frac{c^2 h_0 \rho}{\zeta}}_{\text{universal properties}}, \underbrace{\frac{2a_0}{h_0}, T\sqrt{g/h_0}, \beta, \frac{L_0}{h_0}}_{\text{plain wave run-up}}, \underbrace{\psi_{mr}, \psi_{cs}, \psi_{ls}, \Gamma}_{\text{MR related}} \right) = 0$$

where h_0 indicates water depth over the horizontal bottom being utilized to define various nondimensional variables. First, Froude, Reynolds and Weber number can be written. In hydraulic and coastal engineering, free surface problems are traditionally modelled under the assumption of Froude similitude. Such choice is admissible in case that viscous effects and surface tension are in a negligible order of magnitude. Based on linear wave theory involving the effects of kinematic viscosity, Schüttrumpf (2001) defines a critical Reynolds number to be $Re_{crit} = 10^4$ below which the Froude number of waves propagating over horizontal bottom is no longer independent while values of $Re_{crit} = 10^3$ are recommended during the run-up of waves with decreasing water depth and during overtopping. Besides wave propagation and run-up, the local MR Reynolds number should be above a threshold upon which viscous effects during wave-structure interaction is excluded. Those Reynolds numbers during propagation Re_{pro} and during interaction Re_{int} are defined as follows:

$$Re_{pro} = \frac{ch_0}{\nu} \quad (2)$$

$$Re_{int} = \frac{\nu b_{mr}}{\nu} \quad (3)$$

What is more, in regard to minimizing viscous effects in scaled modelling of water waves Le Méhauté (1976) stipulates minimal flow depth not smaller than 0.02 m and wave periods not shorter than 0.35 s. Due to limited water depth at the wave front it is

NHESSD

1, 1119–1171, 2013

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



anticipated that viscosity adds to overall scale effects in this spatially varying region of the flow.

The effect of surface tension to the propagation and run-up of long waves in scale models was expressed in literature with the help of the extended dispersion relation (Le Méhauté, 1976), which analytically describes the dependency between Weber and Froude number. Dingemans (1997) expressed the effectiveness of the surface tension in an increase of gravitational acceleration but outlined that the influence of surface tension to wave action is mostly dominant when capillary waves are investigated. Due to the nature of surface tension, its effect on fluid motion amplifies where the fluid surface exhibits small radii of the liquid surface or when the distance to solid boundaries is rather small, i.e. for drops or bubbles (Peakall and Warburton, 1996). At most times during wave propagation and run-up, this requirement is not meet, so that scale effects due to surface tension can be excluded.

After discussing the choice of similitude and its implications for scale effects, a second group of dimensionless variables indicates parameters whose variation was often used to study wave run-up on a plain beach. It mainly depends on the relative wave height $2a/h_0$ and the relative wavelength L_0/h_0 . In addition, relative wave height and the slope angle β listed in Eq. (2) as two of the four dimensionless parameters determining the plain undisturbed wave run-up allow for the definition of the often used surf similarity parameter that is further discussed in Sect. 2.2 in connection with analytical run-up laws. It is defined as

$$\xi = \frac{\tan \beta}{\sqrt{\frac{2a_0}{L_0}}} \quad (4)$$

where index 0 refers to deep water wave parameters. In this research focussing on the influence of beachfront development on the maximum wave run-up, only the surf similarity parameter is varied through variation of the wave amplitude.

Finally, a third group of nondimensional variables (Eq. 2) is introduced in order to express the influence of MR elements on the wave run-up. Geometrical definitions are

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



given in Fig. 1. Most importantly, one has to define cross- ψ_{cs} and long-shore rates of obstruction ψ_{ls} . Besides these horizontal obstruction rates, a component expressing the vertical obstruction of the flow cross section ψ_{mr} could be named similarly. Yet, vertical obstruction is only relevant when overtopping of MR elements is concerned which shall be excluded for the moment. The cross-shore obstruction ratio can be defined as

$$\psi_{cs} = \frac{L_{mr}}{n\hat{b}_{mr}} = \frac{n(\hat{b}_{mr} + b_{st}) - b_{st}}{n\hat{b}_{mr}} = \frac{n(b_{mr}/\cos\varphi + b_{st}) - b_{st}}{nb_{mr}/\cos\varphi}$$

where $\hat{b}_{mr} = b_{mr}/\cos\varphi$ denotes the width of the MR elements according to their angle of rotation to the incident wave direction (cp. also Fig. 4). The length of the obstructed beach in the direction of wave run-up is named L_{mr} , which can be defined as follows

$$L_{mr} = n(b_{mr} + b_{st}) - b_{st} \quad (5)$$

with: n denoting the number of MR element rows parallel to the still water line, b_{mr} defines the edge length of a single MR element and b_{st} defines the distance from one MR element to the next one, i.e. street width. The cross-shore rate of obstruction depends on the ratio of the number of element rows times the width to the length of the macro-roughness area in on-shore direction. This geometrical ratio accounts for the obstruction exposed to the in-land propagating flow. The ratio ψ_{cs} decreases if the street width tends to zero, $b_{st} \rightarrow 0$. The ratio grows if the number of MR rows is increased. The long-shore obstruction ratio is given by

$$\psi_{ls} = \frac{\hat{b}_{mr}}{\hat{b}_{mr} + b_{st}} = \frac{b_{mr}/\cos\varphi}{b_{mr}/\cos\varphi + b_{st}} \quad (6)$$

where \hat{b}_{mr} equally expresses the effective width of the MR elements with respect to the approaching wave front. The long-shore rate of obstruction increases according to Eq. (6) when the ratio of the macro-roughness width to the sum of the street width plus

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



the macro-roughness element width tends to unity. A long-shore rate of obstruction of unity is equivalent to a complete obstruction of the beach and results in full reflection of the wave and the run-up tongue.

2.2 Analytical wave run-up

Wave run-up was extensively studied in literature. Most commonly, the work of Hunt (1959) is referenced who reprocessed former published experimental studies, conducted additional experiments and analysed these aiming at design of seawalls and breakwaters. For breaking waves, the non-dimensional run-up becomes a function of beach slope, incident wave height and wave steepness which reads

$$\frac{R}{H_0} = \xi_0 = \tan \beta \left(\frac{H_0}{L_0} \right)^{(-1/2)} \quad \text{for } 0.1 < \xi_0 < 2.3 \quad (7)$$

where R/H_0 denotes the relative run-up normalized by the deep water wave height H , ξ_0 is the surf similarity parameter. Equation (7) was derived on the basis of experimental research and thus the validity is limited to steep beach slopes between 1/10 to vertical walls. In the sequel, also solitary waves (Synolakis, 1987) or N-waves (Tadepalli and Synolakis, 1994) were addressed in connection with long wave run-up. Often, the run-up formulas reported in literature were derived empirically and for the special conditions used in this research with long wave periods and very mild beach slopes it is questionable if these approaches are valid to compare our experimental results with. However, recently Madsen and Schäffer (2010) proposed analytical solutions for the run-up and run-down based on the non-linear shallow water wave theory for various wave forms, where the duration and the wave height could be specified separately in order to overcome the length scale deficiency of available long wave models. This analytical approach resulted in establishing run-up/run-down and velocity formula for various waveforms such as isocetes N-waves, single waves and sinusoidal waves but with the important exception that these waveforms did no longer exhibit a tie to wave

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



height-to-depth ratio. The normalized run-up is summarized as a function of the surf similarity parameter derived, which reads (Madsen and Schäffer, 2010)

$$\frac{R}{a_0} = \chi_{\text{elev}} \pi^{1/4} \left(\frac{a_0}{h_0} \right)^{-1/4} \xi_0^{-1/2} \quad (8)$$

where R denotes run-up, a_0 is offshore amplitude, χ_{elev} and χ_{velo} is shoreline elevation and velocity, $\Omega = 2\pi/T$ is a duration in analogy to the common wave period and h_0 is offshore water depth. Surf similarity ξ is given by Eq. (9), which bases on the previously defined surf similarity Eq. ((4)) but is recast by Madsen and Schäffer (2010) in terms of the effective duration Ω of the wave. It reads

$$\xi = \sqrt{\pi} \left(\frac{a_0}{h_0} \right)^{-1/2} \left(\frac{\Omega^2 h_0}{g \tan \beta^2} \right)^{-1/2} \quad (9)$$

where additionally $\tan \beta$ describes the beach slope. A theoretical breaking criterion is equally given by Eq. (10) (Madsen and Schäffer, 2010)

$$\frac{R_{\text{break}}}{a_0} = \frac{1}{\pi} \xi^2 \quad (10)$$

3 Experimental set-up

Building on the theoretical framework described above, the experimental set-up and the related experimental programme is outlined in the following. Results of dimensional analysis are utilized to decide which variables have to be varied. Details of the varied variables are reported in Sect. 3.2. Froude similitude is chosen to downscale prototype conditions since from previous Sect. 2.1 it can be ruled out that viscosity and surface tension may – under the constraint of to tine length or time scales – falsify run-up measurements significantly. The wave flume, its components to generate long

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



waves as they are possibly generated by tsunamigenic sub-sea earthquakes as well as wave generation methodology have recently been described by Goseberg (2011, 2012); readers are referred to those publications for detailed information whereas a brief outline is presented in Sect. 3.1.

3.1 Wave flume

The experiments were carried out in the closed-circuit wave flume (race-track type) at the Franzius-Institute for Hydraulic, Waterways and Coastal Engineering. A schematic drawing of the 56.94 m long flume illustrates the principal design of the facility (Fig. 2). The clear width of the flume amounts to 1.00 m. The flume is made of two parallel and 19.00 m long straights whose ends are connected by two semicircles of 6.87 m diameter. While one of the straights accommodates the run-up region, the other integrates a pump station that allows for the acceleration and deceleration of the water body. The water depth is kept constant at 0.31 m during all the reported experiments.

For the purpose of generating long waves in the flume, it was separated into a reservoir and a propagation section by a 1 in 40 sloping beach wedge made of wood, sand, plaster and PVC board. Figure 2 also gives the location of the sloping beach section which is located opposite to the pump station. A solid impermeable wall completes the beach wedge at its upper limit to ensure that no water flows from the propagation to the reservoir section or vice versa. The beach is fully covered with poly-vinyl-chloride board and all remaining gaps are sealed with silicone to build a smooth surface for the run-up and drawdown process.

3.2 Test programme

In this study single sinusoidal waves with a fixed period $T = 60$ s were used to study the onshore interaction of long waves and beachfront development. In literature, many studies exist that apply solitary waves when tsunami evolution in the vicinity of the shoreline is investigated (e.g. Synolakis, 1987). Few studies yet base on cnoidal waves

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



(Moronkeji, 2007) while a variety of publications apply dam-break induced bores to study forces on free-standing structures on the beach (Ramsden, 1996). Following the argumentation of Madsen et al. (2008) it recently became questionable if solitons are well suitable to serve as a model for near-shore tsunami. A main argument from the authors was that solitons lack in terms of geophysical time- and length-scales that are typical of tsunami. This argument was recently underlined by offshore GPS and bottom pressure measurements at the Japanese coast originating from the 2010 Chilean earthquake detecting wave periods much longer than those replicable from the solitary wave assumption (Kawai et al., 2012). Hence it was deliberately chosen to step back to the most simple wave form possible in laboratory. With respect to the length and time scales chosen, a period of $T = 60$ s was applied in this study leading to prototype condition periods of $T = 600$ s which is well in the range of measured realistic tsunami periods. Furthermore due to time constraints only leading depression waves were tested so far. Figure 3 presents the investigated waves with amplitude variations. Those applied single period sinusoidal waves consist of a leading wave through followed by a wave crest. Prior and after these single sinusoidal wave periods surface elevation is kept constantly zero. Standard deviation is additionally given in order to show quality of the generated waves and their repeatability.

Table 1 lists the wave properties measured at the wavemaker and the repetitions. Some experiments were repeated quite often in order to investigate the repeatability of the wave generation method; results are reported in Goseberg (2012) who generally found satisfactorily small standard deviation for repeated experiments.

The MR elements were manufactured as solid, cubic blocks with an edge length of 0.10 m manufactured of standard concrete mixture. As form work a standard cube mold was utilized to pour in the concrete. A total number of 100 complete cubes and 20 half cubes were manufactured. Figure 5 depicts an example configuration of the MR elements in the flume. One notices some minor imperfections stemming from the production process of the MR elements (blowholes, dents) but in summary the quality of the concrete cubes and their surfaces is regarded to be satisfactory for the experiments.

Table 2 further outlines the number of experiments and the applied wave conditions during the experiments. With respect to the MR configurations different total numbers of experimental runs were possible.

The varied variables leading to the total number of experimental runs are the distance between the elements b_{st} , the number of MR element rows n , the angle of the elements with respect to the incident wave direction φ and the choice of an offset between the element rows. In case of the aligned configuration (“al” in Table 2), four street widths were realized in the flume which are 2.5 cm, 4.2 cm, 6.6 cm and 10.0 cm whereas for the staggered configuration (“st” in Table 2) only three street widths were used (2.5 cm, 5.9 cm, 10.9 cm). This is because the flume width is limited to 1.0 m. Two element angles $\varphi \in (0^\circ, 45^\circ)$ are further conceivable. Additionally three different numbers of macro-roughness element rows were tested in order to investigate the influence of the obstructed long-shore length, $n \in (1, 5, 10)$. A prerequisite for the experiments was the strict compliance of two-dimensionality and thus symmetry. In preliminary experiments it was found that the application of whole-numbered distances between the MR elements led to large deviations of the maximum measured run-up compared to cases where symmetry conditions directly at the flume walls were met. Figure 4 depicts two examples of MR configuration in the flume where correct application of symmetry was given.

In both cases it was required to correctly close any distances left between the half MR elements and the flume walls in order to circumvent any flow through these gaps. Additional photographs of these wall and symmetry treatment is shown in Fig. 5.

3.3 Measurements and instrumentation

In order to evaluate run-up reduction due to the interaction of the incoming waves with the beachfront development a large set of parameters was recorded. Besides the measurement of surface elevations along the horizontal bottom, water depth were collected on the beach slope. The maximum run-up was captured with a driftline approach.

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Additionally, flow velocities were recoded at selected positions along the propagation path. A comprehensive sketch of positions of sensors is given in Fig. 6.

According to the local water depth wave heights were measured by resistance wave gauges (RWG). Those positions in the vicinity of or on the sloping beach are indicated in Fig. 6. Besides, a few times selected experiments were conducted where the positions of the RWG's were shifted to the bended section of the wave flume in order to investigate how the bend might affect the wave propagation. Details are given in Goseberg (2013). The non-linearity of the RWG (Delft Hydraulics Laboratory, range 0.5 m) is specified 0.5 % of the measurement range. Where still water depth did not exceed 0.2 m, upward facing pressure sensors (PRS) were mounted inside the beach slope (orthogonal to the surface) in order to capture water level variations during the experiments. Two different sorts of sensors were in operation – TM/N sensors by STS AG (range 0.0 to 1.5 m, non-linearity ± 0.25 % FS) and PDCR 1830 by GE Druck (range 0.0 to 0.7 m, non-linearity ± 0.1 % FS). All sensors were calibrated against a calibration unit before mounting. Similarly, exact positions of the PRS are indicated in Fig. 6. Ultrasonic distance sensors (UDS) were utilized to detect water level variations in the landward facing beach area behind the MR elements (Wehrhahn manufacturer, range 0.07 to 2.00 m). Sound transducers of the UDS are mounted 0.40 m above the PVC board surface rectangularly pointing to the beach inclination.

The detection of the maximum wave run-up is of key importance in this study. Besides visual detection, beach parallel capacitive type wave gauges (Oumeraci et al., 2001) or camera-based techniques are feasible. Alternatively to these methods, in this study a drift line based approach is chosen to detect the maximum run-up of the waves. The basic idea resembles drift lines along natural beaches where deposited debris is one potential demarcation of the highest run-up of the incoming sea condition. Analogously, in this experiments run-up values are determined with measuring tape manually after each experiment. As tracer material a special artificial granulate was applied. It is a low-density polyethylene granulate of 0.91 to 0.94 g cm⁻³ density with an average diameter of 2 to 3 mm. Even though the experiments were designed fully

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures



Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



two-dimensional by observing the symmetry in long-shore direction, measured run-up results vary along the flume width. The maximum run-up indicated by the applied tracers is not distributed along a straight line but parabolic with the apex in the flume axis. This phenomena arises from the fact that wall roughness of the plastering contributes to the energy dissipation in the boundary layer of the flume walls. The reported wave run-up, that is referred to in the following, is therefore defined as the mean of the lowest and the highest maximum run-up measured in each experimental run. Figure 7 exemplifies how the lowest R_{\min} and the highest maximum run-up R_{\max} was defined. In order to circumvent influences from wall effects, areas as close as 0.05 m to the flume walls were excluded from the run-up determination. After determination of minimal and maximal run-up, average values were recorded.

3.4 Experimental procedure and data processing

A typical experimental procedure encompasses the wave generation described in Sect. 3.1 along with the run-up measurements, the wave height and water depth measurements at various positions as well as the velocity measurements described in Sect. 3.3. After a long wave was generated it propagated along the flume and climbed up the sloping beach until its maximum run-up elevation was reached. Thereafter the particle velocity turned negative and the run-down phase followed.

Data processing aimed at converting, filtering, smoothing and decimating recorded signals. For some reasons an internal sampling rate of 1.0 kHz was applied; thus decimation was accomplished by extracting every fiftieth data point from the signal subsequently after filtering and smoothing took place. Data were stored in a digital experimental notebook where in addition to the measured data, processed variables such as the configuration tag Γ , the vertical run-up $R_{v,\min}$, $R_{v,\max}$, the long- and cross-shore obstruction ratio ψ_{ls} , ψ_{cs} , the wave period T , crest a_{cr} and trough a_{tr} amplitudes and wavelength L_0 at the wavemaker, the wave number k , the frequency Ω and the surf similarity parameter ξ_0 . Filtering with good results was obtained on the basis of the empirical mode decomposition (EMD) technique (Huang et al., 1998).

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



4 Results

The wave run-up of single, sinusoidal waves was investigated by means of scale experiments. Basically wave height was varied and surf similarities between $\xi = 7.69 - 10.49$ were achieved. Firstly 78 experiments were conducted to analyse the run-up on a plain 1 in 40 sloping beach without any surface-piercing MR objects. These results are addressed in Sect. 4.1. Subsequently, experiments were conducted that focus on the long wave run-up reduction due to MR elements which were arranged in four different patterns on the shore, details are listed in Table 2. Resulting reduced run-up is described in Sect. 4.2. Additionally, in Sect. 4.3 it is qualitatively presented how the surface elevation evolves with and without MR elements on the shore in longitudinal direction.

4.1 Undisturbed run-up

As outlined in Sect. 2.1, the wave run-up directly relates to a number of dimensionless variables as presented in Eq. (2). In this research basically the wave nonlinearity $2a_0/h_0$ was varied leading to the variation of surf similarity as mentioned before. Figure 8 depicts the maximum relative run-up $R_{\text{mean}}/a_{\text{cr}}$ as a function of the nonlinearity $\epsilon = a_{\text{cr}}/h_0$ of the investigated waves. Here, the maximum relative run-up is defined as the average of the measured run-up outlined in Sect. 3.3. The general trend of the measured data set is expressed by the polynomial fit of the degree $n = 1$ which basically reduces to a linear function.

Since the overall aim was to investigate wave run-up reduction due to MR elements onshore, no further variations of the wave period T and the beach slope $\tan(\beta)$ were accomplished in this research. In tendency, it becomes clear that an increasing nonlinearity results in increasing relative run-up. Yet, only nonbreaking conditions were tested. It is also obvious that the relative run-up related to the positive crest amplitude is remarkably high. It ranges between $R_{\text{mean}}/a_{\text{cr}} = 3.1 - 3.4$. In order to check whether our data sets are consistent with existing research, Fig. 9 presents the maximum relative

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



5 wave run-up in the framework of the surf similarity parameter (cp. Eq. 8 and (Madsen and Schäffer, 2010)). The maximum relative run-up, which is normalized by the positive amplitude at the wavemaker, decreases with an increasing surf similarity parameter ξ_0 . The increase in non-linearity, which can be expressed as the ratio of the applied wave amplitude of the crest a_{cr}/h_0 , reduces the surf similarity of the applied wave condition at the beach under investigation. Higher non-linearity of the wave also indicates higher energy content. This finding underlines the statement of Madsen and Fuhrman (2008) after which critical flow depth and flow velocities occur for surf similarity in the order of $\xi_0 = 3 - 6$ and in turn require relatively mild beach slopes.

10 The resulting measurements depicted by the point clouds are in good agreement with the analytical framework. Yet, it is apparent that the results do not perfectly fit the expected analytical behaviour even though the results already indicate that the theory and the functional relation of the nondimensional terms is well established. The analytical approach fully neglects the influence of roughness and even though the experiments aim at eliminating roughness influences to the greatest extent, energy losses due to roughness influence are still existent. From the degree of agreement between the experimental data and the analytical approach it can thus be reasoned that the experimental set-up is in principal suitable to reproduce long wave motion in scaled models and in the following it shall be demonstrated how MR elements contribute to reduced long wave run-up quantitatively.

4.2 Run-up reduction

25 Cross- ψ_{cs} and long shore obstruction ratio ψ_{ls} as well as the MR configuration Γ was found to influence run-up reduction of long waves climbing up a plain, sloping beach as described in Sect. 2. According to Table 2 four MR configurations were tested and results are presented in this section in parts. Additional plots are presented in the annex. Firstly, the relative run-up of long waves normalized by the positive wave amplitude at the wavemaker is shown in Fig. 10 as a function of the long-shore obstruction ratio

(Eq. 6) where an long-shore obstruction ratio of unity denotes full obstruction of the beach section.

The MR configuration $\Gamma = 1$ comprises MR elements in aligned arrangement with non-rotated elements. The resulting run-up reduction can generally be described by a downward opening parabolic function. Note that two of the experiments were excluded from the analysis because measured run-up appeared to be non-physical. Four different long-shore obstruction ratios were tested. Long wave run-up diminishes with increasing ψ_{JS} . It appears that the influence of the incident wave amplitude indicated by the marker symbols tends to increase with increasing long-shore obstruction ratio. Higher incident wave amplitudes result in higher water levels in the vicinity of the macro-roughness area. A reason for this finding could be that higher velocities in between the MR openings result in higher energy exchange between the predominant inland-directed flow and the calmed areas in the back of MR elements; thus leading to increasing energy dissipation. This observation agrees with the findings of Goto and Shuto (1983) who investigated the effects of wooden pillars exposed to a steady flow and similarly reported momentum exchange between the main flow and eddies in the back of their obstacles. Figure 10 also indicates that the number of MR element rows n has a much smaller influence on the maximum run-up reductions. For values of $\psi_{JS} = 0.8$ only marginal differences exist between the relative run-up reduction $R_{\text{mean}}/a_{\text{cr}}$. Yet, it appears that the scatter of the measured run-up slightly increases to some extent, which could be attributed to an increase in turbulence when the flow interacts with a higher number of MR element rows.

In contrast to MR configuration $\Gamma = 1$, Fig. 11 depicts MR configuration $\Gamma = 4$ where MR elements are arranged in a staggered manner and rotated by the angle $\varphi = 45^\circ$. The functional dependency between normalized run-up and long-shore obstruction ratio is similar to the aligned, non-rotated configuration. A parabolic polynomial is also suitable to fit the measured data. Yet, the run-up reduction is higher than in the previous case.

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



In comparison with MR configuration $\Gamma = 1$ a substantial additional run-up reduction could be found for configuration $\Gamma = 4$ (cp. Fig. 4, right panel). It appears that for cases with $n = 5$ MR element rows in staggered, rotated configuration the normalized run-up decreases by 4.2% for $\psi_{1s} = 0.5$ and by 21.4% for $\psi_{1s} = 0.8$ while in case of $n = 10$ MR element rows those reductions are 13.1% for $\psi_{1s} = 0.5$ and 19.6% for $\psi_{1s} = 0.8$ respectively. Additionally, some scatter is introduced into the data in the case of $\Gamma = 4$. This fact underlines the hypotheses that the run-up flow experiences higher turbulence production in this case. Additionally, the approaching wave front is obliquely reflected at every even MR element row directing the wave front to the back side of the uneven MR element rows and thus leading to additional energy dissipation through vortex formation. A summary of the results and the comparison of investigated MR configurations is presented in Table 3.

According to the tabulated results presented in Table 3 the highest run-up reduction is achieved by MR configuration $\Gamma = 2$ which depicts the aligned, rotated case (cp. Table 2). This configuration leads to additional run-up reduction of 21.3% for $\psi_{1s} = 0.8$ when $n = 5$ MR element rows are tested. This reduction slightly decreases for the case with $n = 10$ MR element rows which at first contradicts the statement that energy reduction increases with increasing number of MR element rows. Although here, this opposed trend is rather related to the scatter of data that leads to variations in the polynomial fitting curves applied for the summary table.

In order to incorporate the non-dimensional variable cross-shore obstruction ratio ψ_{cs} it is convenient to define an additional non-dimensional variable that accounts for the entire horizontal obstruction ratio. It may read

$$\psi_{2d} = \psi_{cs} \psi_{1s} \quad (11)$$

Equation (11) relates the long- and the cross-shore obstruction ratio. While the long-shore obstruction ratio resides between zero and unity, the cross-shore obstruction ratio, by definition, reaches unity if only one macro-roughness row is considered and decreases with a growing number of macro-roughness element rows. For one MR

element row the product of ψ_{cs} and ψ_{ls} reduces to the long-shore obstruction ratio ψ_{ls} whereas it grows with an increasing number of element rows. It has to be noted that the 2-D obstruction ratio bases on the assumption that the street widths in the cross- and long-shore direction are equal. As an example, Fig. 12 presents the normalized run-up as a function of the two-dimensional obstruction ratio ψ_{2d} according to Eq. (11) for the case of $\Gamma = 2$ with rotated elements in an aligned arrangement.

The upper panel of Fig. 12 is identical to those results where $\Gamma = 2$ and $n = 1$ plotted only as a function of ψ_{ls} while the lower two panels reveal the functional dependency on the two-dimensional obstruction ratio as defined by Eq. (11). It is obvious that with increasing MR element rows n the values of 2-D obstruction ratio ψ_{2D} approach unity. As clearly seen, 2-D obstruction ratio is a non-linear variable and here it is defined so as to acknowledge that the run-up reduction effect is most pronounced for few MR element rows whereas this effect diminishes with increasing number of MR element rows.

From an engineering perspective it is preferable to either yield empirical equations or derive nomograms which are suitable to predict the reduced run-up due to a given MR configuration in comparison with the undisturbed case. As outlined above, besides the specific MR element configuration it is the long-shore obstruction ratio that basically dominates the run-up reducing effect of the investigated MR elements while cross-shore obstruction ratio. On that basis it is convenient to focus on the dominating long-shore obstruction ratio ψ_{ls} to deduce conclusions from the experimental findings. Figures 13 and 14 illustrate the relation of the relative run-up R_{mean}/a_{cr} to the surf similarity parameter ξ_0 within the analytical framework of Madsen and Schäfer (2010) for the MR configurations $\Gamma = 2, 4$. Nomograms referring to the configurations $\Gamma = 1, 3$ are presented in the annex for completeness. The reduced relative run-up according to the MR elements on the shore is expressed by means of the long-shore obstruction ratio and is plotted as a error bar line where the line denotes the averaged run-up of the cases $n = 1, n = 5$ and $n = 10$. The error bars mark the standard deviation of the measured relative run-up. Generally, a smaller relative run-up is associated with a higher

NHESSD

1, 1119–1171, 2013

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



number of macro-roughness element rows. The relative run-up associated with the long-shore obstruction ratio slightly decreases with an increasing non-linearity of the applied waves, which is most likely contributed to greater momentum exchange within the MR element area. Six error bars for each long-shore obstruction ratio indicate those wave conditions tabulated in Table 1.

For both nomograms presented in Figs. 13 and 14, the reduction of the relative run-up decreases slightly with decreasing surf similarity ξ . The reduction of the relative run-up compared to undisturbed conditions is less significant for a long-shore obstruction ratio of $\psi_{1s} = 0.393$; yet, it appears that the maximum run-up tends to zero when the long-shore obstruction ratio is further increased. Standard deviation of measured relative run-up is rather small indicating that most of the experimental results are fairly repeatable. Secondly, it can be observed that the relative run-up again decreases with an increase in non-linearity, which is probably contributed to the growth in momentum exchange when waves of higher non-linearity interact with the MR elements. Higher turbulence production is assumed to affect the increase in run-up reduction.

4.3 Longitudinal evolution of surface elevation

Finally, the longitudinal evolution of the free surface elevation of the undisturbed long wave run-up is compared to the MR configuration $\Gamma = 4$, a staggered and rotated configuration, with $\psi = 1s = 0.739$ and $n = 5$ MR element rows. Here, the qualitative influence of MR elements configuration to the spatial distribution of the water level at distinct time steps should be assessed. Out of the sequence of available time steps of the wave evolution along the beach wedge Fig. 15 depicts the situation during wave run-up while Fig. 16 focuses on the draw-down phase. Time is given non-dimensional with reference to the beginning of the wave generation at the wavemaker; thus, non-dimensional time $\hat{t} = t\sqrt{g/h_0}$ relates to the total time for the wave needed to propagate along the horizontal bottom section and climbing up the 1 in 40 sloping beach. Both figures illustrate the development of the surface elevation profile along the beach wedge. The

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



normalized surface elevation $\hat{\eta} = \eta/h_0$ is given as a function of the non-dimensional horizontal coordinate $\hat{x} = x/h_0$. The position and the height of the obstacle area is given by the dash-dotted gray lines while the beach surface is presented by the full gray line. A wave with a non-linearity of $\epsilon = 0.093$ is used. Generally, three qualitative phases can be determined during the evolution of the wave. First of all the initial draw-back of the water, which is due to the leading depression wave character, characterizes the wave evolution. In this phase no significant differences can be found for the cases with and without MR elements. In a second phase when the water surface elevation is rising on the beach, in tendency higher water levels are determined for the obstructed case. Thirdly, during the draw-down of the long wave the situation inverts and the surface profile is lower for the case with MR elements present. The flow of water which is in-between and behind the MR area is retarded significantly (Fig. 16).

An interesting feature is demonstrated in Fig. 15. The approaching wave tongue collides with the first MR element row and a shock wave is formed which propagates seawards. This flow feature is observed for all investigated long-shore obstruction ratios and for all MR configurations. The height of the shock wave qualitatively depends on the long-shore obstruction ratio, though this dependency has not been investigated in detail in this research. The front of the initiated reflection wave is comparable with a monoclinal rising wave which though rapidly varies during its seaward propagation. In the presented case the surface gradient of the positive surge, whose front is initially steep, decreases rapidly and it flattens out noticeably (lower panel of Fig. 15). The situation inverts as soon as the long wave starts the draw-down process (Fig. 16). In the vicinity of the MR elements the undisturbed case faces higher water levels than in the case with MR elements. This is a direct result of the effect of the MR elements causing decreasing maximum wave run-up, since obstacles are causing partially reflections of the shoreward transported wave energy. Finally, an additional feature appears within the MR element area. There, the water levels are higher with MR elements compared to the undisturbed case. This is due to the fact that the volume of water which at this time is situated in-between and behind the obstacles is retarded significantly and is

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



only displaced during the draw-down phase with a notable time lag. In the case of MR it is apparent that the negative, seaward directed surface gradient in the vicinity of the obstacles is considerably higher than that without obstacles. In consequence, the flow velocities are qualitatively increased compared to the undisturbed run-up.

5 Discussion

Section 4 presents experimental results which were obtained to address the question of how MR elements on a shore influence the run-up of a long wave approaching a plain, sloping beach. Based on the parameters long- and cross-shore obstruction ratio as well as on four different MR element configurations run-up reductions were determined compared with an analytical approach after Madsen and Schäffer (2010). Prototype conditions were approximated by means of a plain, sloping beach and MR elements consisting of concrete cubic blocks representing beachfront development with close proximity to the shoreline. Instead of an unstructured urban building pattern, regular MR element patterns were chosen. Additionally, the three-dimensional prototype character was substituted by a 2-D scale model, since laboratory space was limited. Yet, this 2-D approach provides reasonable results for prototype conditions in case that homogeneity in long-shore direction can be assumed. Larger deviations are likely when up-scaled results are going to be utilized to predict run-up extents of tsunami in coastal regions where strongly inhomogeneous urban patterns exist. However, the presented results are considered suitable for the calibration of analytical or numerical models that also aim at predicting the energy dissipation and run-up reduction due to beachfront development. Long waves were generated by means of a novel long wave generator which utilized high-capacity pumps under the control of an industrial controller (Goseberg, 2012). Due to limited laboratory space and some related restrictions, wave propagation was through a 180° bend. Therefore, the question arises whether this laboratory setting negatively influences the quality and shape of the applied wave conditions. In this regard, Goseberg (2011) researched how transverse and

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



propagation-wise velocity cross profile evolved at times of maximum positive or negative flow. It was found that rather uniform vertical velocity profiles developed which is in accordance with the long wave theory and that, more importantly, transverse velocity components inside the flume bend were marginal, so that no direct influence on the wave propagation is expected.

The basic flow features which were observed in the course of the run-up and draw-down phase of the wave climbing up the beach and interacting with the MR elements consist of acceleration and deceleration due to sudden narrowing and widening in between the MR element area. This process increases with reduction of long-shore obstruction ratio and with increasing non-linearity of the incident waves. These features directly influence the maximum run-up that was under investigation. These findings are in close agreement with previous findings for steady-state flows. For example, Soares-Frazão et al. (2008) qualitatively found the same flow pattern (sudden narrowing, widening and deceleration) in the case of a flood wave impinging an urban settlement of houses which are arranged in an aligned and staggered configuration. Similar findings were reported by Goto and Shuto (1983); Soares-Frazão and Zech (2008).

In order to validate the scale model approach using Froude similitude it seems practical to check whether threshold limits of Reynolds and Weber number are met. Reynolds number during wave propagation over horizontal bottom yields $Re_{pro} = ch_0/\nu = \sqrt{9.81 \cdot 0.3} \cdot 0.3/10^{-6} = 5.1 \cdot 10^5$. This Reynolds number is well above the threshold value defined by Schüttrumpf (2001). During long wave interaction with beachfront development, it is recommended to refer the calculation of Reynolds number to the width of the MR elements b_{mr} which equals 0.1 m. On the basis of PIV-based velocity measurements inside the MR area local velocities v_{loc} in the range of 0.2 to 0.35 m s⁻¹ are reported (Goseberg, 2011). With reference to Eq. (3), it yields

$$Re_{int} = \frac{v_{loc} b_{mr}}{\nu} = 2.0 \cdot 10^5 - 3.5 \cdot 10^5 \quad (12)$$

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Hence, viscous effects appear to be in-significant during the run-up phase of the long wave since inertial forces outweigh the viscous forces in the down-scaled model. It is also obvious that the Reynolds number is decreasing with decreasing horizontal velocity. This fact also suggests that viscosity develops into a significant parameter only when smaller flow velocities occur. The error stemming from these scale effects is present though negligible. What is more, with the aim to excluding surface tension effects biasing scale model results Novak and Cabelka (1981) state that down-scaled waves should always comprise a minimum wave length of 0.017 m, minimum surface velocities of 0.23 m per s as well as a minimum water depth h_0 of 0.015 m. These conditions are approximately met in this research though at some stages of the wave run-up or draw-down of the long wave, surface velocities or water depth could fall below these limits, while the wave length is always well above. Therefore, scale effects due to viscosity and surface tension appear to be considerably small during most of the tested conditions.

Besides the fundamental flow features that were determined, four different MR element configurations were tested with varying number of obstacle rows. A summary of absolute run-up and differences between those variations is presented in Table 3. It appears that absolute run-up reduction for a given obstruction ratio depends on the kind of MR element arrangement (i.e. aligned or staggered). Aligned and staggered arrangements with identical angle of rotation yield differences of 8.4 to 8.8 % for 5 MR rows and 11.9 to 17.1 % for 10 MR rows. Yet, in case of the aligned configuration ($\Gamma = 1, 2$) the angle of rotation seems to have a greater contribution to the overall run-up reduction compared with the staggered configuration ($\Gamma = 3, 4$) where difference are less dominant. This effect is clearly seen for increasing number of MR element rows. an increasing effect of MR element rotation angle to the overall run-up reduction is also observed with increasing obstruction ratio which can be substantiated by differences for configuration $\Gamma = 1, 2$ between $\psi_{/s} = 0.5, 0.8$. Differences in this case are almost doubled from 12.1 % to 23.3 %. A likely explanation of this fact is that an angle of $\varphi = 45^\circ$ results in much greater local momentum deflection at the MR elements than in non-rotated

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



cases. The momentum deflection in turn leads to greater turbulence generation in the zone where two opposing currents are colliding with each other and therefore a larger amount of energy is dissipated. Due to the greater energy dissipation, less energy is maintained in the run-up tongue resulting in smaller overall run-up values. Hence from the results obtained by an comparison of effects due to angle of rotation and due to MR element configuration it is apparent that the angle of rotation stronger influences the overall run-up reduction than the distinction of aligned or staggered configuration.

Nomograms have been elaborated on the basis of the analysed experimental research. Those nomograms reflect the ability of the tested MR element configurations to reduce the overall run-up of the researched waves for a 1 in 40 sloping beach. Contrary to prototype conditions, beachfront development was idealized by concrete cubic blocks which were arranged in a regular pattern. This approach significantly differs from prototype development in terms of structuring, where houses and infrastructure is mostly randomly distributed. Nevertheless, the roughness effect to an incoming flow which is either exposed to a randomly build or a structured development is maintained by defining cross- and long-shore obstruction ratios. This approach bases of course on the assumption that roughness effects considered over a larger onshore area cancel out local variability in MR element distribution. In order to compare the obtained results with a realistic prototype situation is it thus under this assumptions possible to check for the given obstruction ratios of an area under research. With those obtained values one might consider the nomograms for estimating possible run-up reduction; this method offers a rough estimation at least for beach slopes within the investigated range.

6 Conclusions and outlook

The experiments reported herein address the problem of determining the reduction of maximum tsunami run-up on a 1 in 40 plane beach when the run-up process is disturbed by beachfront development in close proximity of the shoreline. The resistance of urban structures on the beach is modelled by means of MR elements made

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



of cubic concrete blocks. The experimentally determined relative run-up of single sinusoidal waves on the plane beach without MR elements agrees well with an analytical approach by Madsen and Schäffer (2010). For the first time this theoretical approach is successfully validated by experimental means. The relative run-up of the waves under consideration is significantly influenced by the presence of MR elements onshore. As a result, nomograms are presented to facilitate the determination of the relative run-up reduction for four different MR element configurations, various long- and cross-shore obstruction ratios and a range of wave non-linearities. Simple yet idealized predictions of the reduction of long waves due to MR element configurations are therewith made possible. The experimental datasets of long wave run-up are also capable to be utilized for the calibration and validation of numerical models.

It belongs to the character of laboratory idealizations that some aspects of the prototype process are still neglected. Among these aspects are influences stemming from the collapse of buildings, the transport of sediment and debris, the volumetric disequilibrium between solid blocks (in the present model) and common hollow buildings (in prototype), three-dimensional influences of river estuaries or heterogeneous existing developments, different building orientations and dimensions, the overflow of developments and spatially or temporally varying surface gradients. Moreover, the presented material is believed to owe the potential to facilitate future modelling approaches of beachfront development, though it is beyond the intention of this paper. The above mentioned aspects remain untreated and further laboratory work is needed to elucidate those additions to the interaction process of long waves and beachfront development.

Appendix A

Additional nomograms

Some more nomograms are provided in this section that add information about the remaining MR configurations $\Gamma = 1, 3$ for the sake of completeness.

Supplementary material related to this article is available online at:
[http://www.nat-hazards-earth-syst-sci-discuss.net/1/1119/2013/
nhessd-1-1119-2013-supplement.zip](http://www.nat-hazards-earth-syst-sci-discuss.net/1/1119/2013/nhessd-1-1119-2013-supplement.zip).

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NHESD

1, 1119–1171, 2013

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



- Goseberg, N.: Experimental run-up determination of single sinusoidal, solitary and N-waves, in: Proceedings of the International Offshore and Polar Engineering Conference, in press, 2013. 1133
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**Reduction of
maximum tsunami
run-up**

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



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NHESSD

1, 1119–1171, 2013

Reduction of maximum tsunami run-up

N. Goseberg

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

Table 1. List of sinusoidal waves used, values refer to measurements at the wavemaker, water depth h is 0.31 m, period T is 60 s, linear wave length L_0 becomes 104.63 m, ξ_0 surf similarity, no. of repetitions is given under “Rep”.

No. ID	a_{cr} [-]	H_0 [m]	a_{cr}/h_0 [s]	H_0/L_0 [m]	ξ_0 [-]	Rep. [-]
1	0.016	0.031	0.052	2.963×10^4	10.49	2
2	0.018	0.036	0.058	3.441×10^4	9.79	39
3	0.020	0.041	0.064	3.918×10^4	9.10	2
4	0.023	0.046	0.074	4.396×10^4	8.57	2
5	0.026	0.052	0.083	4.970×10^4	8.06	2
6	0.030	0.059	0.097	5.639×10^4	7.69	31

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

Table 2. Number of conducted experiments classified according to the utilized wave properties listed in Table 1 and to the MR combinations Γ , “al” aligned, “st” staggered.

ID	w/o MR	al	al	st	st	Σ
		$\varphi = 0^\circ$ $\Gamma = 1$	$\varphi = 45^\circ$ $\Gamma = 2$	$\varphi = 0^\circ$ $\Gamma = 3$	$\varphi = 45^\circ$ $\Gamma = 4$	
1	2	12	9	8	6	37
2	39	12	9	8	6	74
3	2	12	9	8	6	37
4	2	12	9	8	6	37
5	2	12	9	8	6	37
6	31	13	10	8	6	68
Σ	78	73	55	48	36	290

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

Table 3. Summary of normalized run-up and relative differences with regard to configuration $\Gamma = 1$ for MR element rows with $n = 5; 10$, listed values base on the polynomial fitting curves.

Γ	n	$R_{\text{mean}}/a_{\text{cr}}$ $\psi_{Is} = 0.5$	Δ [%]	$R_{\text{mean}}/a_{\text{cr}}$ $\psi_{Is} = 0.8$	Δ [%]
1	5	3.30	–	2.80	–
2	5	2.90	12.1	2.15	23.3
3	5	3.01	8.8	2.32	17.1
4	5	2.83	14.2	2.20	21.4
1	10	3.20	–	2.70	–
2	10	2.87	10.3	2.18	19.3
3	10	2.93	8.4	2.38	11.9
4	10	2.78	13.1	2.17	19.6

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

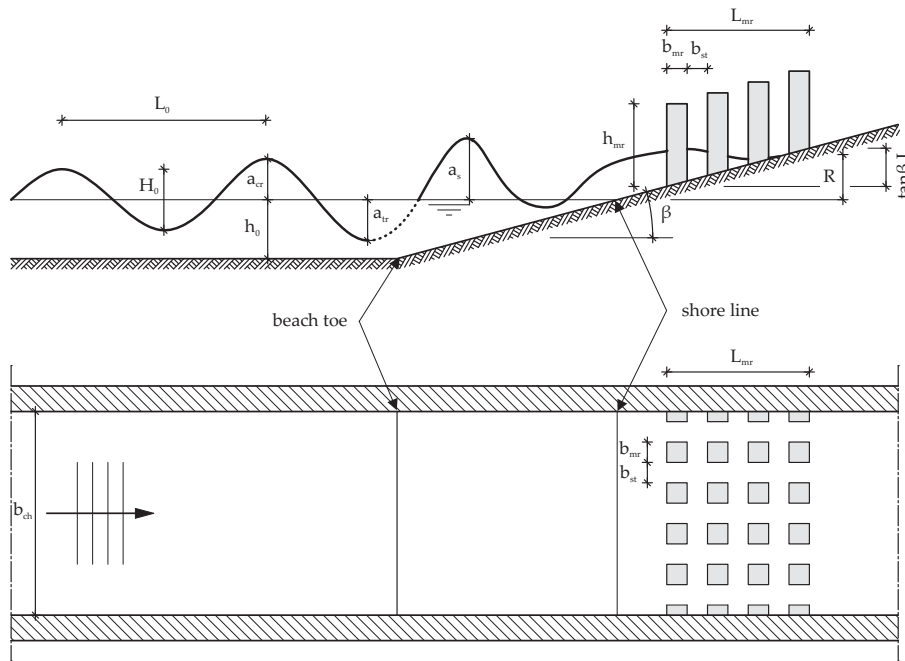


Fig. 1. Schematic sketch of the problem set-up and the definition of the used variables, h_0 water depth, H_0 wave height and L_0 wave length at the even bottom section, a_{cr} and a_{tr} amplitude of the wave crest and trough, β slope angle, R wave run-up, b_{mr} and h_{mr} width and height of the macro-roughness elements, b_{st} distance between the macro-roughness elements. Spatial extent of the shown wave is not to scale. Experimentally investigated wave length significantly exceeds beach length.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

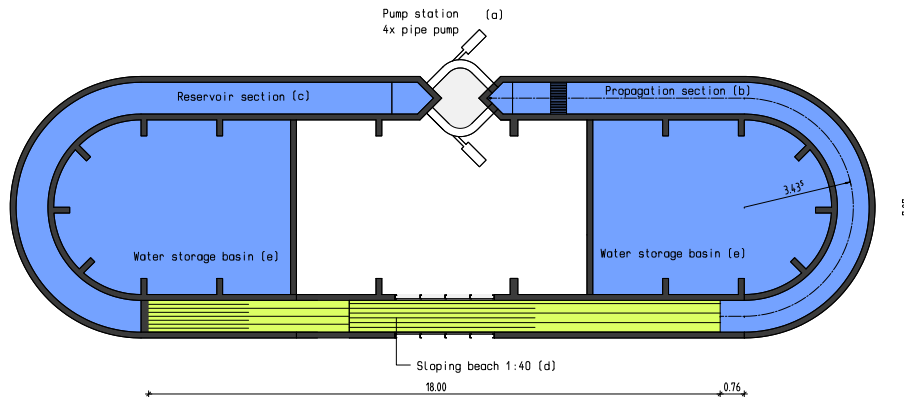


Fig. 2. Schematic drawing of the closed-circuit flume and its equipments, **(a)** pump station, **(b)** propagation section, **(c)** reservoir section, **(d)** sloping beach, **(e)** water storage basin (units in meter).

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

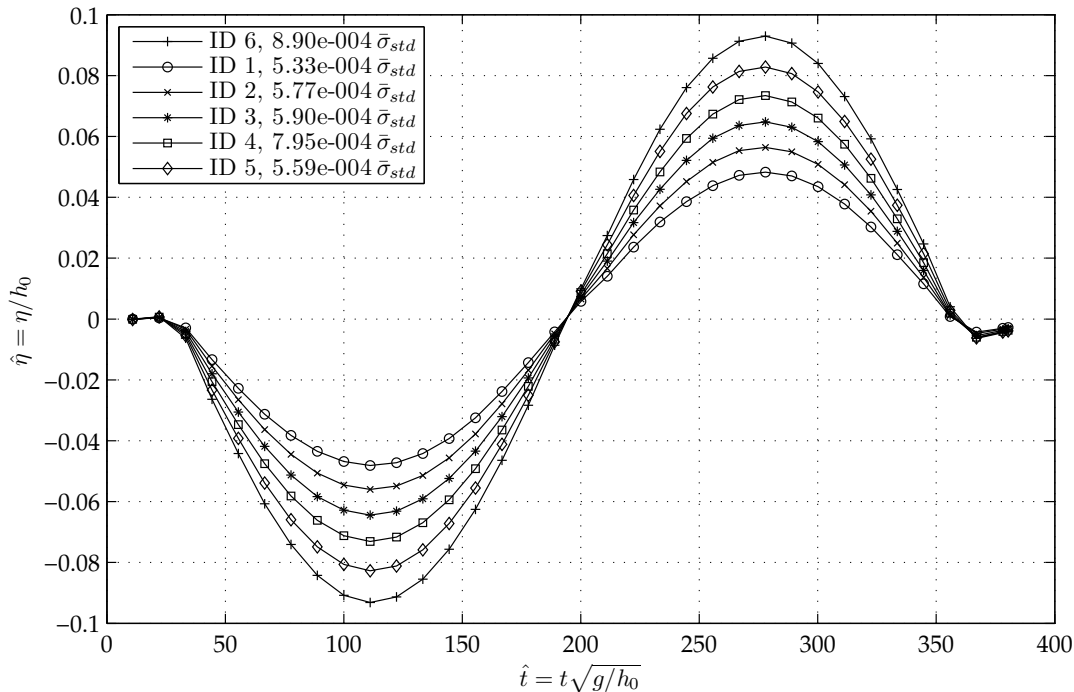


Fig. 3. Averaged non-dimensional surface elevation history for the six utilized waves, the markers denote the different waves heights, $\bar{\sigma}_{std}$ is the averaged standard deviation calculated from the entire number of experiments in each group.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

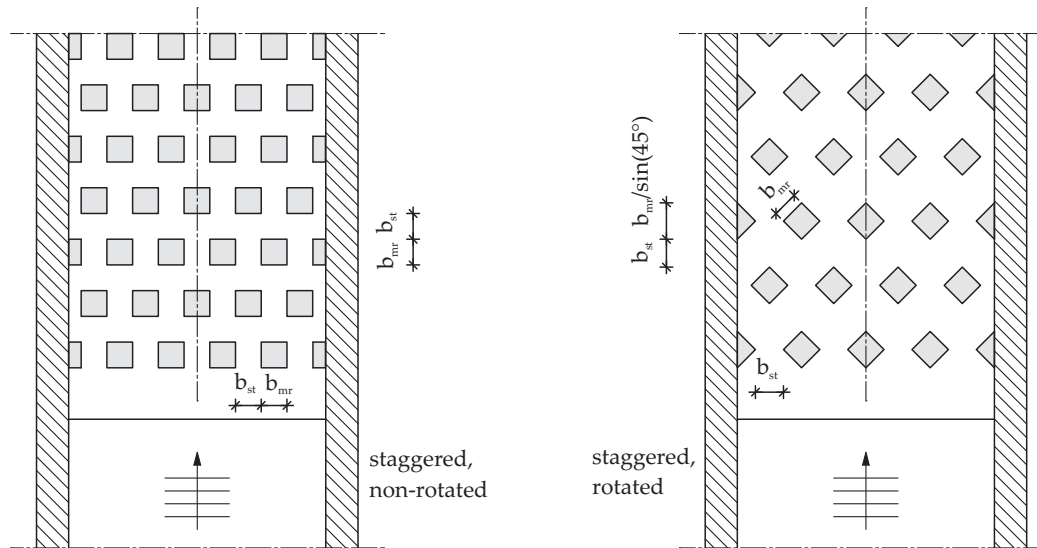


Fig. 4. Symmetry requirements for the configuration of the MR elements for a staggered, non-rotated configuration (left) and a staggered rotated configuration (right), drawing not to scale, demonstration of the axis of symmetry at the flume wall, which always cuts an MR element in the half.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



**Reduction of
maximum tsunami
run-up**

N. Goseberg



Fig. 5. Photographs of the MR element implementation. Left: Example of a staggered MR element configuration, Right: Half cubes used at the flume walls.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

I◀

▶I

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

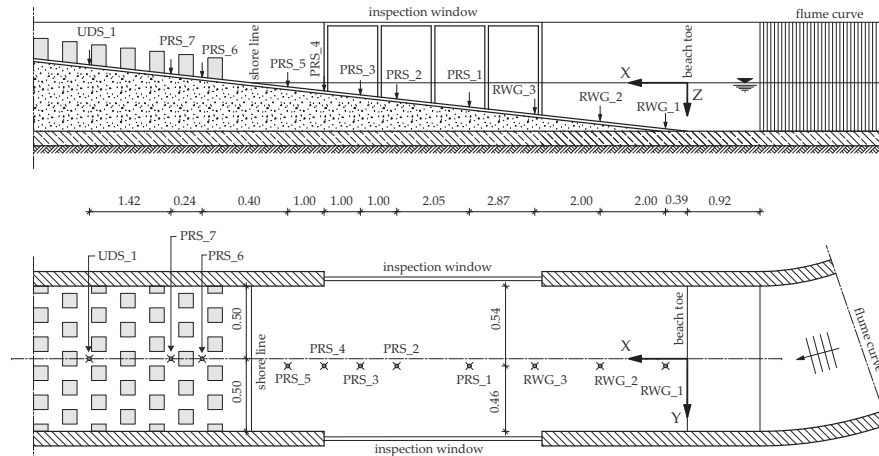


Fig. 6. Layout drawing of the laboratory instrumentation with used pressure sensors and wave gauges for the experimental set-up, abbreviations denote the following devices: PRS – pressure sensor, RWG – resistance wave gauge, UDS – ultrasonic distance sensor; for more than one device consecutive numbers are given, not to scale (units in meter), note that sensors are not analysed in this paper.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



**Reduction of
maximum tsunami
run-up**

N. Goseberg

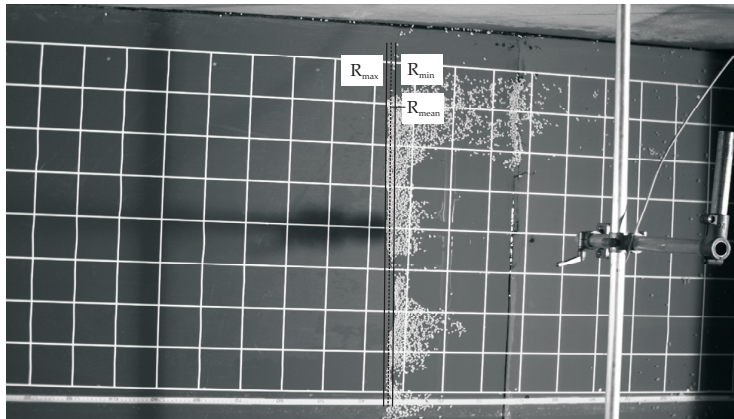


Fig. 7. Maximum wave run-up detection in the laboratory defined as the mean R_{mean} of the lowest R_{min} and the highest maximum run-up R_{max} .

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

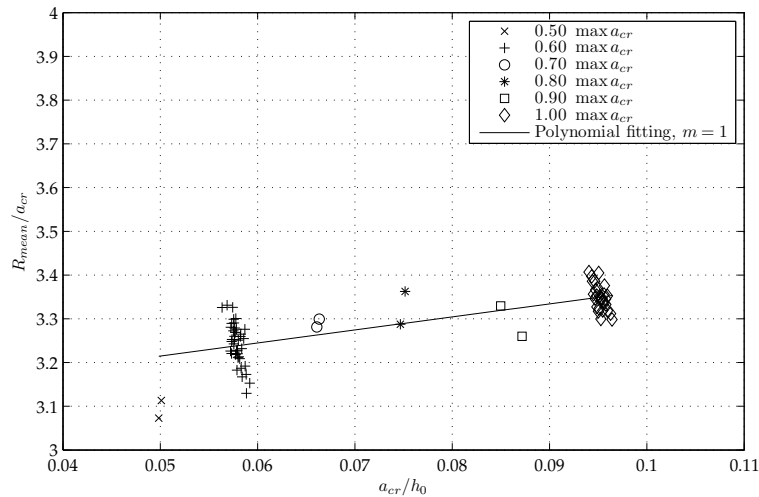


Fig. 8. Maximum relative run-up as a function of the nonlinearity ϵ of the waves depicted by the markers, a linear polynomial fit denotes the increase in maximum run-up with increasing nonlinearity

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

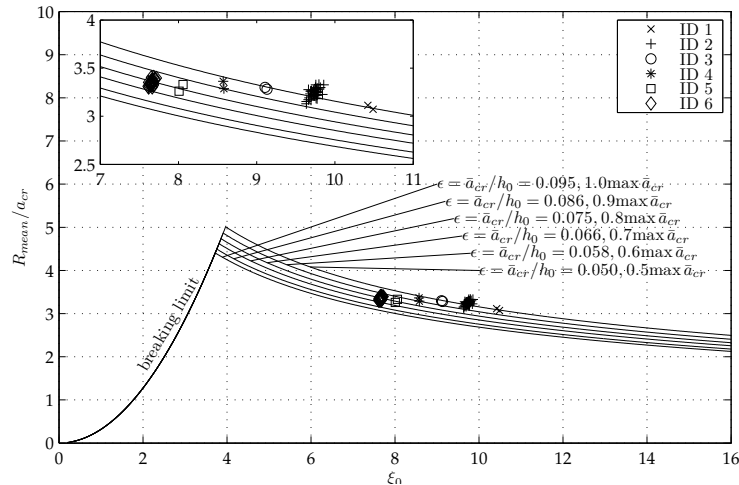


Fig. 9. Maximum relative run-up as a function of the surf similarity parameter ξ_0 depicted by the markers and according to the wave conditions applied, the theoretical breaking limit given by Eq. 10 and the theoretical run-up functions given by Eq. 8 (Madsen and Schäffer, 2010) based on the non-linearity ϵ . The inset zooms to the experimental data.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

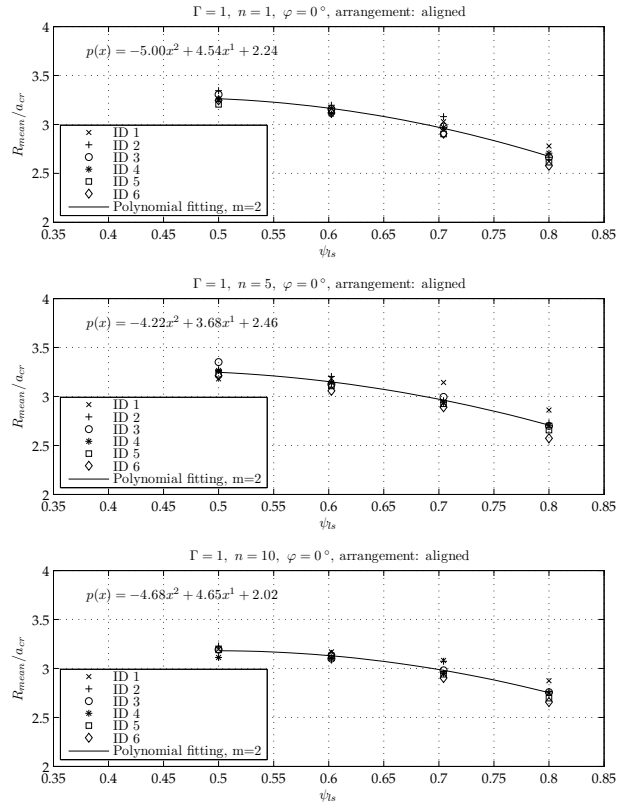


Fig. 10. The relative wave run-up R_{mean} as a function of the long-shore obstruction ratio ψ_{cs} , $\Gamma = 1$. The experimental data are fitted to a polynomial of degree $m = 2$ and the equations are given according to the investigated number of MR rows $n = 1, 5, 10$

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

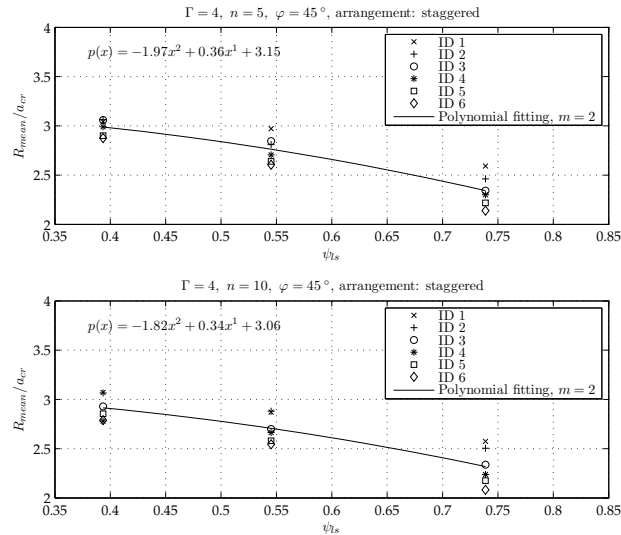


Fig. 11. The relative wave run-up R_{mean} as a function of the long-shore obstruction ratio ψ_{ls} , $\Gamma = 4$. The experimental data are fitted to a polynomial of degree $m = 2$ and the equations are given according to the investigated number of MR rows $n = 5, 10$

Title Page

Abstract Introduction

Conclusions References

Tables Figures

◀ ▶

◀ ▶

Back Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

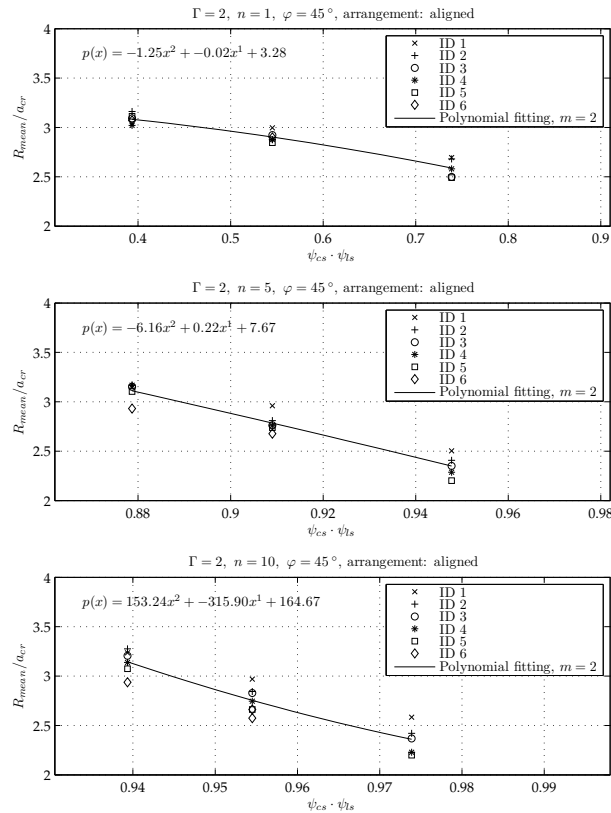


Fig. 12. The relative wave run-up R_{mean} as a function of the cross-shore obstruction ratio ψ_{CS} times the long-shore obstruction ratio ψ_{LS} , $\Gamma = 2$. The experimental data are fitted to a polynomial of degree $m = 2$ and the equations are given according to the investigated number of macro-roughness rows $n = 1, 5, 10$

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

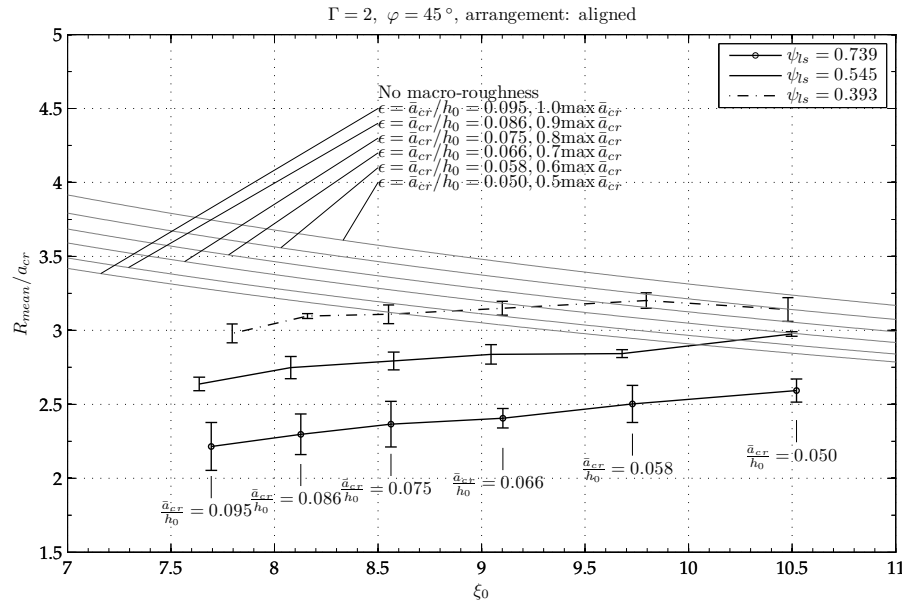


Fig. 13. The relative wave run-up $R_{\text{mean}}/a_{\text{cr}}$ as a function of the surf similarity parameter ξ_0 , $\Gamma = 2$. The experimental data are plotted according to the individual long-shore obstruction ratio. The error bars indicate the variation of the number of macro-roughness element rows. The grey parallel graphs indicate the relative run-up without the interaction with macro-roughness. The non-linearity of the applied waves is indicated as text which always refers to all error bars in the vertical above.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

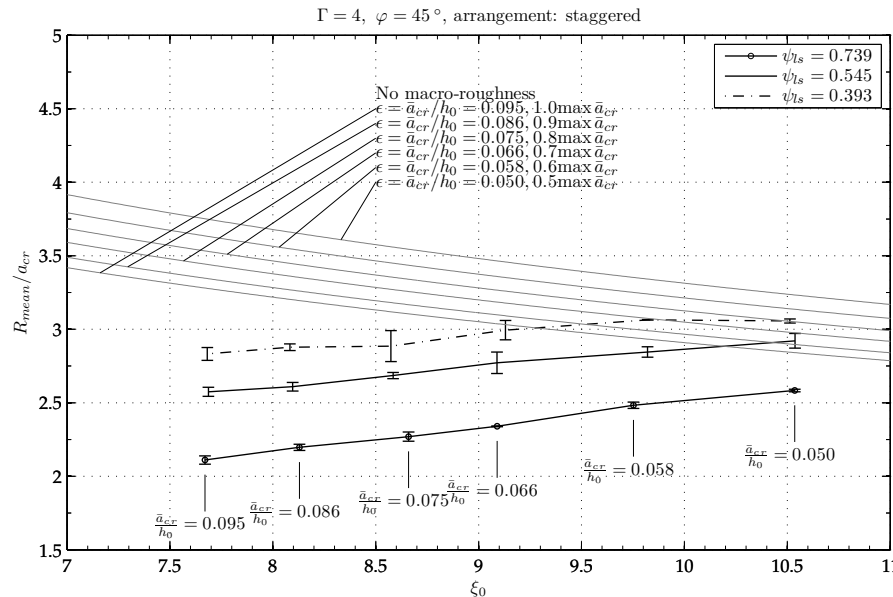


Fig. 14. The relative wave run-up $R_{\text{mean}}/a_{\text{cr}}$ as a function of the surf similarity parameter ξ_0 , $\Gamma = 4$. The experimental data are plotted according to the individual long-shore obstruction ratio. The error bars indicate the variation of the number of macro-roughness element rows. The grey parallel graphs indicate the relative run-up without the interaction with macro-roughness. The non-linearity of the applied waves is indicated as text which always refers to all error bars in the vertical above.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

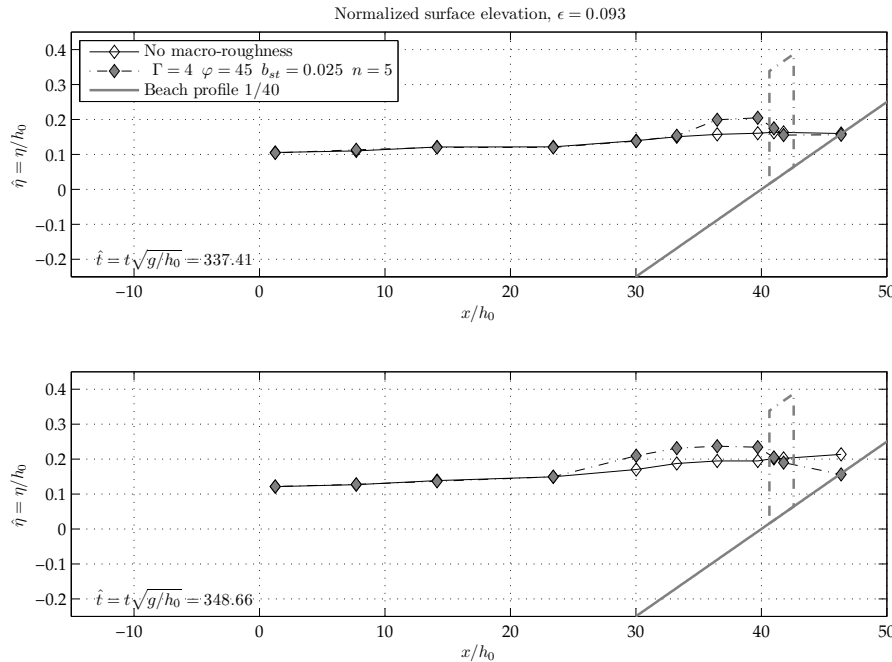


Fig. 15. Comparison of the normalized surface elevation for a sinusoidal wave with $\epsilon = 0.093$ at $\hat{t} = 337.41$ and $\hat{t} = 348.66$ with and without the presence of macro-roughness elements. The macro-roughness element configuration is specified by $\Gamma = 4$, $\varphi = 45^\circ$, $b_{st} = 0.025$ and $n = 5$. Surface elevation and spatial coordinates are normalized by the invariant water depth h_0 . The dash-dotted line depicts the macro-roughness element area.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

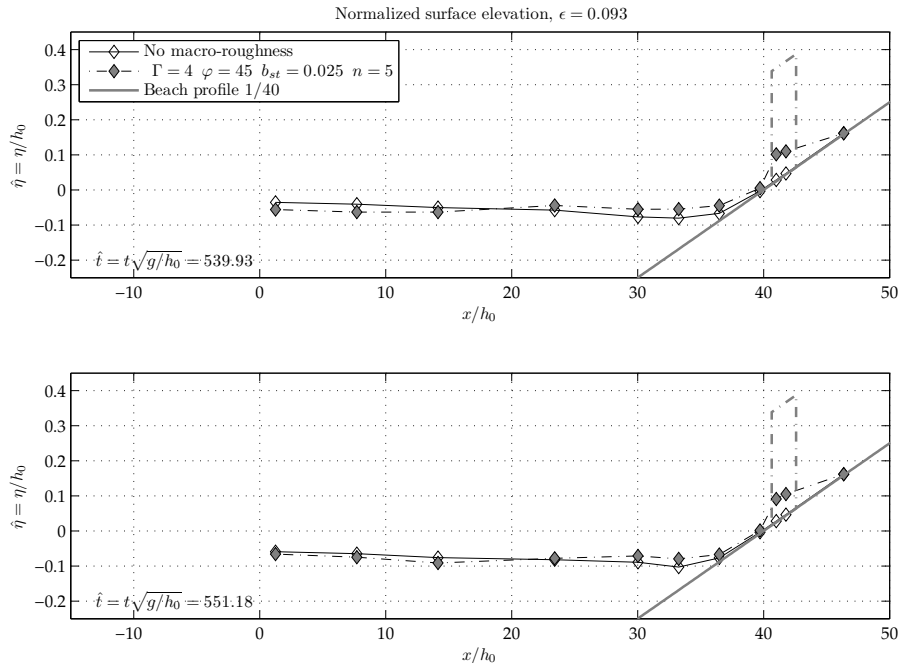


Fig. 16. Comparison of the normalized surface elevation for a sinusoidal wave with $\epsilon = 0.093$ at $\hat{t} = 539.93$ and $\hat{t} = 551.18$ with and without the presence of macro-roughness elements. The macro-roughness element configuration is specified by $\Gamma = 4$, $\varphi = 45^\circ$, $b_{st} = 0.025$ and $n = 5$. Surface elevation and spatial coordinates are normalized by the invariant water depth h_0 . The dash-dotted line depicts the macro-roughness element area.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

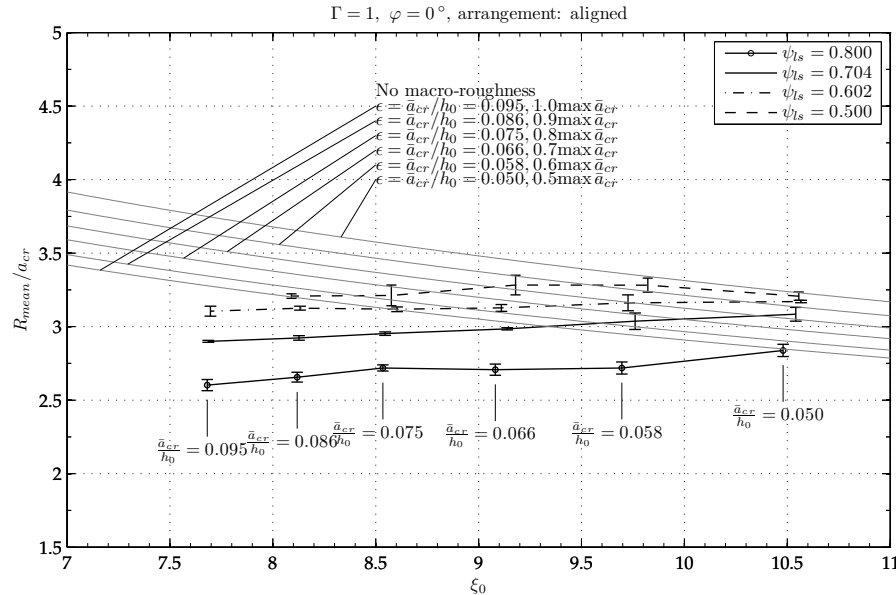


Fig. A1. The relative wave run-up $R_{\text{mean}}/a_{\text{cr}}$ as a function of the surf similarity parameter ξ_0 , $\Gamma = 1$. The experimental data are plotted according to the individual long-shore obstruction ratio. The error bars indicate the variation of the number of macro-roughness element rows. The grey parallel graphs indicate the relative run-up without the interaction with macro-roughness. The non-linearity of the applied waves is indicated as text which always refers to all error bars in the vertical above.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion



Reduction of maximum tsunami run-up

N. Goseberg

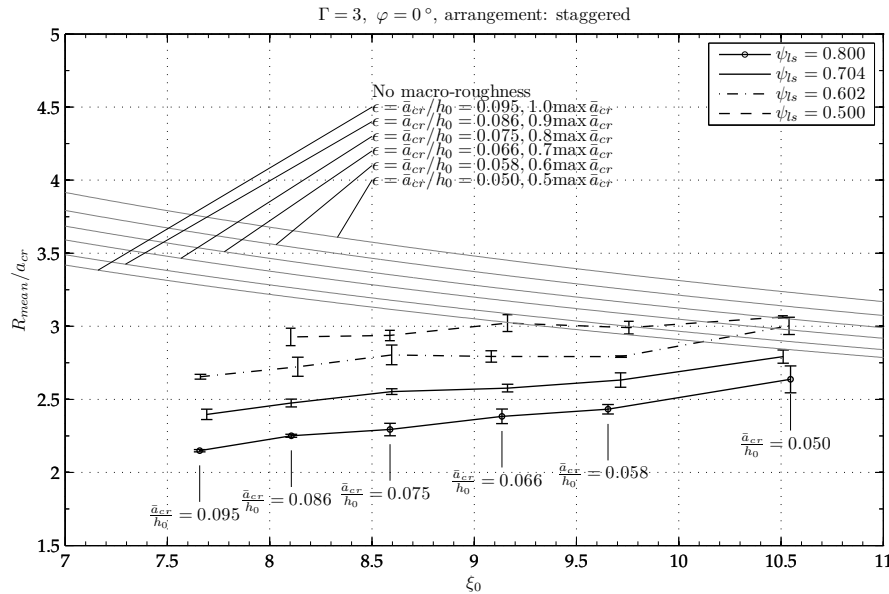


Fig. A2. The relative wave run-up $R_{\text{mean}}/a_{\text{cr}}$ as a function of the surf similarity parameter ξ_0 , $\Gamma = 3$. The experimental data are plotted according to the individual long-shore obstruction ratio. The error bars indicate the variation of the number of macro-roughness element rows. The grey parallel graphs indicate the relative run-up without the interaction with macro-roughness. The non-linearity of the applied waves is indicated as text which always refers to all error bars in the vertical above.

Title Page

Abstract

Introduction

Conclusions

References

Tables

Figures

◀

▶

◀

▶

Back

Close

Full Screen / Esc

Printer-friendly Version

Interactive Discussion

